

CHAPTER 8: MECHANICAL FAILURE



ISSUES TO ADDRESS

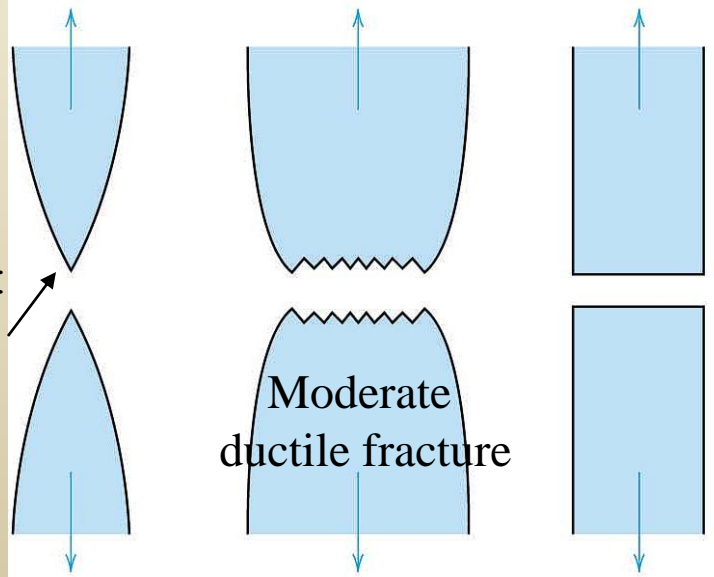
- How do **flaws in a material** initiate failure?
- How is fracture resistance quantified; how do different **material classes compare**?
- How do we estimate the ***stress to fracture***?
- How do **loading rate, loading history, and temperature** affect the failure stress?

DUCTILE VS BRITTLE FAILURE

Simple fracture is the separation of a body into two or more pieces in response to an imposed static stress at temperatures that are lower than material's melting point.

Two limiting fracture **modes** are possible for engineering materials: **ductile** and **brittle**

Highly ductile fracture:
neck down to a point

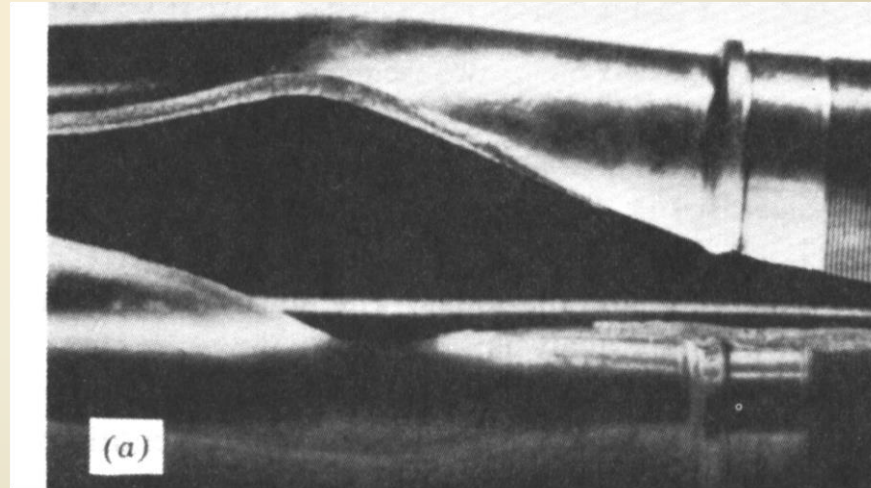


Absolute brittle fracture:
without any plastic deformation

Example: FAILURE OF A PIPE

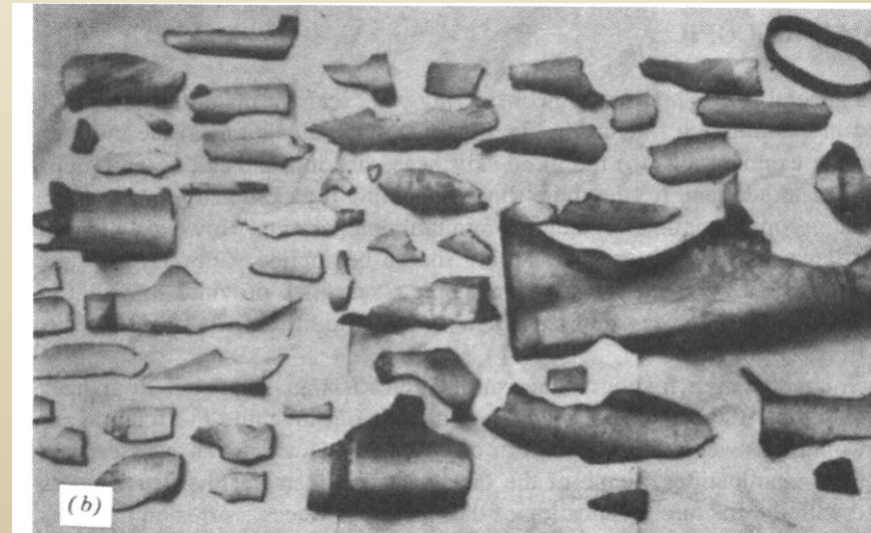
- **Ductile failure:**

- one piece
- large deformation



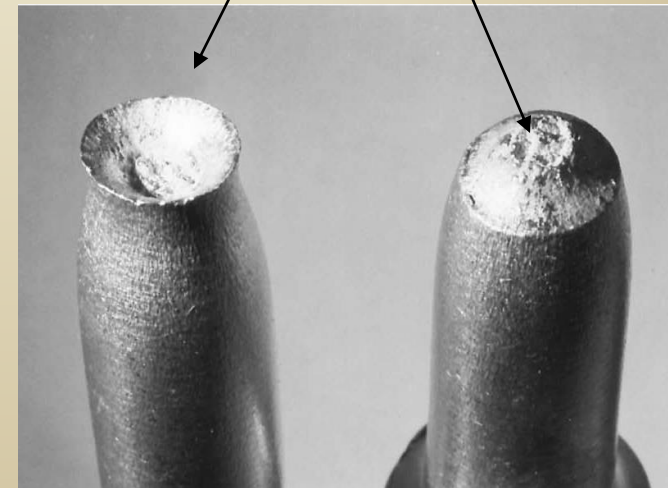
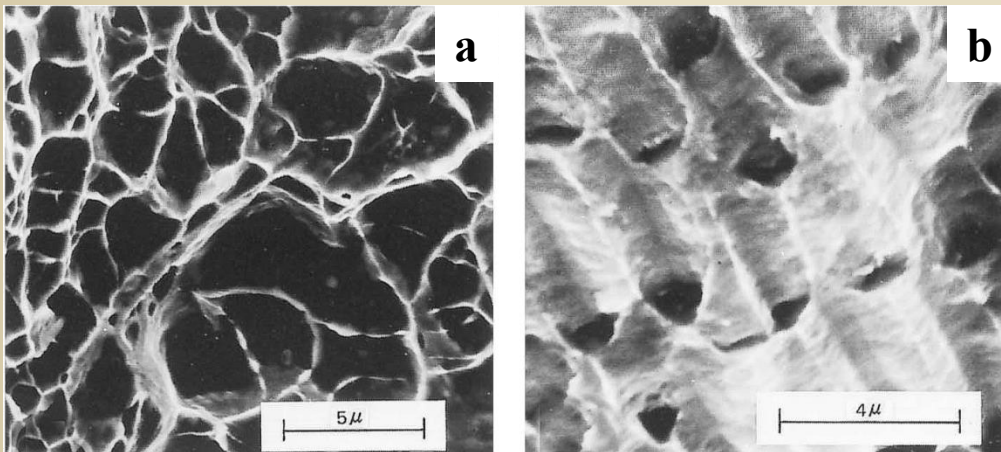
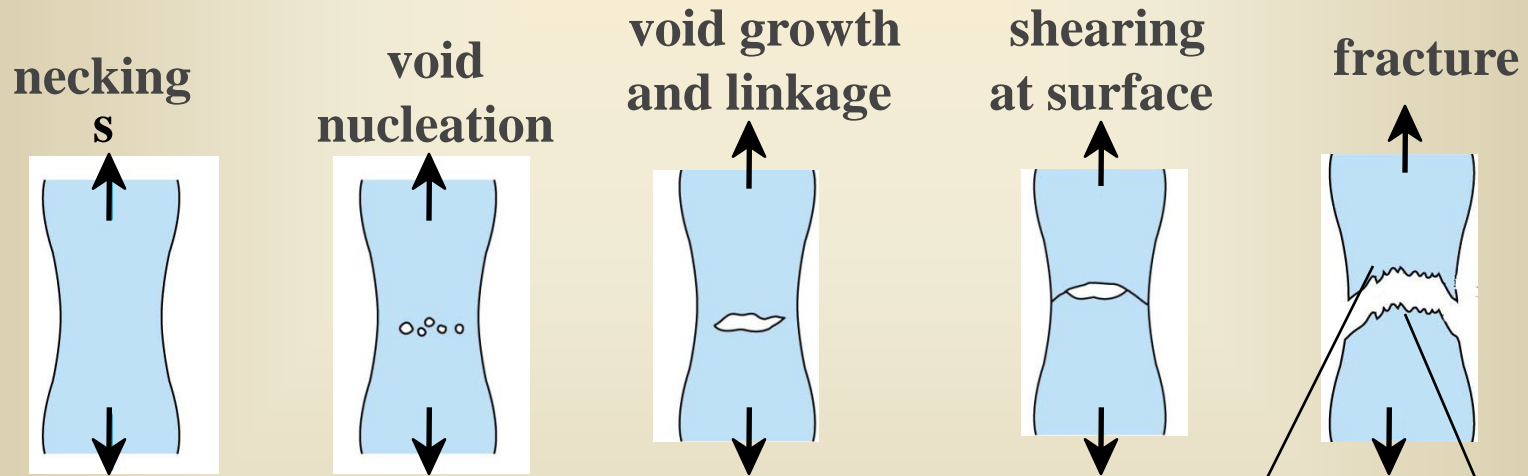
- **Brittle failure:**

- many pieces
- small deformation



MODERATELY DUCTILE FAILURE

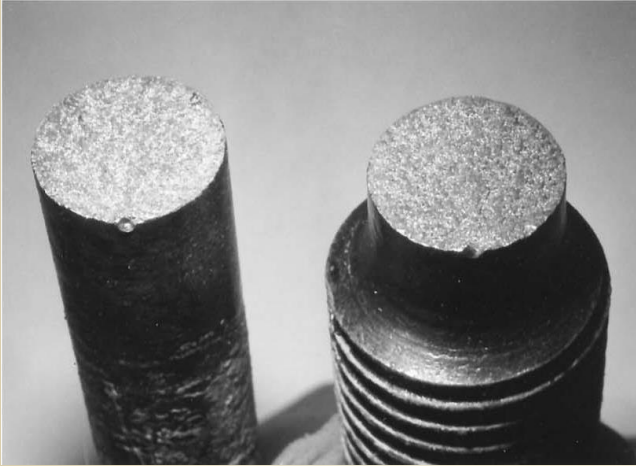
- Evolution of **cup-and-cone** fracture to failure:



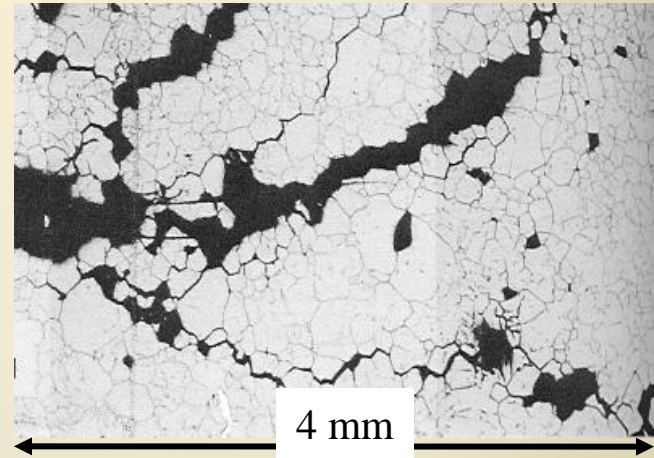
Spherical (a) and parabolic (b) “dimples” characteristic for ductile fracture from uniaxial tensile and shear loadings

Cup-and cone fracture of Al

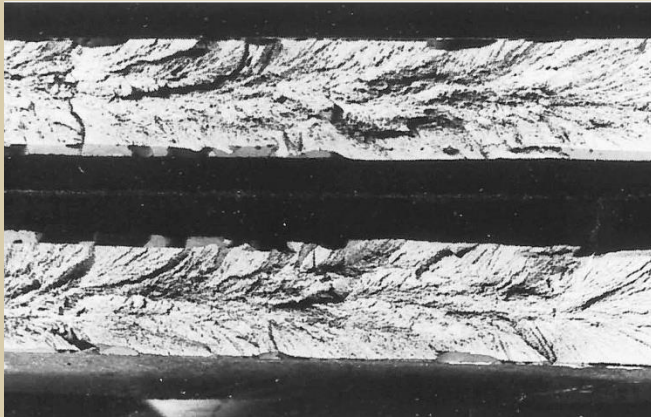
BRITTLE FAILURE



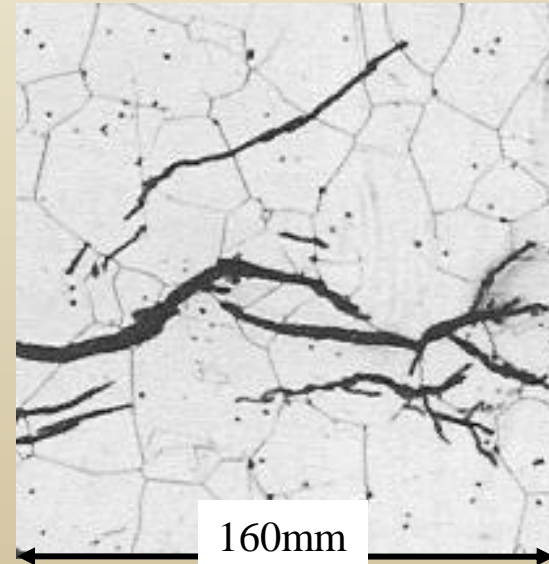
Brittle fracture in a mild steel



Intergranular fracture (between grains)



