

Mean arterial pressure nonlinearity in an elastic circulatory system subjected to different hematocrits

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Abstract The level of hematocrit (Hct) is known to affect mean arterial pressure (MAP) by influencing blood viscosity. In the healthy population, an increase in Hct (and corresponding increase in viscosity) tends to raise MAP. However, data from a clinical study of type 2 diabetic patients indicate that this relationship is not universal. Instead, individuals in the lower levels of Hct range display a decrease in MAP for a given rise in Hct. After reaching a minimum, this trend is reversed, so that further increases in Hct lead to increases in MAP. We hypothesize that this anomalous behavior occurs due to changes in the circulatory autoregulation mechanism. To substantiate this hypothesis, we develop a physically based mathematical model that incorporates autoregulation mechanisms. Our model replicates the anomalous U-shaped relationship between MAP and Hct found in diabetic patients in the same range of Hct variability.

Keywords Blood pressure · Hematocrit · Blood viscosity · Model · Diabetes · Blood vessels · Elastic properties · Hypertension · Hypotension · Autoregulation

Nomenclature

P	Pressure	17
P_z	Pressure gradient	18
T	Tension	19
r	Vessel radius	
A	Activation	
C_j	Constant associated with the subscript j	
r_0	Reference vessel radius	21
S_t	Stimulus	
σ	Shear stress	
Q	Flow rate	
μ	Effective blood viscosity	
MAP	Mean arterial pressure	
Hct	Hematocrit	

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1 Introduction

Mean arterial pressure (MAP) depends on physical factors including blood volume, elastic properties of the circulatory tree, length and radius of blood vessels, and blood viscosity, a direct function of hematocrit (Hct). In this work, we focus exclusively on the influence of hematocrit. According to Poiseuille's law, vascular resistance is predominantly determined by vessel diameter, which in healthy individuals is controlled by vascular smooth muscle tone. The latter, in turn, depends on myogenic and chemical autoregulation, influenced by extrinsic neurological and hormonal controls (Levick 2003). Under physiological conditions, the variability of blood viscosity is not sufficient to produce anomalous changes in MAP (Salazar-Vázquez et al. 2006).

36 However, [Martini et al. \(2005\)](#) found that MAP decreased
37 significantly when Hct was raised acutely in healthy ham-
38 sters, leading to a U-shaped MAP versus Hct distribution—a
39 finding attributed to the competition between nitric oxide
40 (NO) production by shear stress due to changes in blood
41 viscosity and changes in resistance due to changes in blood
42 viscosity. [Salazar-Vázquez et al. \(2006\)](#) found a similar rela-
43 tionship in type 2 diabetic patients, when MAP was plot-
44 ted as a function of the naturally occurring variability of
45 Hct in the sampled population. Both studies show that there
46 is a range of Hct where increases of Hct cause MAP to
47 decrease, suggesting the presence of blood-rheology sensing
48 mechanisms.

49 Autoregulation of vascular tone is either attenuated or
50 absent under conditions of endothelial dysfunction, e.g., in
51 the case of endothelial and intimal damage characteristic of
52 diabetes. Diabetic angiopathy lowers myogenic autoregula-
53 tion through increased collagen cross-linkage by increasing
54 vascular stiffness ([Cameron and Cruickshank 2007](#)). The lat-
55 ter results in secondary hypertension and a MAP that is raised
56 sufficiently to maximize the myogenic response ([Carlson et
57 al. 2008](#)). The shear stress response is impeded by the depo-
58 sition of advanced glycation end products on intact extra-
59 cellular matrix which can absorb, hinder diffusion of, and
60 inactivate NO ([Shore and Tooke 1997](#)). Glycated hemoglo-
61 bin, free radicals, and oxidized low-density lipoproteins also
62 quench the effect of NO. In diabetics, autonomic autoreg-
63 ulation is impaired as shown by the deficient endothelium-
64 dependent relaxation to acetylcholine and bradykinin ([Malik
65 et al. 2005](#)).

66 Under these conditions, microcirculatory regulation
67 resembles that of a system of rigid tubes. Blood viscosity
68 affects peripheral resistance and hence MAP. Maintaining
69 flow rates requires varying driving pressure in proportion to
70 changes in viscosity. The behavior of pressure in the circula-
71 tion also approximates that within a system of rigid tubes for
72 extreme changes of blood viscosity, such as pathologically
73 high Hcts ([Bertinieri et al. 1998](#); [Jefferson et al. 2002](#)) and
74 extreme Hct reductions ([Toy et al. 2000](#)).