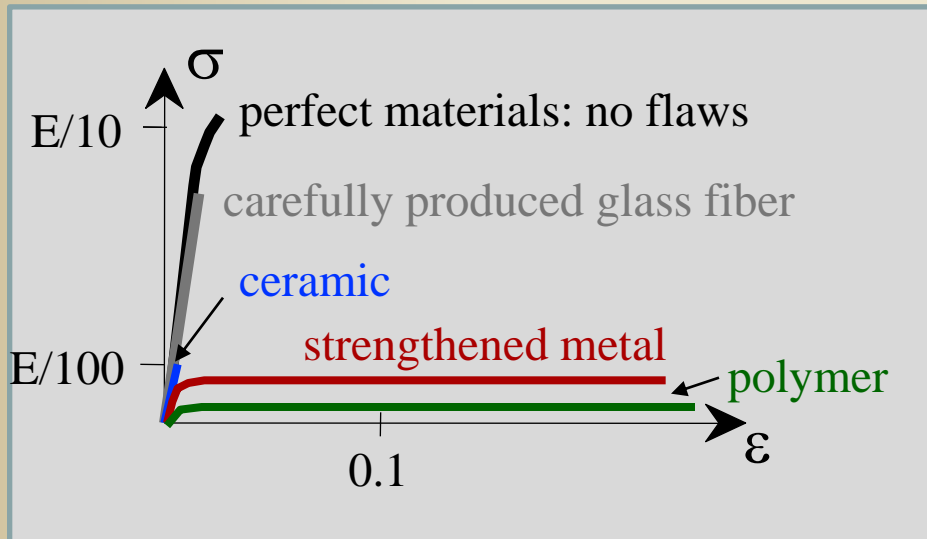
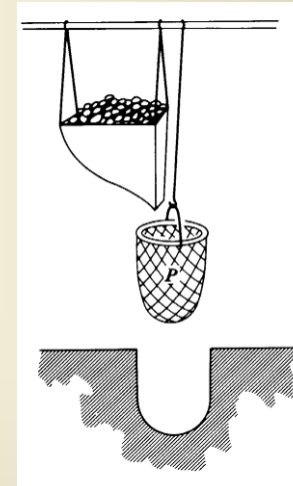
V-shaped “chevron” markings
Characteristic of brittle fracture



Intragranular fracture (within grains)

Again: IDEAL vs. REAL MATERIALS

- **History:** **Leonardo DaVinci** (500 years ago!) observed: the **longer the wire**, the **smaller the failure-causing load**.



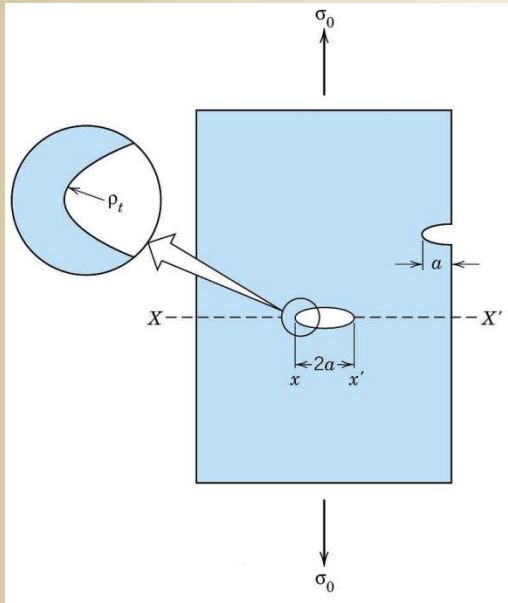
- **Reason:**

- **flaws** cause premature failure.
- larger samples are more flawed!

$$TS_{\text{Eng. Mater.}} \ll TS_{\text{Perfect Mater}}$$

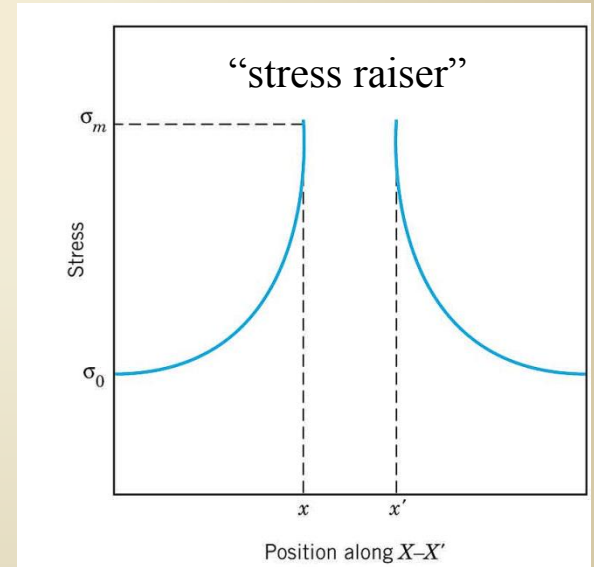
FLAWS ARE STRESS CONCENTRATORS!

- Schematic of surface and internal cracks
- Stress distribution in front of a crack:



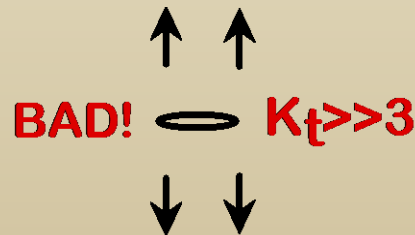
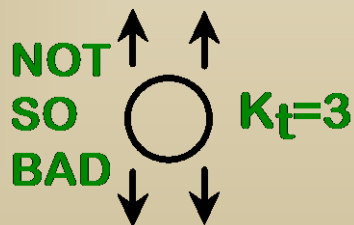
Characterized by:
length ($2a$)
curvature (ρ_t)

$$\sigma_m = 2\sigma_0 \left(\frac{a}{\rho_t} \right)^{1/2}$$



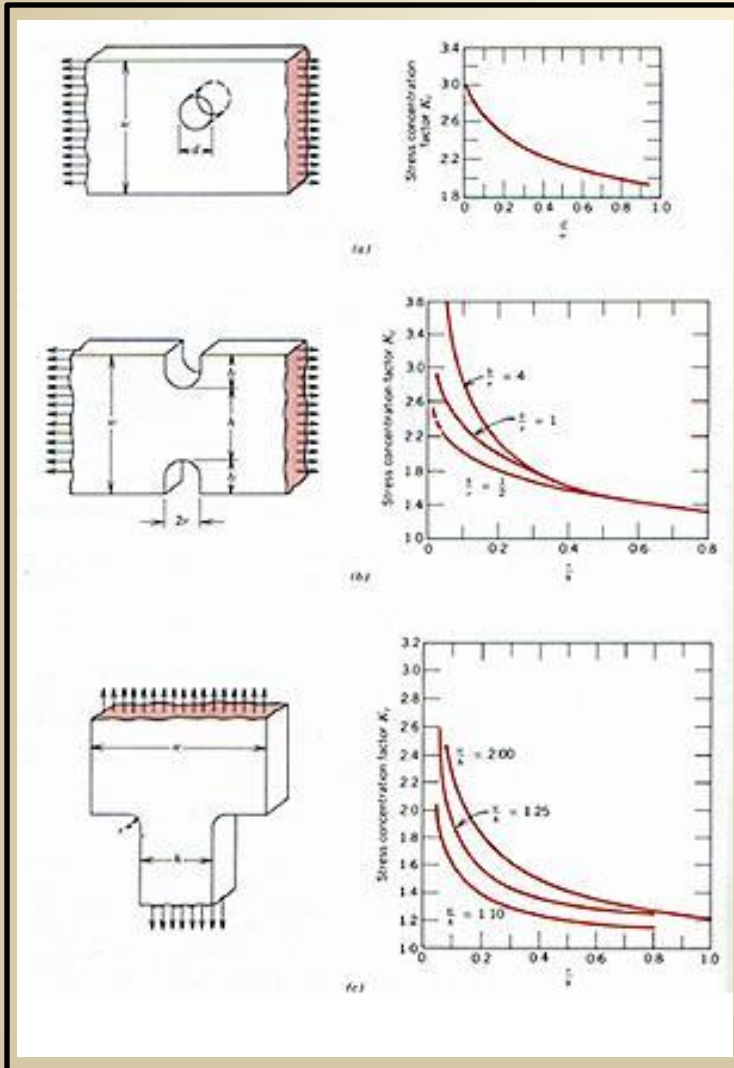
- Stress concentration factor: $K_t = \sigma_{\max} / \sigma_0$ - stress amplification!

Large K_t promotes failure:



Plastic deformation ($\sigma_m > \sigma_y$)
leads to more uniform
stress distribution!!

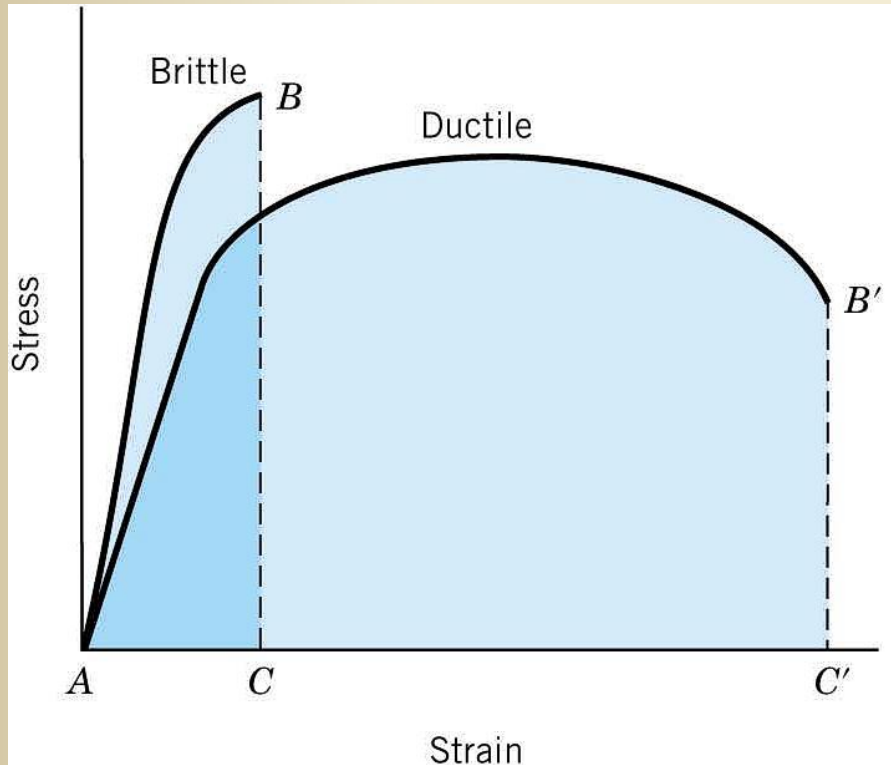
STRESS CONCENTRATORS: Macroscopic Level



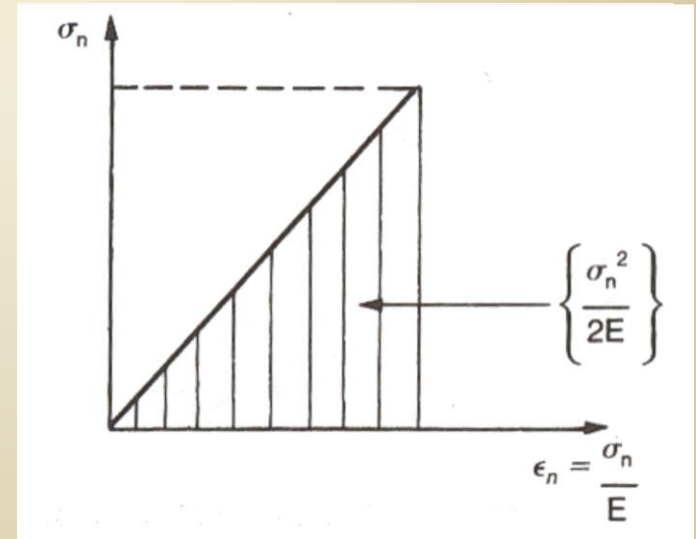
- It is important to remember that stress amplification not only occurs on a **microscopic level** (e.g. small flaws or cracks,) but can also take place on **the macroscopic level** in the case of sharp corners, holes, fillets, and notches.
- The figure depicts the theoretical stress concentration factor curves for several simple and common material geometries.

Reminder: TOUGHNESS

- Energy to break a unit volume of material
- Approximate by the area under the stress-strain curve.
- Units: [J/m²]



For Elastic Strain

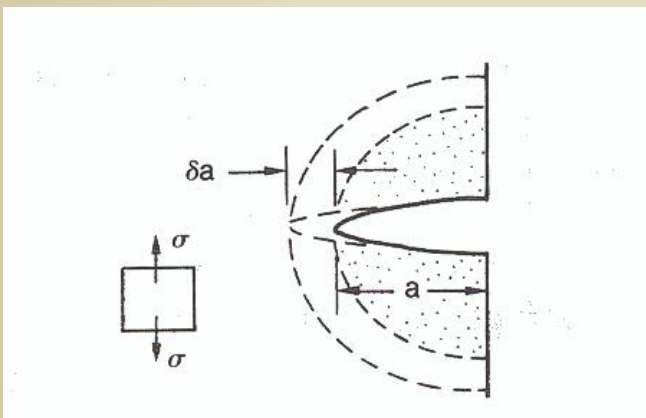
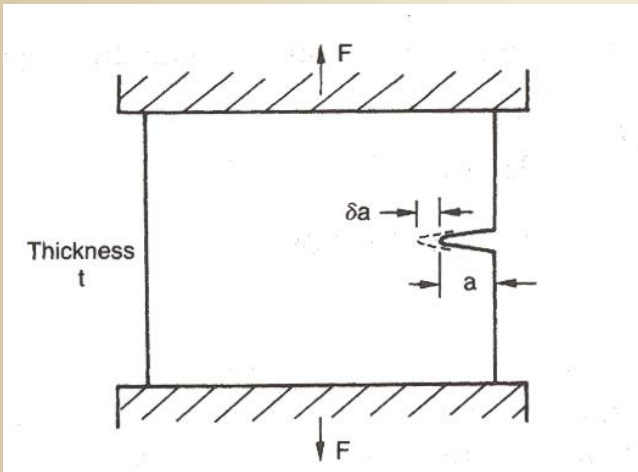


$$\sigma_n = E\epsilon_n$$

$$U^{el} = \int \sigma_n d\epsilon_n = \int \sigma_n \frac{d\sigma_n}{E} = \left\{ \frac{\sigma_n^2}{2E} \right\}$$

Material Fast Fracture

Fixed displacements case
(the boundary of the plate are fixed)



$$\delta W \geq \delta U^{el} + \gamma_c t \delta a$$

$$-\delta U^{el} = \gamma_c t \delta a$$

$$U^{el} = -\frac{\sigma^2 \pi a^2 t}{2E}$$

$$\delta U^{el} = \frac{\delta U^{el}}{\delta a} \delta a = -\frac{\sigma^2}{2E} \frac{2\pi a t}{2} \delta a$$

$$\frac{\sigma^2 \pi a}{2E} = \gamma_c$$

$$\sigma_c = \left(\frac{2E\gamma_c}{\pi a} \right)^{1/2}$$

where E – module of elasticity; σ_c – critical stress;
 γ_s – **toughness** (J m^{-2} , specific surface energy);
 W -work; U^{el} - energy related to elastic deformation

1: FRACTURE TOUGHNESS

- Thus it is possible to show that critical stress for crack propagation is:

Stress intensity factor \rightarrow $\sigma\sqrt{\pi a} = \sqrt{E\gamma_c}$ \leftarrow Material properties only!!!

- **Fracture toughness** is a property that is a measure of a materials' resistance for brittle fracture when cracks are present. This property can be defined by parameter K_c that relates the critical stress for crack propagation and geometry of the crack:

$$K_c = Y\sigma_c\sqrt{\pi a}$$

$$K_{ic} = [\text{Pa} \cdot \text{m}^{1/2}]!!!$$

where **Y** – a dimensionless parameter that depends on both crack and specimen sizes and geometries (is tabulated for different crack-specimen geometries), as well as the type of load application

GEOMETRY, LOAD, & MATERIAL

- Condition for crack propagation:

Stress Intensity Factor: \longrightarrow

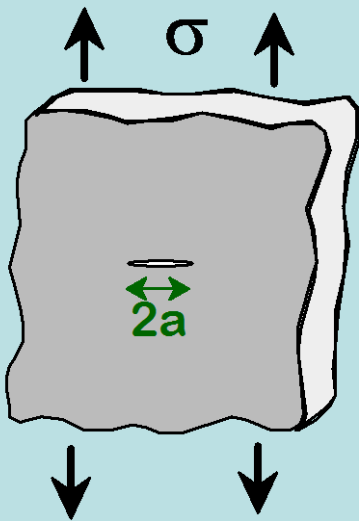
$$K \geq K_c$$

-Depends on load & geometry

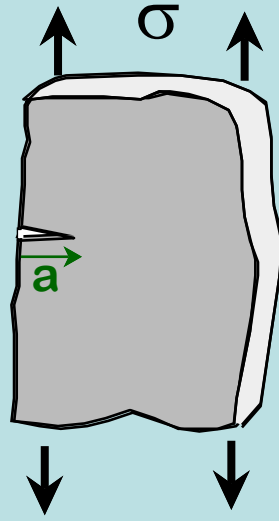
Fracture Toughness:

-Depends on the material, temperature, environment, & rate of loading.

- Values of K for some standard loads & geometries:



$$K = \sigma \sqrt{\pi a}$$



$$K = 1.1 \sigma \sqrt{\pi a}$$

When sample thickness \gg crack dimensions

plane-strain conditions occur,

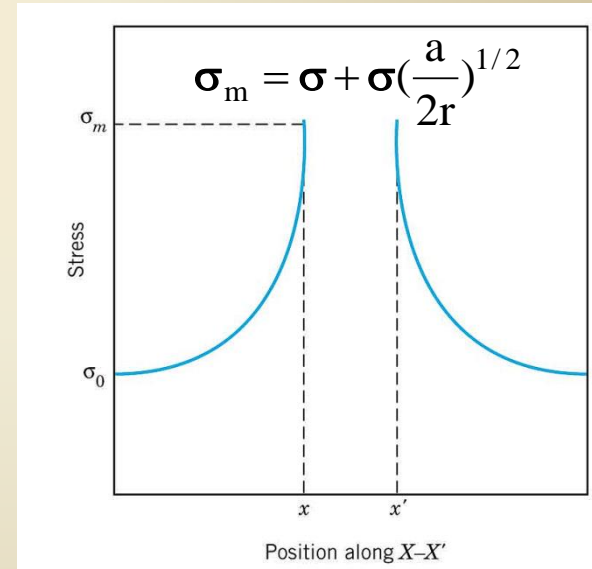
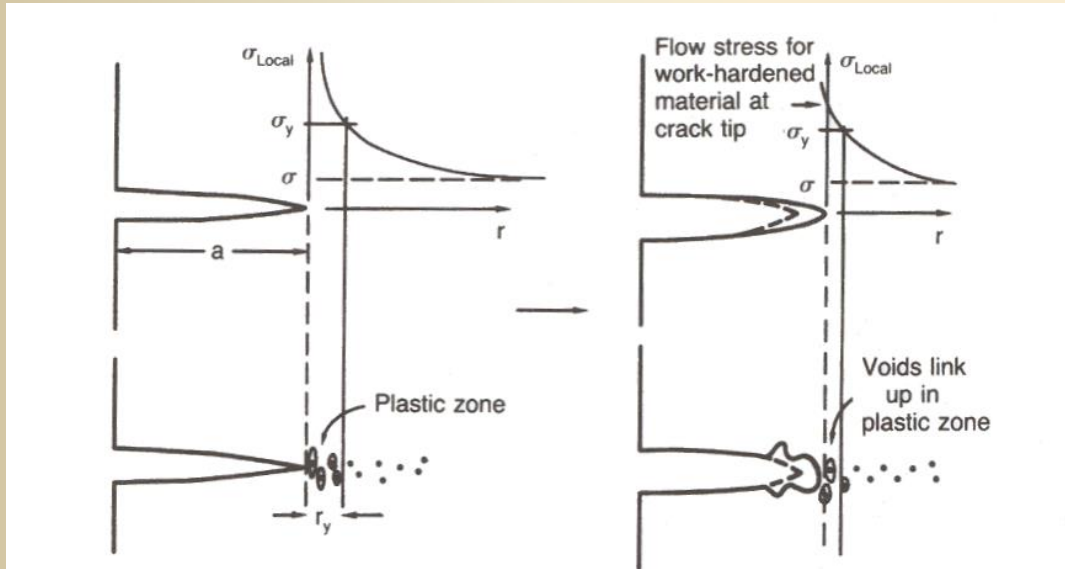
$$K_c = K_{1c}$$

that is independent of thickness

$$K_{1c} = Y \sigma_c \sqrt{\pi a}$$

Plane strain fracture toughness

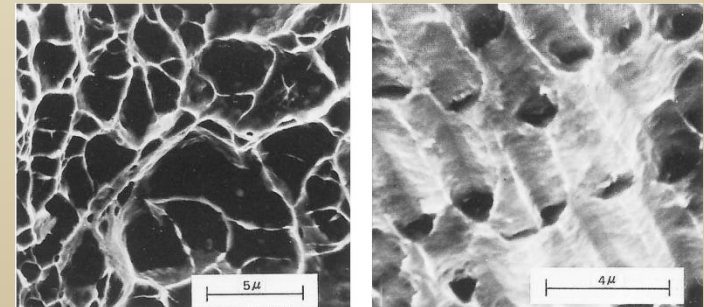
Mechanisms of Crack Propagation: ductile tearing



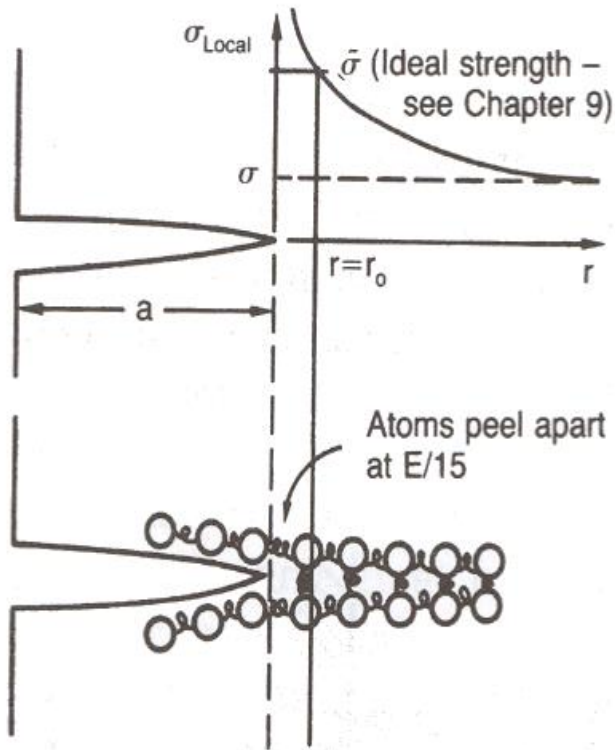
The plastic flow at the crack tip naturally turns the initially sharp crack into a **blunt crack**. Crack blunting decreases σ_m so that crack tip itself can keep on plastically deforming. Thus ductile tearing **consumes a lot of energy by plastic flow**. This is why ductile materials are so **tough**.

$$r_y = \frac{\sigma^2 a}{2\sigma_y^2} = \frac{K^2}{2\pi\sigma_y^2}$$

$$K = K_c$$



Mechanisms of Crack Propagation: cleavage



Blunting of the sharp crack does not occur/
The local stress at the crack tip is large enough to **break apart the inter-atomic bonds!!**

The crack spreads between a pair of atomic planes leading to the formation of flat surfaces by **cleavage**.

$$r = \frac{\sigma^2 a}{2\sigma_y^2} = \frac{K^2}{2\pi\sigma_y^2}$$

$$K = K_c$$

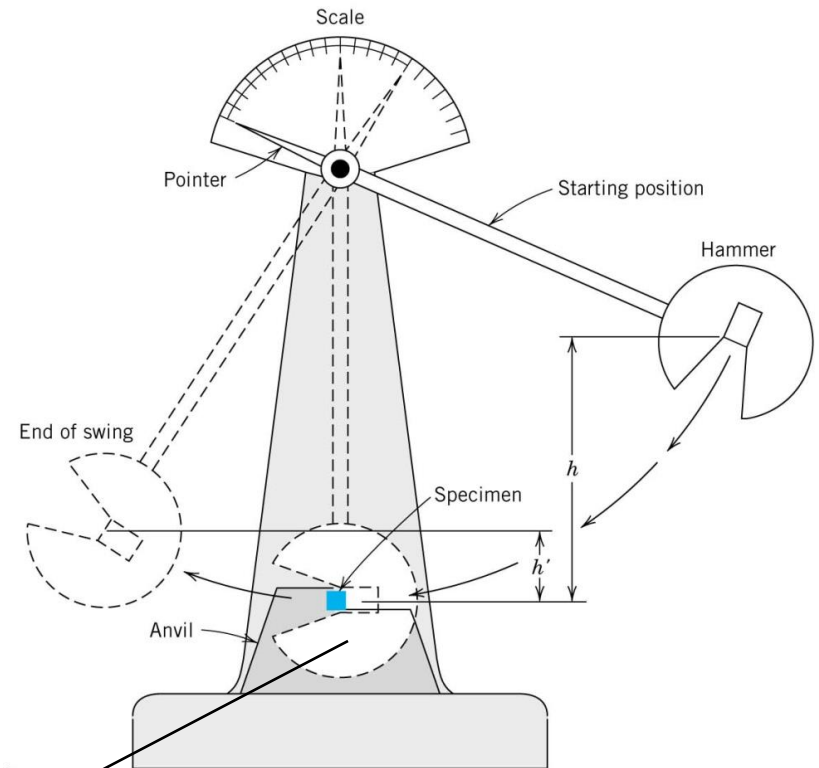
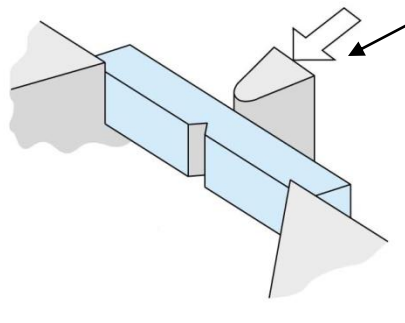
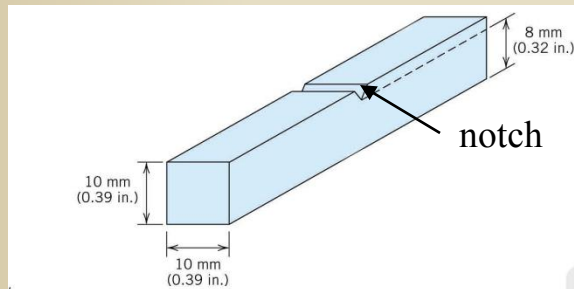


IMPACT FRACTURE TESTING

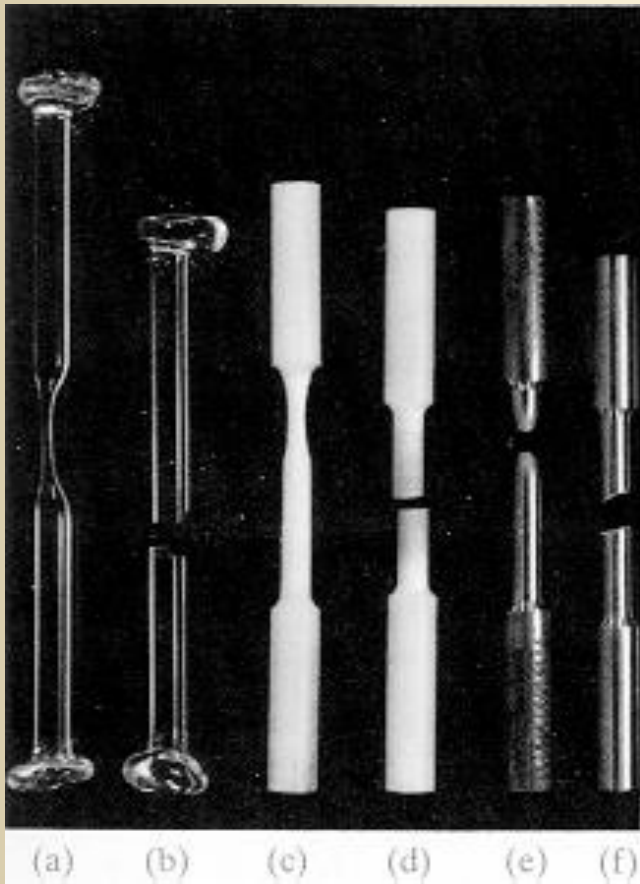
Charpy V-notch (CVN) technique

represents the most severe relative to the potential for fracture conditions:

- deformation at low temperature;
- a high strain rate
- a tri-axial stress state introduced by a notch and is used to measure the **impact energy** or notch toughness



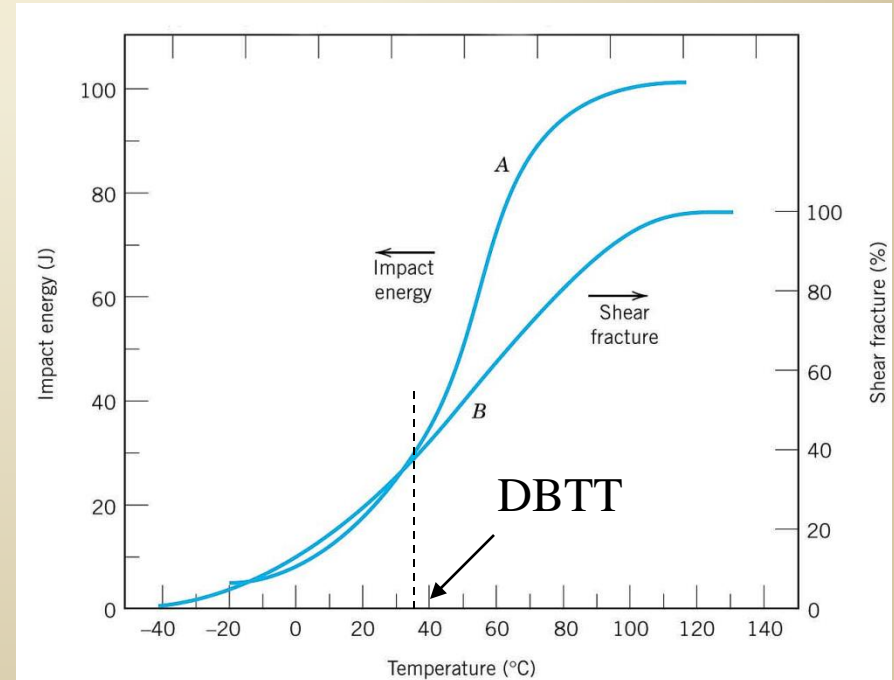
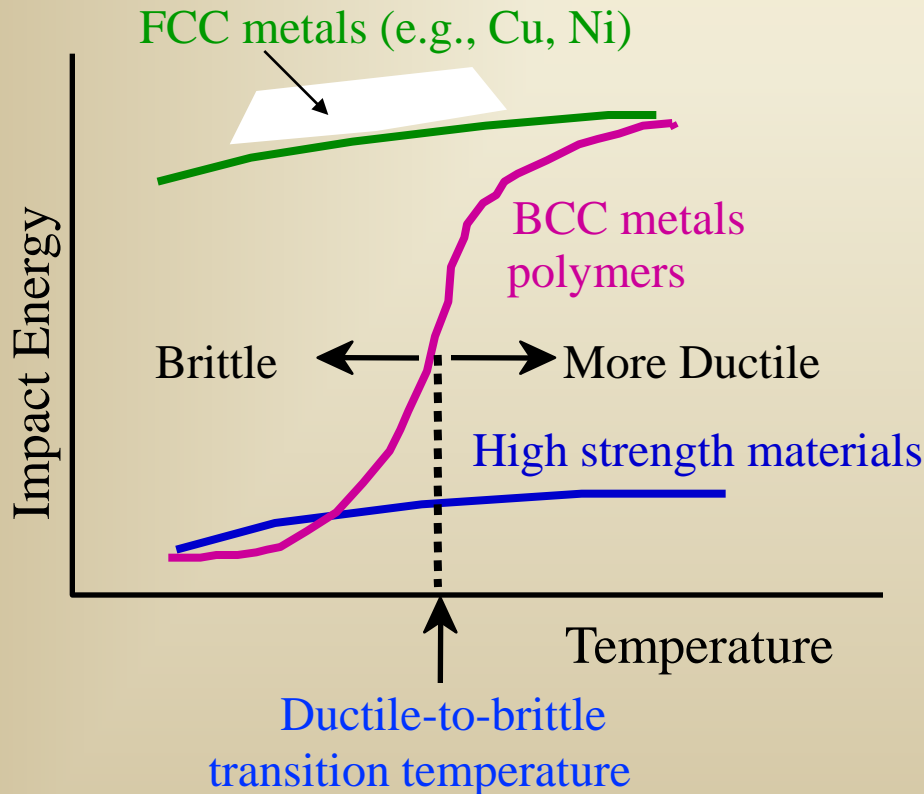
Temperature and Strain Rate



- Samples (a) and (b) are both glass, but (a) was tested at **800 K** and behaved in **a ductile** (viscoelastic) way whereas (b) was tested at **273 K** and showed **brittle** behavior.
- Poly-propylene samples (c) and (d) were tested at the same temperature but at **different strain rates**. At low strain rates (c) this polymer is **ductile** but at high strain rates (d) it shows **brittle** fracture.

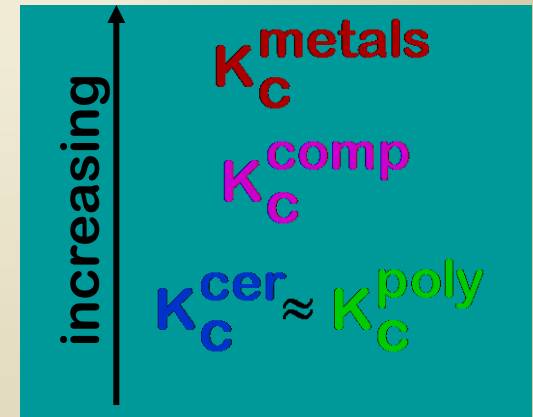
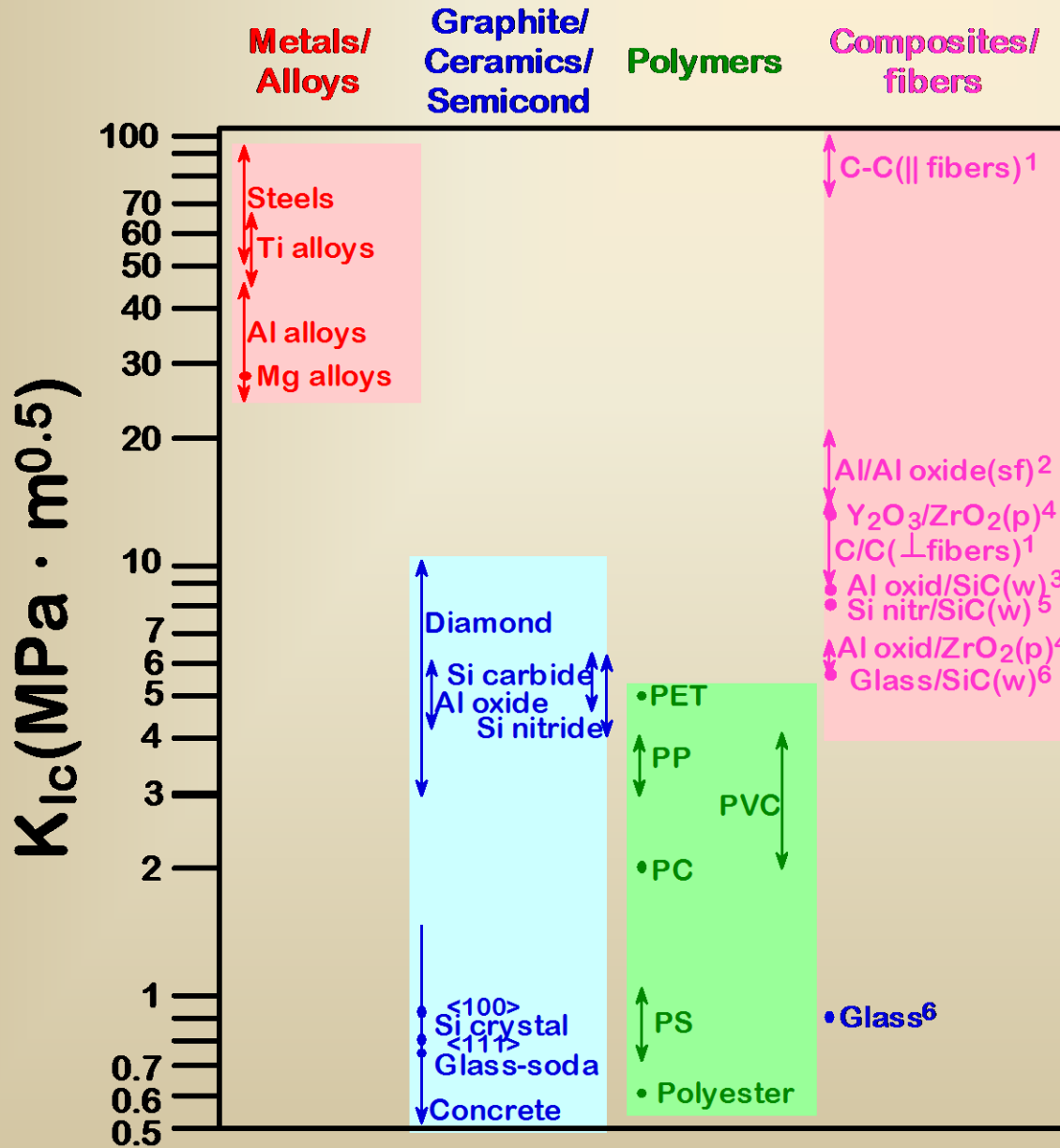
TEMPERATURE

- Increase of temperature leads to the increase of shear fracture and K_{Ic}
- **Ductile-to-brittle transition temperature (DBTT)**



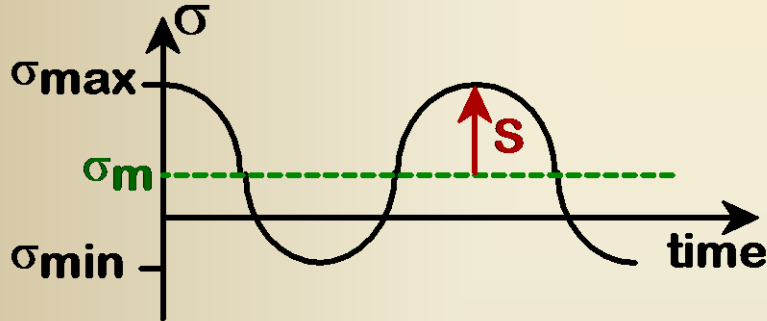
Example: A283 Steel

FRACTURE TOUGHNESS



FATIGUE

• **Fatigue** is a failure **under cyclic stress** (bridges, aircraft etc.).



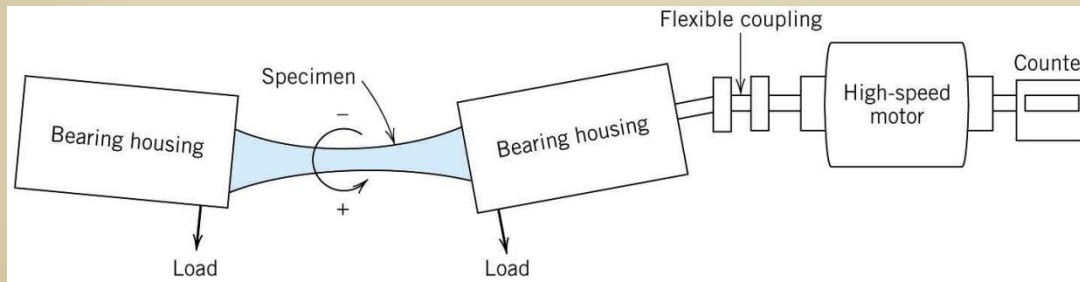
- Stress varies with time:
 - key parameters are S , σ_{mean} , $\Delta\sigma$