75 In an elastic tube, MAP and blood viscosity are related
76 in part by the material properties of the tube (vessel) wall,
77 the viscosity of the flowing fluid, and the tube's autoreg-
78 ulatory capacity. The relationship between changes in Hct
79 and average pressure in an elastic tube is explored by
80 analyzing the interactions between vessel wall mechanics
81 and the pressure distribution in blood vessels at different
82 viscosities.

83 We use mathematical modeling to investigate the relation-
84 ship between pressure and the viscosity of a fluid flowing in
85 an autoregulating elastic tube. The model employs param-
86 eters and boundary conditions related to the arteriolar cir-
87 culation and sheds new light on the anomalous U-shaped
88 MAP–Hct distribution found in diabetic patients.

2 Mathematical model

89 Standard analytical models of blood flow in vessels account
90 for their elasticity by invoking the concept of compliance
91 ([Loscalzo and Schafer 2002](#); [Secomb 2008](#)). Such mod-
92 els typically result in monotonically increasing relations
93 between pressure and Hct. The proposed *autoregulatory elas-
94 tic compliance model* illustrates that under certain conditions,
95 the MAP–Hct curve can exhibit a U-shape observed experi-
96 mentally by [Salazar-Vázquez et al. \(2006\)](#).

97 Our model accounts for the possibility of such a behav-
98 ior by invoking two postulates. First, the degree to which a
99 blood vessel dilates or constricts in response to changes in
100 blood pressure is determined, at least in part, by its elastic
101 properties. Second, pressure at the tube inlet varies with the
102 vessel inner radius in an autoregulatory fashion. The pro-
103 posed autoregulatory elastic compliance model allows for an
104 arbitrary dependence of blood viscosity on Hct and is valid
105 for either linear or pseudo-linear descriptions of the blood
106 vessel, with both descriptions yielding analogous MAP ver-
107 sus Hct dependence.

108 We develop a mathematical model that describes how
109 blood pressure P changes with hematocrit (Hct), in a manner
110 that accounts for the elasticity of blood vessels and incorpo-
111 rates various autoregulation mechanisms. We assume steady-
112 state blood flow with constant volumetric flow rate Q in a
113 single blood vessel. This approximation is justified by the
114 very definition of autoregulation which is “the maintenance
115 of a constant supply of blood to an organ in spite of varying
116 arterial pressure”¹. Furthermore, we focus on blood vessels
117 whose effective inner radii are sufficiently large for the Stokes
118 equation to provide an adequate description of blood flow.
119 This is a standard approach that can be found in, e.g., Chapter
120 17 in [Loscalzo and Schafer \(2002\)](#) and Section 5.1 in [Keener
121 and Sneyd \(2004\)](#).

122 Under these assumptions, the volumetric flow rate is esti-
123 mated by Poiseuille's law,
124

$$125 \quad Q = -\frac{\pi r^4 P_z}{8\mu}, \quad (1)$$

126 where r is the radius of a blood vessel, P_z is the axial pressure
127 gradient along the vessel, and μ is the effective blood viscos-
128 ity. The radius of an elastic blood vessel, r , is determined by
129 a balance between blood pressure P and tension in the vessel
130 wall, T from Laplace's law for thin-walled vessels ([Carlson
131 et al. 2008](#); [Carlson and Secomb 2005](#); [Secomb 2008](#)),

$$132 \quad P = \frac{T}{r}. \quad (2)$$

133 The tension in the vessel wall, T , consists of passive (T_{passive})
134 and active (T_{active}) components and can be represented as

¹ <http://www.merriam-webster.com/medlineplus/autoregulation>.

135 their weighted sum,

$$136 \quad T = T_{\text{passive}} + AT_{\text{active}}, \quad (3)$$

137 where A is the fraction of the total autoregulatory potential
138 being exerted by the vessel (i.e., a number between 0 and
139 1). The passive tension is a function of the vessel's mechan-
140 ical properties, and the active tension reflects the vessel's
141 autoregulatory capacity.

142 Following Carlson et al. (2008), Carlson and Secomb
143 (2005), Secomb (2008), we adopt the following phenome-
144 nological relations for T_{passive} , T_{active} , and A . The passive
145 component of tension T_{passive} satisfies an exponential model,

$$146 \quad T_{\text{passive}} = C_{\text{passive}} \exp \left[C'_{\text{passive}} (r/r_0 - 1) \right], \quad (4)$$

147 which, similar to the often-used linear compliance model,
148 reflects a monotonic relation between passive tension and
149 radius, i.e. any increase in tension causes an increase in
150 radius. The active component of tension T_{active} is described
151 with a Gaussian model,

$$152 \quad T_{\text{active}} = C_{\text{active}} \exp \left[- \left(\frac{r/r_0 - C'_{\text{active}}}{C''_{\text{active}}} \right)^2 \right]. \quad (5)$$