$$\sigma_m = \frac{\sigma_{\max} + \sigma_{\min}}{2}$$

$$\Delta\sigma_a = S = \frac{\sigma_{\max} - \sigma_{\min}}{2}$$

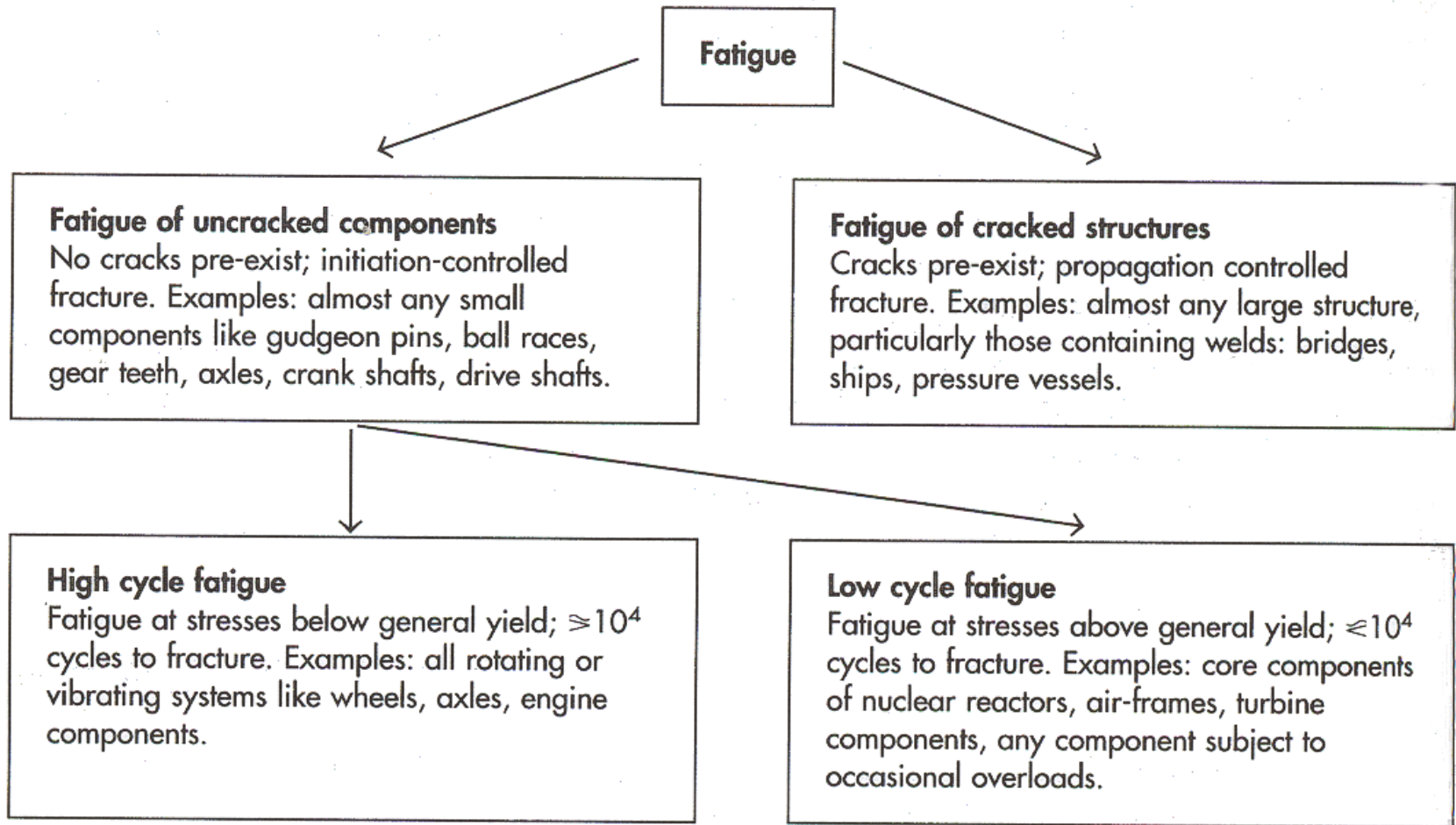
• **Key points:**

- Fatigue can cause part failure, even though $\sigma_{\max} < \sigma_c$, i.e. at lower strength than for a static conditions;
- Fatigue causes ~ 90% of mechanical engineering failures.

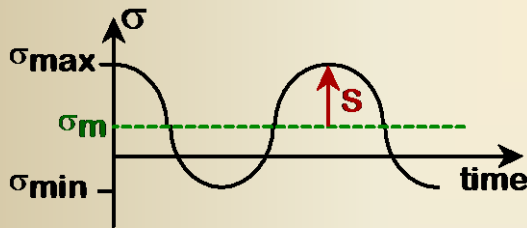


Fatigue-testing apparatus for rotating-bending test: main parameters S and number of cycles (N)

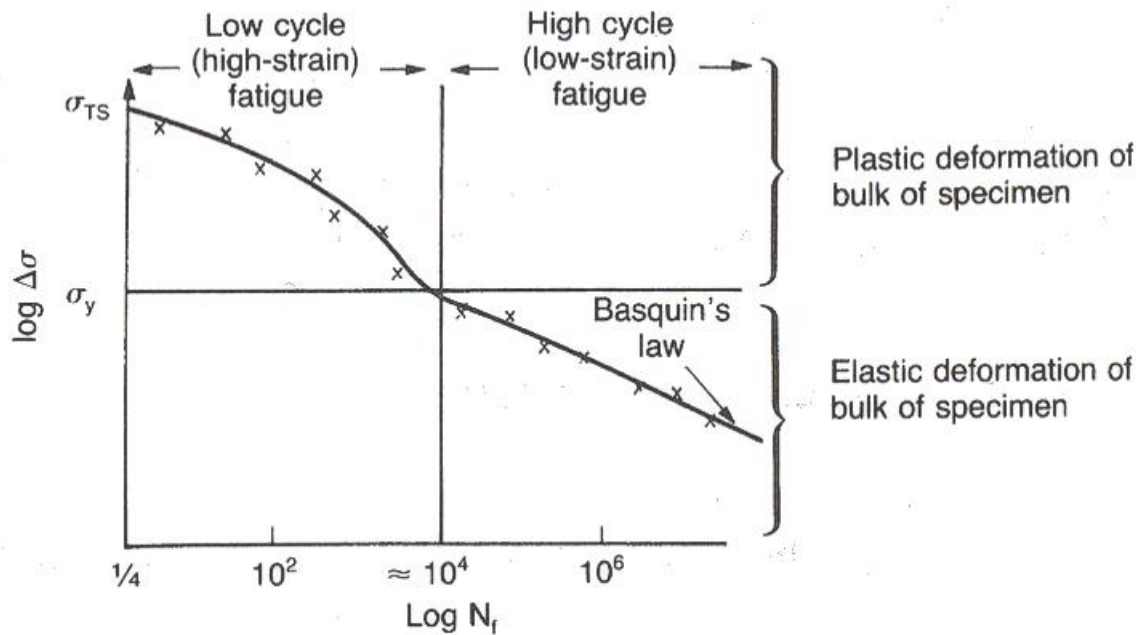
Fatigue Failure: Classification



Fatigue of Un-Cracked Component: high cycle fatigue



Both σ_{\max} and $|\sigma_{\min}|$ are **below yield stress**.

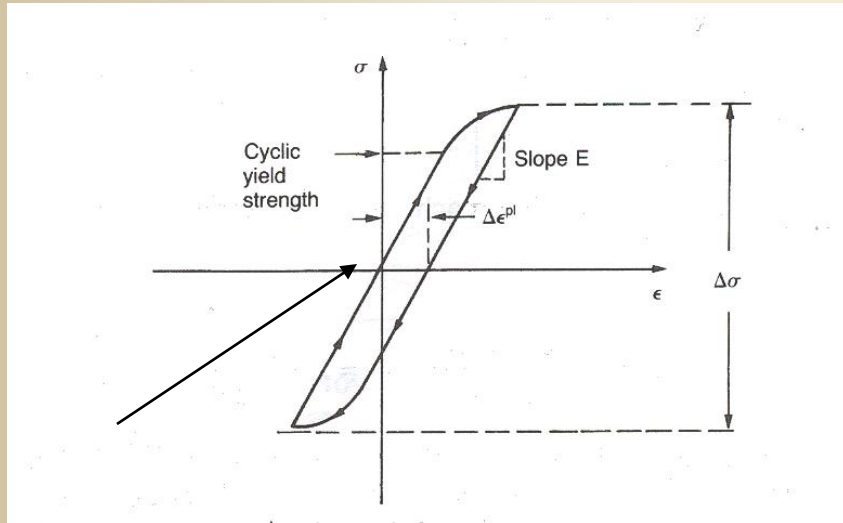


Empirical formula:

$$\Delta\sigma N^a = C_1$$
Basquin's law

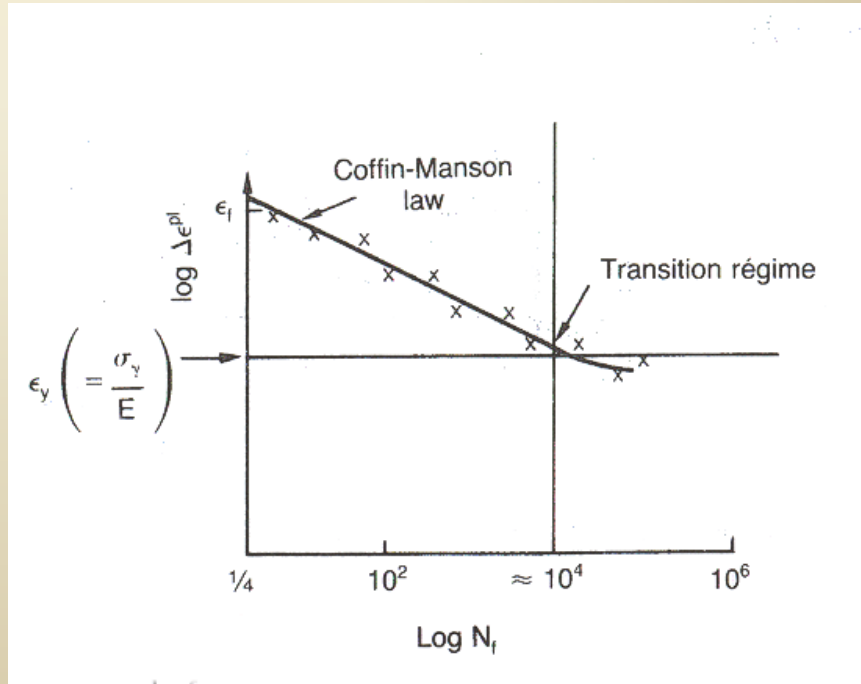
N - number of cycle
 C_1 - constant
 a - constant
 $\Delta\sigma = (\sigma_{\max} - \sigma_{\min}) / 2$

Fatigue of Un-Cracked Component: low-cycle fatigue



σ_{\max} and/or σ_{\min} are **above yield stress**.

Plastic strain range $\Delta\epsilon^{pl}$

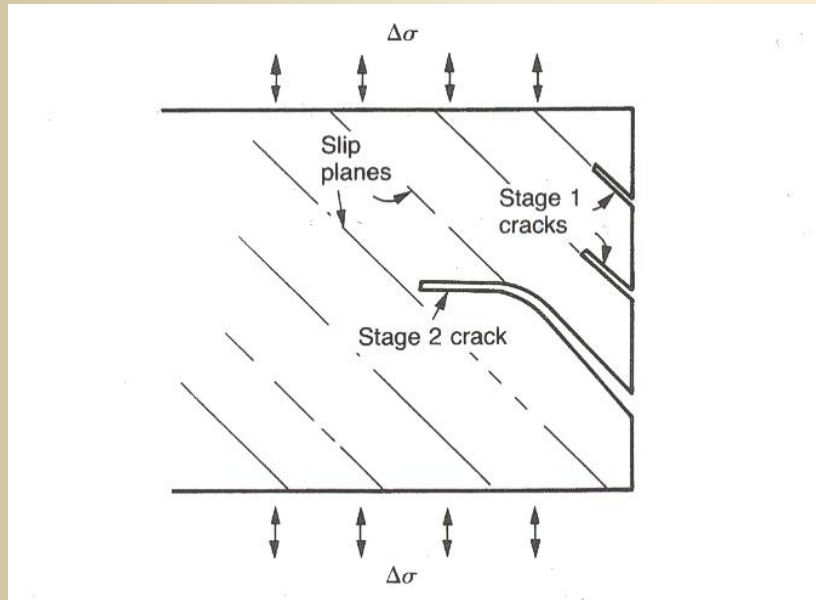


Empirical formula:

$$\Delta\epsilon^{pl} N^b = C_2$$
Coffin-Manson law

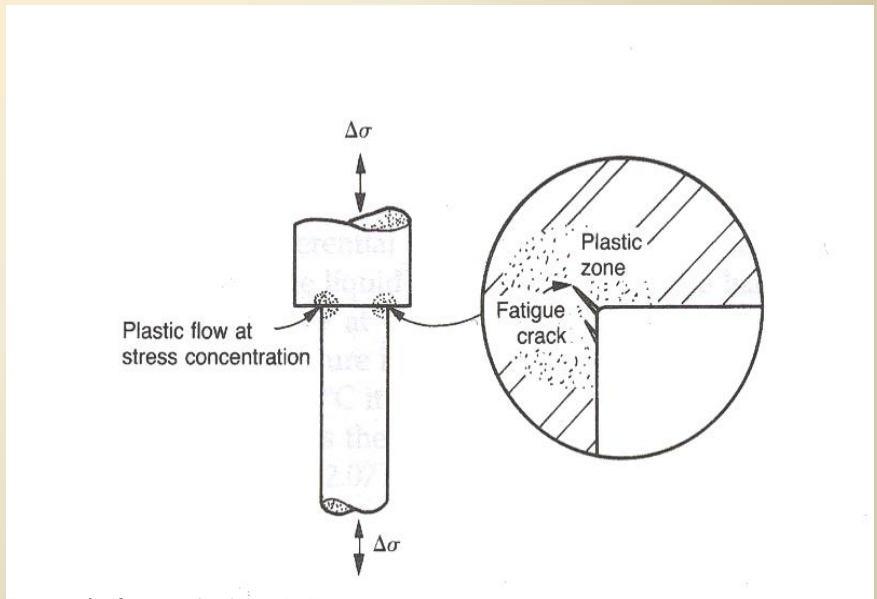
N - number of cycle
 C_2 – constant
 b - constant

Fatigue Mechanism: un-cracked structures



Low-cycle fatigue:

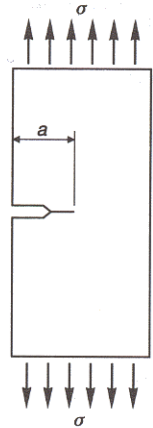
The general plasticity quickly roughens the surface and crack *forms there*, first propagating along a *slip plane* and then normal to the tensile axis



High-cycle fatigue:

Stress is below σ_y , thus essentially all of the life is time up in *initiating* a crack in the place of local plasticity, which is related to the *zones of stress concentrations*. Formed crack propagates slowly at first and then faster, until component fails.

Fatigue of Cracked Component

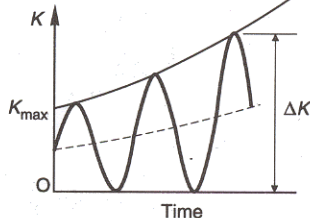
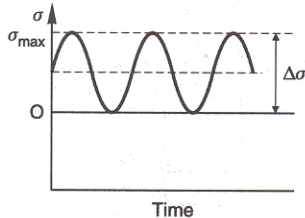
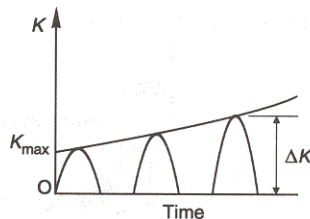
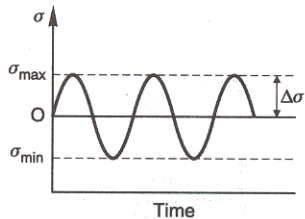


$$K = \sigma \sqrt{\pi a}$$

$$K_{max} = \sigma_{max} \sqrt{\pi a}$$

$$K_{min} = \sigma_{min} \sqrt{\pi a} \text{ for } \sigma_{min} > 0$$

$$K_{min} = 0 \text{ for } \sigma_{min} \leq 0$$



$$\Delta K = K_{max} - K_{min} = \Delta \sigma (\pi a)^{1/2}$$

The cycle stress intensity ΔK increases with time (at constant S) because the crack growth increases tension.

For example in steady-state regime the crack growth rate is:

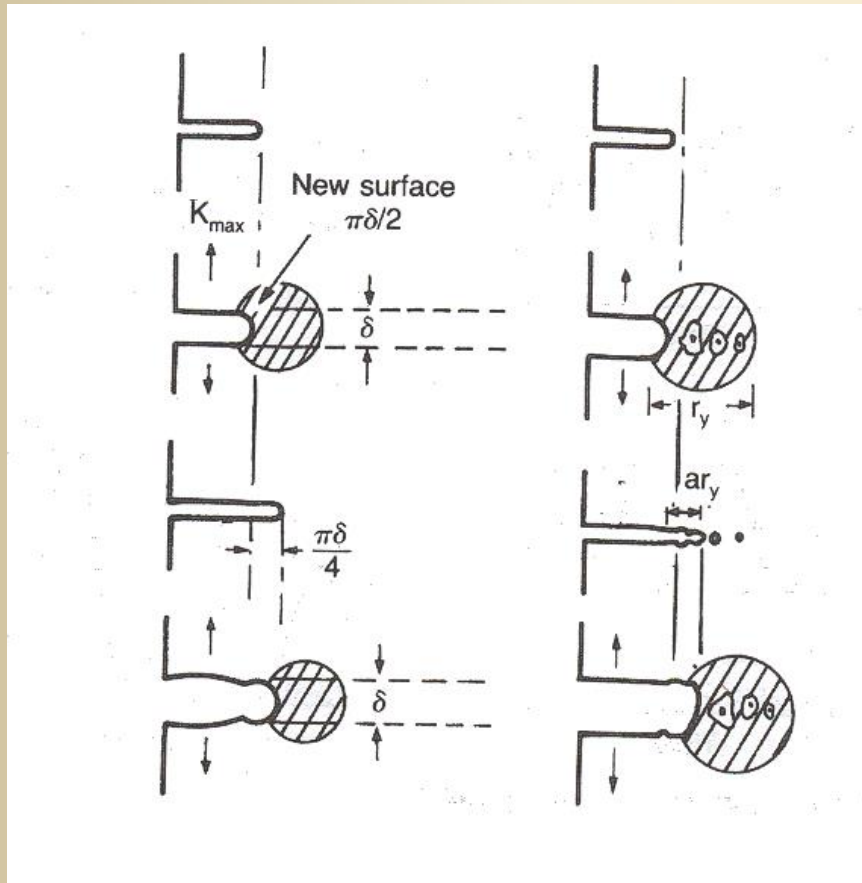
$$da/dN = A \Delta K^m$$

where A and m are materials' constant

$$N_f = \int_0^{N_f} dN = \int_{a_0}^{a_f} \frac{da}{A(\Delta K)^m} = \int_{a_0}^{a_f} \frac{da}{A[\Delta \sigma (\pi a)^{1/2}]^m}$$

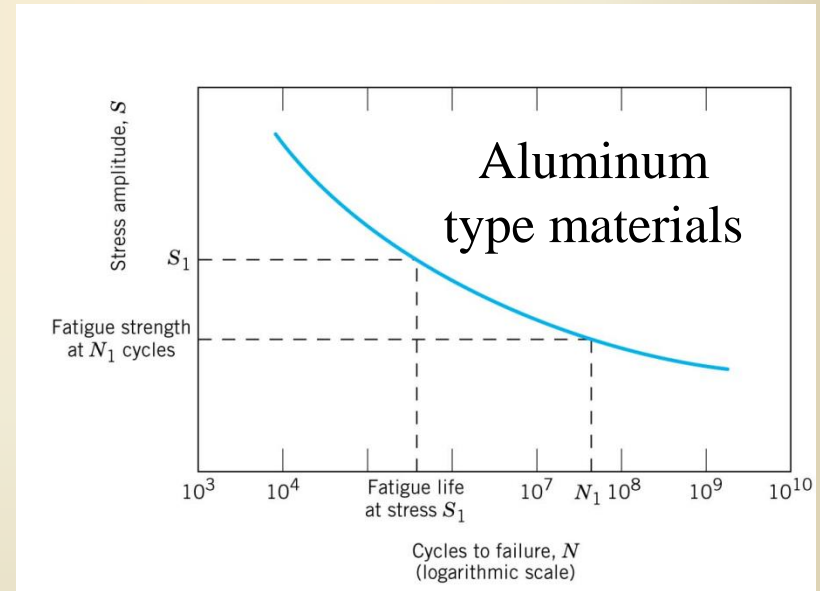
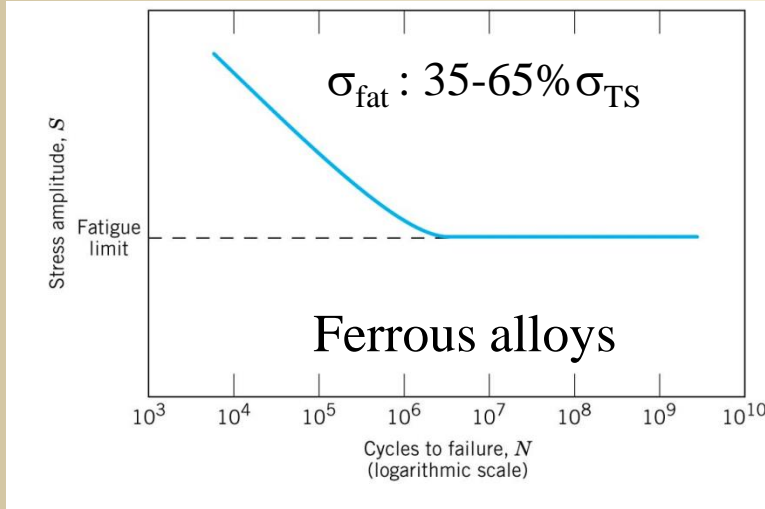
Large, particularly *welded structures* (bridge, ships, nuclear reactors)

Fatigue Mechanism: pre-cracked structures

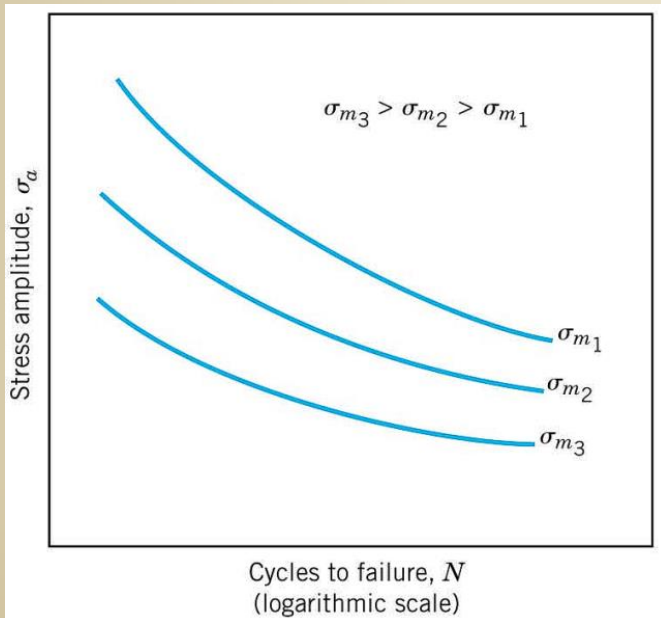


- In pure metals or polymers the tensile stress produces a plastic zone which makes the crack tip stretch open by the amount of δ , creating **a new surface** there.
- As the stress is removed the crack closes and new surface folds forward extending crack $\sim \delta$.
- On the next cycle the same happens again and **crack moves with rate**
 $da/dN \sim \delta$.
- **Inclusions** make the crack propagate even **faster**.

S-N DIAGRAMS: FATIGUE DESIGN PARAMETERS



- Fatigue limit, S_{fat} : -no fatigue if $S < S_{fat}$



- Sometimes, the fatigue limit is **zero!**

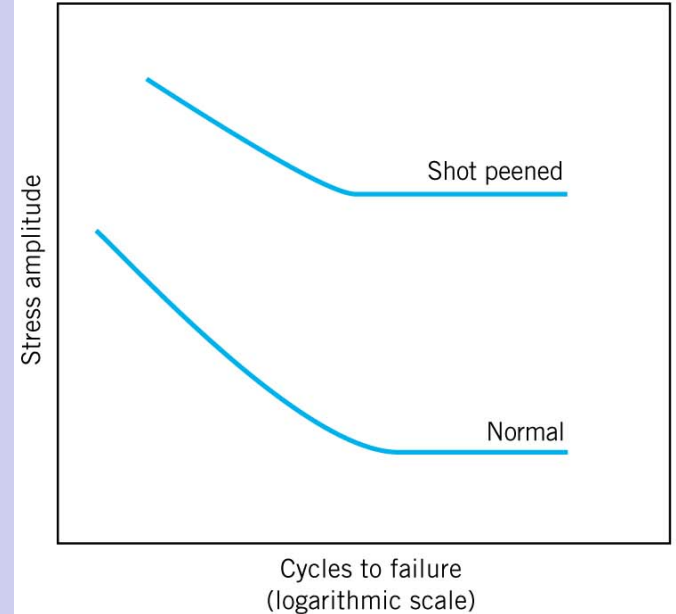
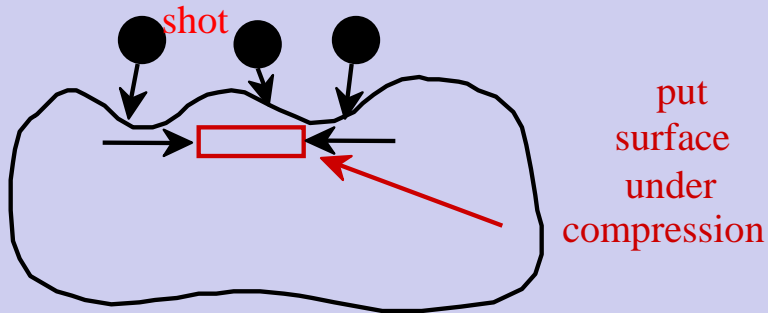
Factors affect fatigue life:

- increasing the **mean stress** leads to a decrease in fatigue life
- **surface effects:** design factor, surface treatment

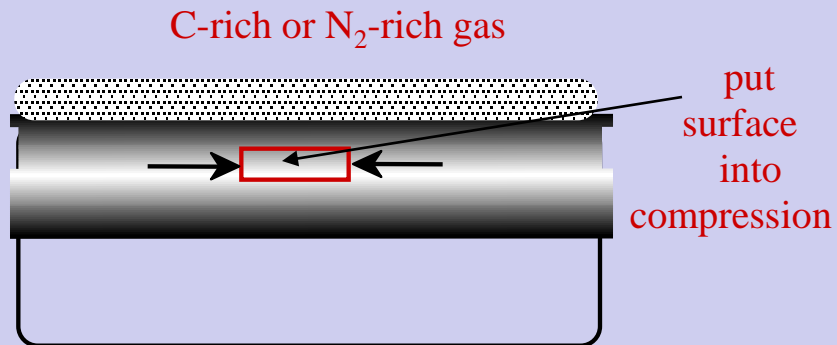
IMPROVING FATIGUE LIFE

1. Impose a **compressive surface stress**:
to suppress surface cracks from growing.

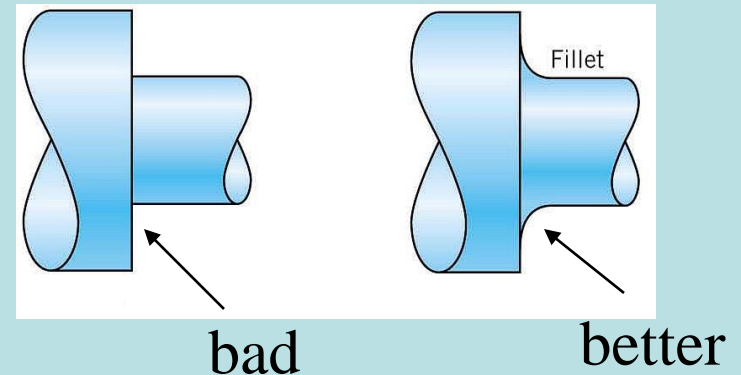
-Approach 1: **shot peening**



-Approach 2: **carburizing or nitriding**

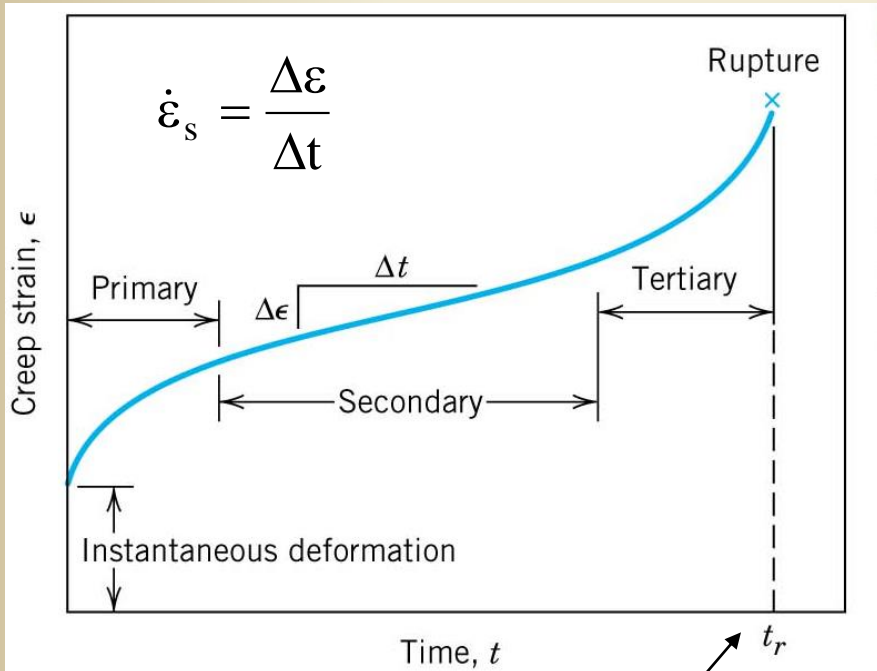


2. **Remove stress concentrators.**



3: CREEP

- **Creep** is deformation at elevated temperature ($T > 0.4 T_{\text{melt}}$) but under static mechanical stress

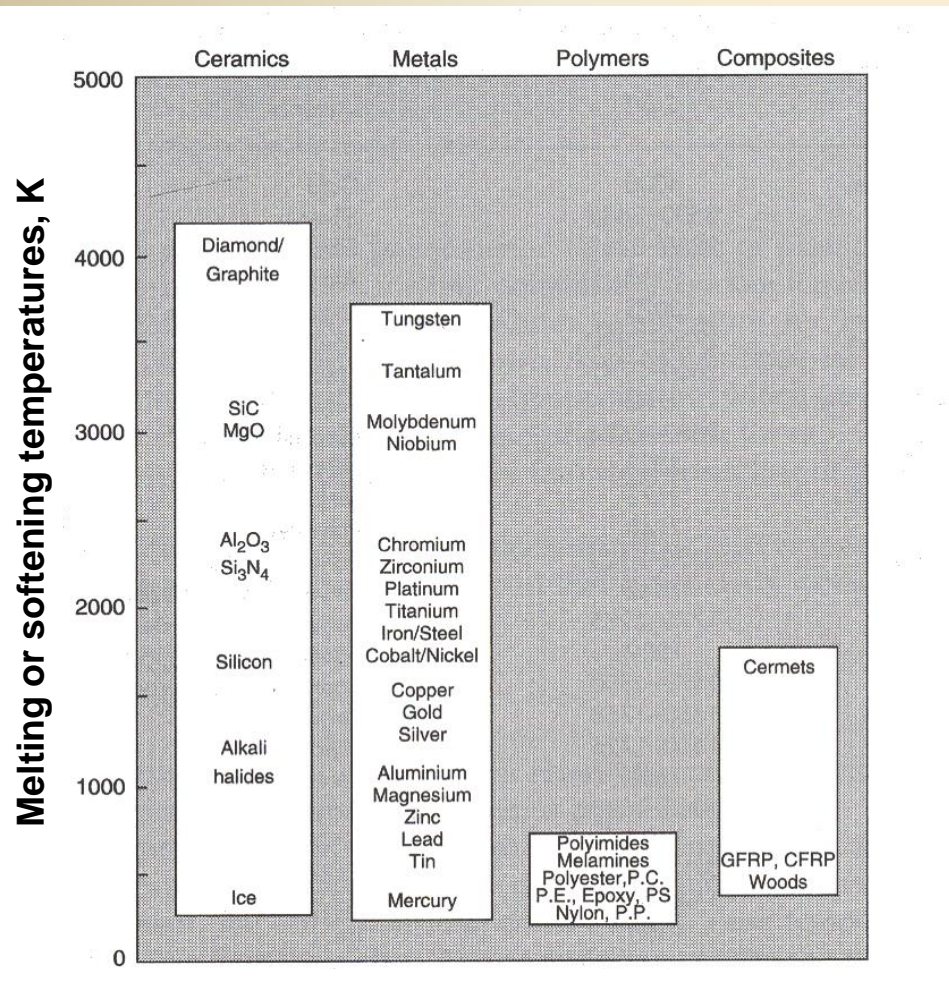


t_r , Creep rupture time

• But in this case deformation changes with time!!

- Three characteristic regions:
 - **primary creep**: creep rate decreases
 - **secondary creep**: steady-state creep,
 - **tertiary creep**: creep rate acceleration

What is low T and what is high T?



Examples:

Tungsten ~ 3000 °C

Lead ~ 400° C

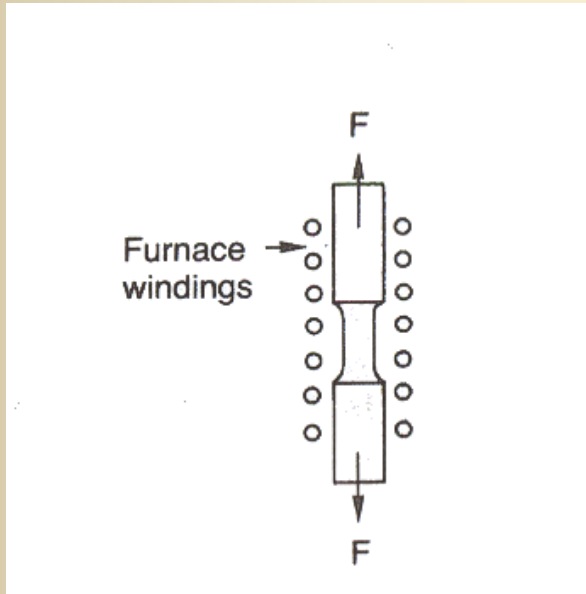
Ice ceramic 0 °C - !!

$T > 0.3$ to $0.4 T_M$ for **metals**

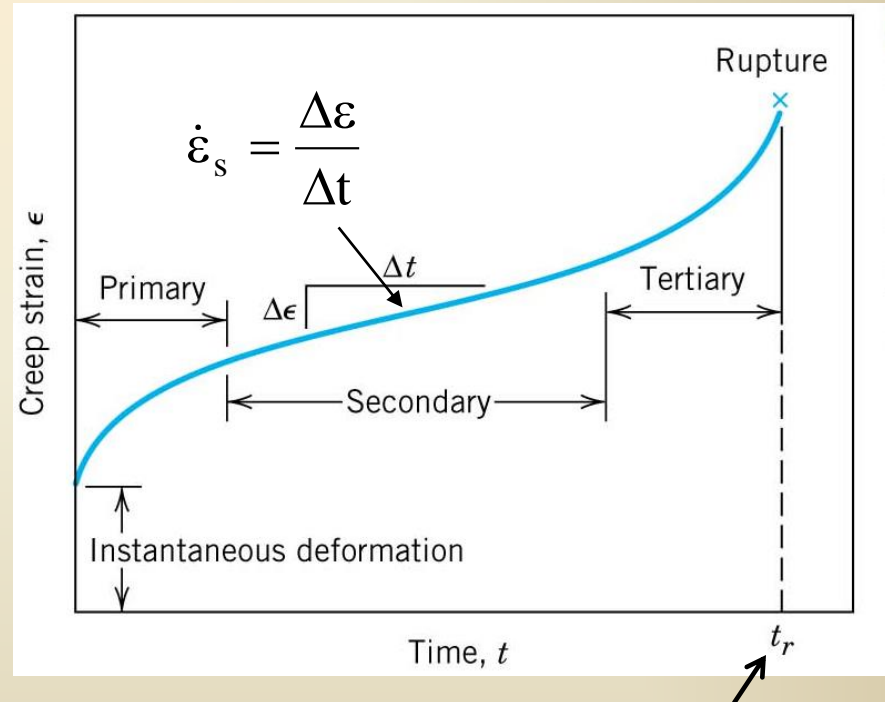
$T > 0.4$ to $0.5 T_M$ for **ceramics**

Polymers also creep – many of them do so at room temperature

Creep Curve



Schematics of Creep Curve Measurements



t_r , Creep rupture time

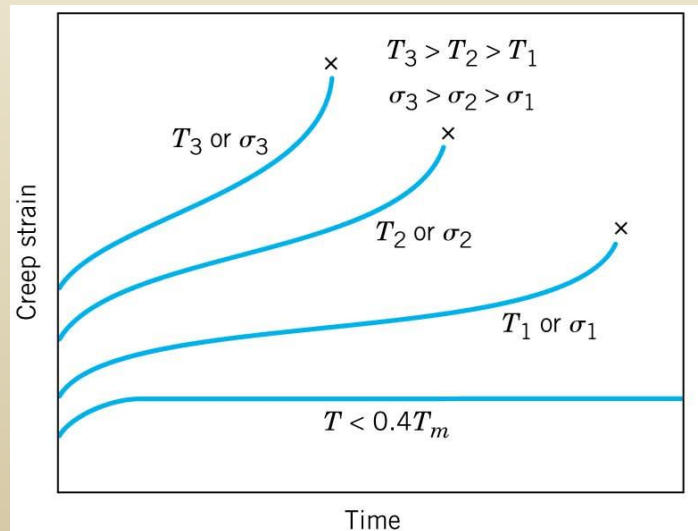
- Three characteristic regions:
 - **primary creep**: creep rate decreases
 - **secondary creep**: steady-state creep,
 - **tertiary creep**: creep rate acceleration

Steady-State (Secondary) Creep

- Typically the **longest duration** stage of creep
- **Strain rate is constant** at a given T and σ , because **strain hardening** is **balanced by recovery!!!**

$$\dot{\epsilon}_s = K_2 \sigma^n \exp\left(-\frac{Q_c}{RT}\right)$$

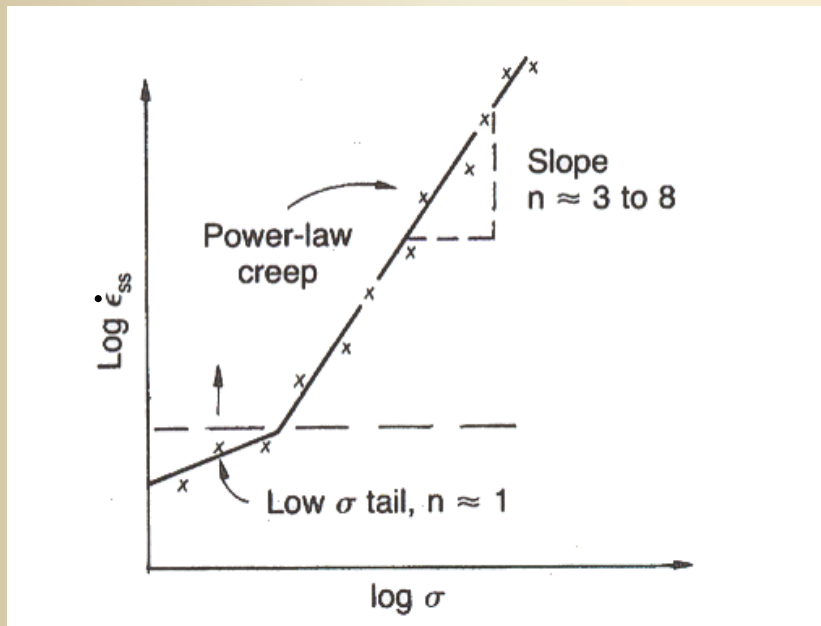
strain rate $\dot{\epsilon}_s$ material const. K_2 applied stress σ stress exponent (material parameter) n activation energy for creep (material parameter) Q_c RT



the strain rate increases for larger T, σ

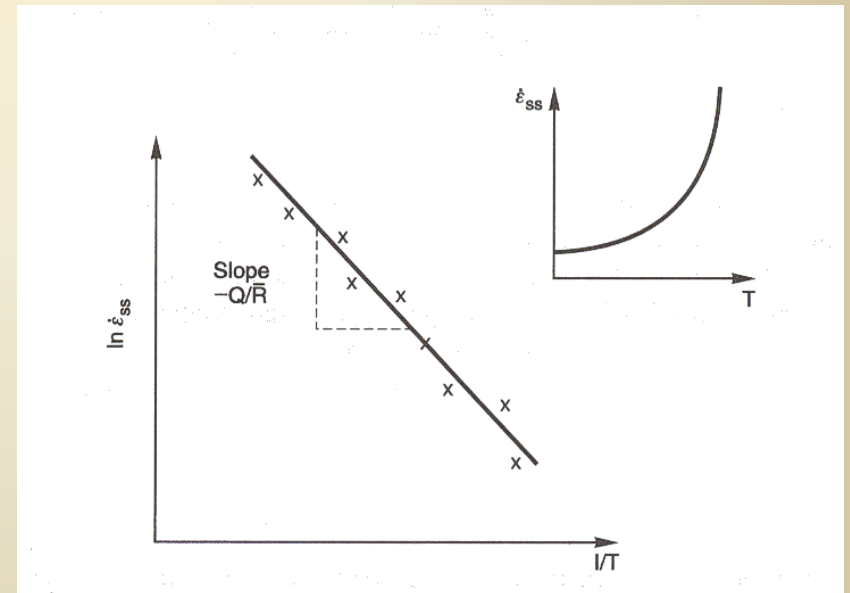
Steady-State Creep

Stress Dependence



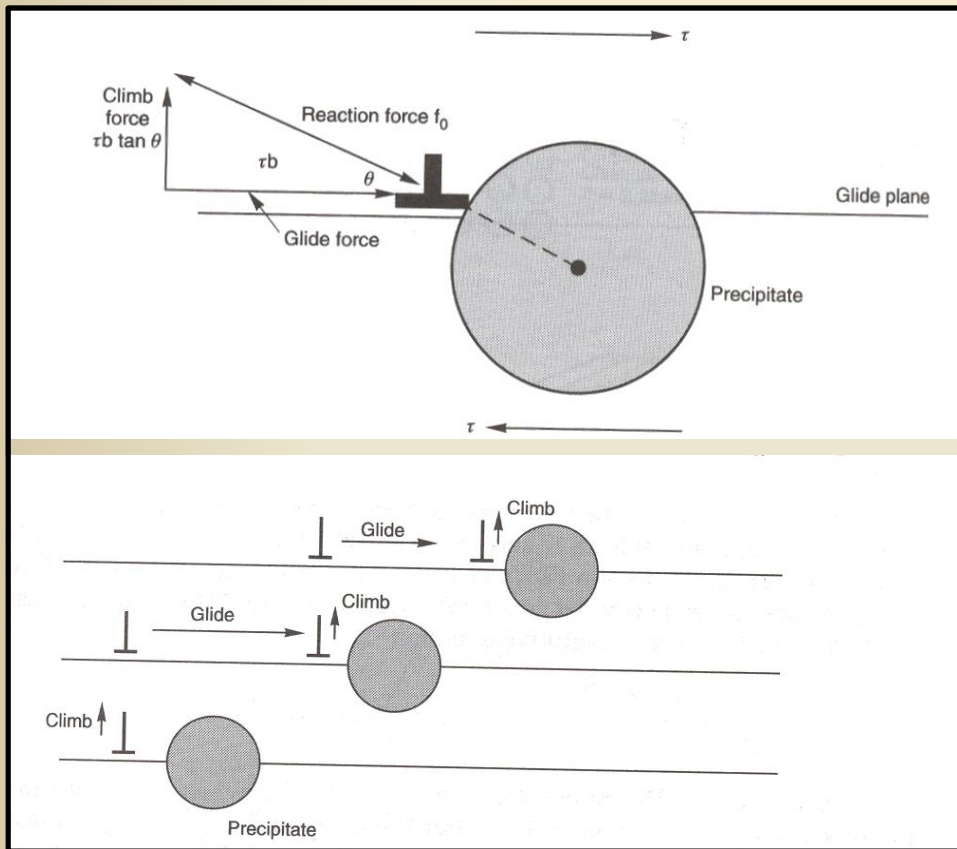
$$\dot{\epsilon}_{ss} = B\sigma^n$$

Temperature Dependence



$$\dot{\epsilon}_{ss} = Ce^{-(Q/RT)}$$

Creep Mechanisms: Dislocation Creep



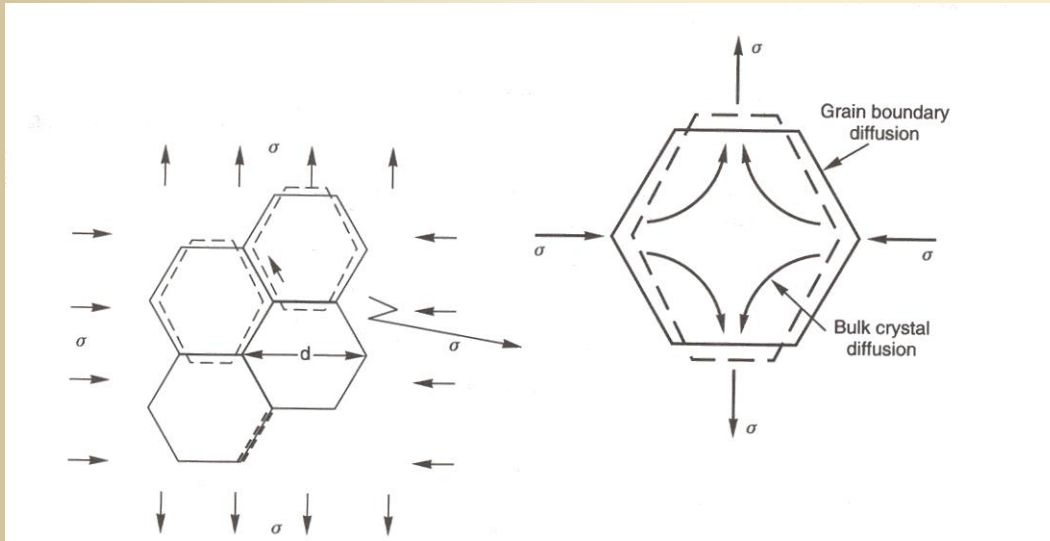
Climb unlocks dislocations from precipitate which pin them
And further slip may take place!!

$$D = D_0 \exp(-Q/RT)$$

$$\dot{\epsilon}_{ss} = A \sigma^n e^{-(Q/RT)}$$

- At lower end of the creep regime (0.3-0.5 T_m) **core diffusion** tends to dominate; at the higher end (0.5-0.9 T_m) it is **bulk diffusion**

Creep Mechanisms: Diffusion Creep



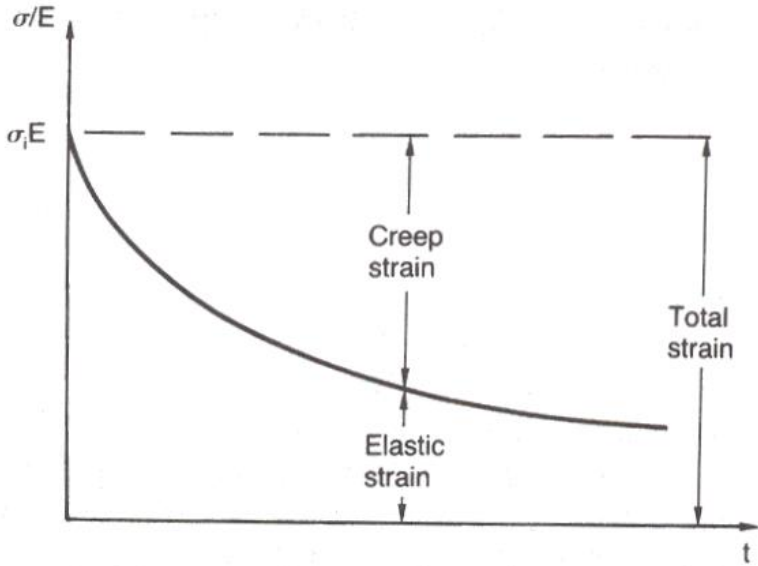
At **lower stresses** the rate of power-law creep falls quickly (from $n = 3 - 8$ to $n=1$). Creep does not stop, but rather proceeds by an **alternative mechanism**, i.e. a polycrystal can extend in response to applied stress by **grain elongation!**

In this case **atoms diffuse** from one set of the grain faces to the another and dislocations are not involved.

$$\dot{\epsilon}_{ss} = CD\sigma / d^2 = C'\sigma e^{-(Q/RT)} / d^2$$

grain size

Creep Relaxation



The relaxation time (arbitrarily defined as time taken for the stress to relax to half of original value) can be estimated as follows:

$$\varepsilon^{\text{total}} = \varepsilon^{\text{elastic}} + \varepsilon^{\text{creep}}$$

$$\varepsilon^{\text{el}} = \sigma / e$$

$$\dot{\varepsilon}^{\text{cr}} = B\sigma^n$$

$$\varepsilon^{\text{total}} = \text{const}$$

$$\frac{1}{e} \cdot \frac{d\sigma}{dt} = -B\sigma^n$$

Integrating from $\sigma = \sigma_{\text{in}}$ to $\sigma = \sigma_{\text{in}} / 2$

$$t_r = \frac{(2^{n-1} - 1)}{(n - 1)BE\sigma_{\text{in}}^{n-1}}$$

At constant displacement, creep causes stresses to relax with time!

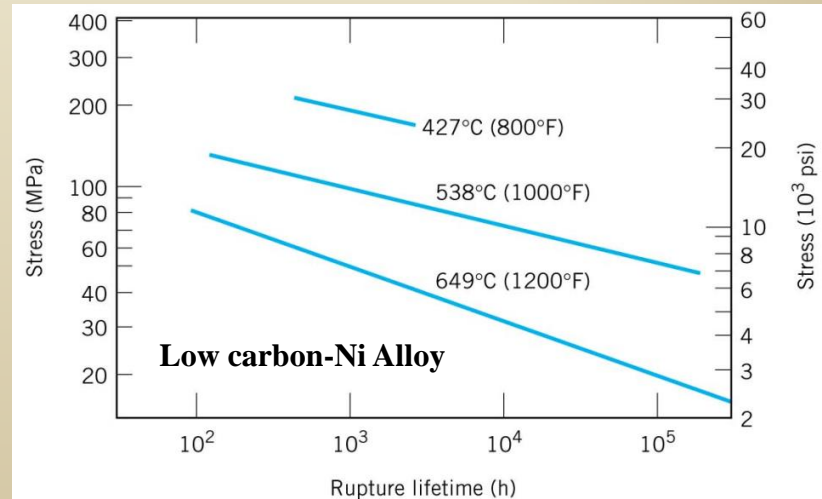
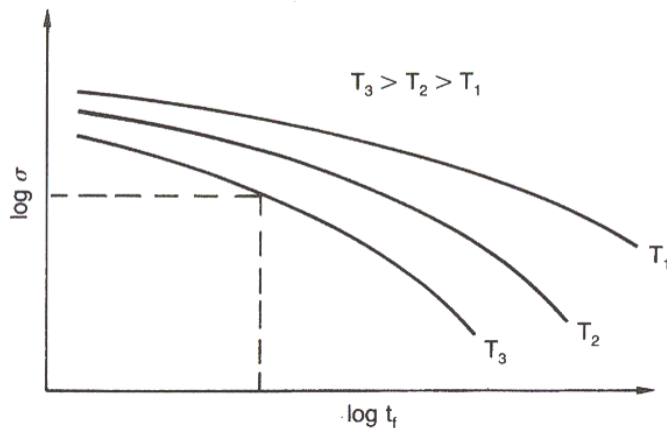
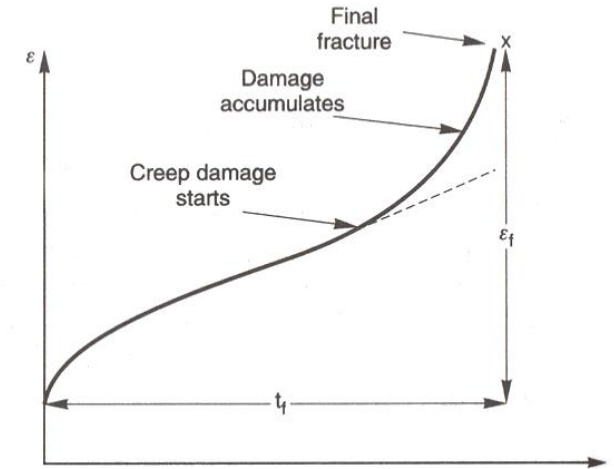
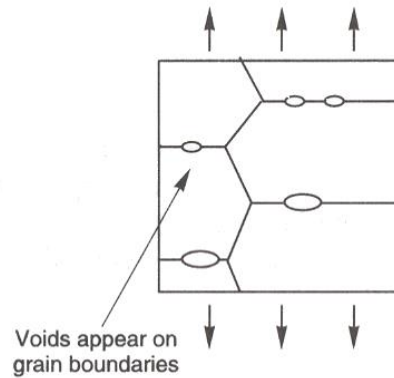
Example: Bolts in hot turbine must be regularly tightened; plastic paper-clips are not, in the long term, as good as a steel ones.

Creep Damage and Creep Fracture

Damage in the form of internal cavities, accumulates during Tertiary Stage of creep

Time-to-failure is described by equation which looks like that for creep itself:

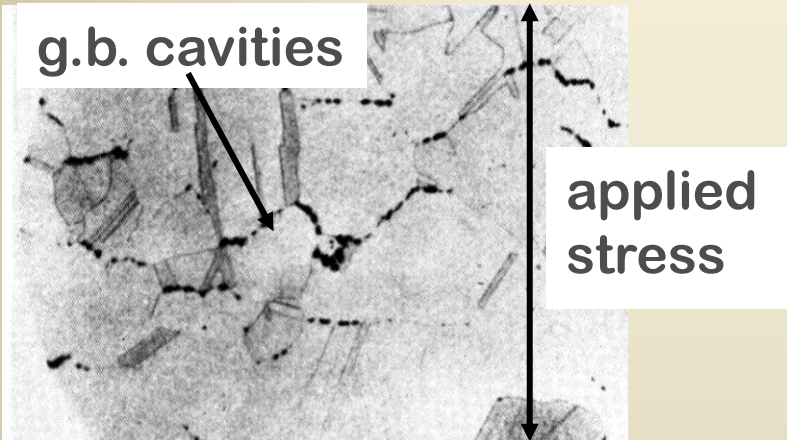
$$t_f = A' \sigma^{-m} e^{Q/RT}$$



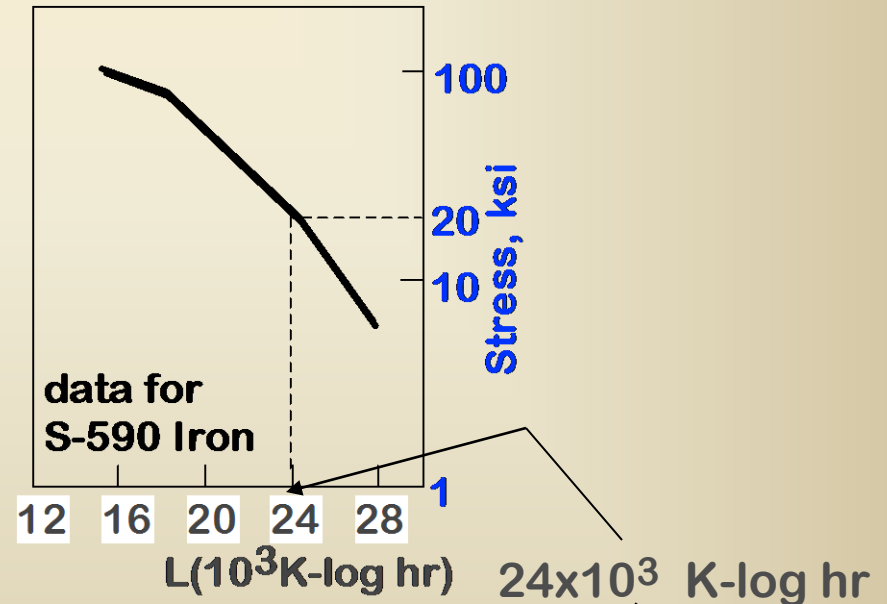
- the rupture time diminishes for larger T , σ

CREEP FAILURE

- Failure: along grain boundaries.



- Estimate rupture time
S-590 Iron, $T = 800^{\circ}\text{C}$, $\sigma = 20$ ksi



- Time to rupture, t_r

$$T(20 + \log t_r) = L$$

temperature

time to failure (rupture)

function of applied stress

$$T(20 + \log t_r) = L$$

1073K

Answer: $t_r = 233\text{hr}$

SUMMARY

- Engineering materials don't reach **theoretical strength**.
- **Flaws** produce **stress concentrations** that cause premature failure.
- Sharp corners also produce large stress concentrations and premature failure.
- **Failure type depends on stress and T:**
 - for non-cyclic σ and $T < 0.4T_m$, failure stress decreases with:
increased maximum flaw size, decreased T,
 - for cyclic:
cycles to fail decreases as $\Delta\sigma$ increases.
 - for higher T ($T > 0.4T_m$):
time to fail decreases as σ or T increases.