153 The activation A is related to the intensity of the autoregula-
154 tory stimulus S_t via a sigmoidal function

$$155 \quad A = \frac{1}{1 + \exp(-S_t)}. \quad (6)$$

156 According to Levick (2003), the intensity of the
157 autoregulatory stimulus S_t depends primarily on the fol-
158 lowing three phenomena (although additional ones can eas-
159 ily be included). *Myogenic response*, $S_{\text{myog}} = C_{\text{myog}}T$,
160 accounts for a vessel's response to changes in intraluminal
161 pressure P (or in accordance with Eq. 2 to changes in
162 tension T), whereby increased pressure causes smooth mus-
163 cle stretching. Myogenic autoregulation depends on stretch
164 activated ion channels in vascular smooth muscle that, upon
165 stretching, allow calcium ions to enter, inducing contraction.
166 *Shear-dependent response*, $S_{\text{shear}} = -C_{\text{shear}}\tau$, represents
167 changes in endothelial NO release (a potent vasodilator)
168 caused by changes in endothelial shear stress τ . *Metabolic*
169 *response* S_{meta} accounts for the fact that tissue hypoxia stim-
170 ulates formation of vasodilator metabolites, such as NO. This
171 metabolic response of the vasculature to alterations in oxy-
172 gen demand is proportional to the conducted response signal
173 S_{CR} , so that $S_{\text{meta}} = -C_{\text{meta}}S_{\text{CR}}$. Combining these three
174 response mechanisms, $S_t = S_{\text{myog}} + S_{\text{shear}} + S_{\text{meta}} + C_{\text{tone}}$,
175 where C_{tone} is a dimensionless constant, Carlson et al. (2008)
176 obtained a phenomenological model for the intensity of the
177 autoregulatory stimulus,

$$178 \quad S_t = C_{\text{myog}}T - C_{\text{shear}}\tau - C_{\text{meta}}S_{\text{CR}} + C_{\text{tone}}. \quad (7)$$

179 The constants C_{passive} , C'_{passive} , C_{active} , C'_{active} , C''_{active} ,
180 C_{myog} , C_{shear} , C_{meta} , and C_{tone} in Eqs. (4, 5, 6, 7) are fitting
181 parameters whose values depend on a vessel's radius. Below
182 we will conduct a sensitivity analysis to quantify the effects
183 of uncertainty in these values on modeling predictions.

184 Combining Eqs. (2, 3, 4, 5, 6, 7) yields a relationship
185 between intraluminal blood pressure P and vessel radius r ,

$$186 \quad P = \frac{C_{\text{passive}}}{r} \exp \left[C'_{\text{passive}} (r/r_0 - 1) \right] \\ 187 \quad + \frac{C_{\text{active}}}{r} \frac{1}{1 + \exp(-C_{\text{myog}}T + C_{\text{shear}}\tau + C_{\text{meta}}S_{\text{CR}} - C_{\text{tone}})} \\ 188 \quad \times \exp \left[- \left(\frac{r/r_0 - C'_{\text{active}}}{C''_{\text{active}}} \right)^2 \right], \quad (8)$$

189 which accounts for both the natural elastic properties of blood
190 vessels and autoregulation. To relate intraluminal blood pres-
191 sure P (and MAP) to blood viscosity μ (and hematocrit
192 (Hct)), we supplement the static relation (8) with the dynamic
193 equation (1). This step requires both a closure assumption to
194 relate pressure gradient P_z to pressure P and a constitutive
195 law, $\mu = \mu(\text{Hct})$, to relate blood viscosity to hematocrit.

196 We explore two alternative closures. The first is based on
197 the assumption that any increase in overall pressure is propor-
198 tional to the increase in pressure drop across a given segment
199 according to

$$200 \quad P \propto P_z. \quad (9)$$

201 Physically, this approximation implies that a greater resis-
202 tance to flow requires a greater pressure to maintain flow
203 rate. Combining (1) and (9) yields

$$204 \quad P = C_p \frac{8Q\mu}{\pi r^4}, \quad (10)$$

205 where C_p , the constant of proportionality between P_z and
206 P , must be fit to data by regression or any other appropriate
207 method. The second closure, which is based on the pressure-
208 shear hypothesis of Pries and Secomb (2000), is discussed in
209 Appendix A.

210 Finally, it remains to specify a functional relationship $\mu =$
211 $\mu(\text{Hct})$. While the analysis and conclusions presented below
212 are valid for any monotonically increasing functional depen-
213 dence, in lieu of example we use the power-law relationship
214 (Pries et al. 1992):

$$215 \quad \frac{\mu}{\mu_{\text{plasma}}} = 1 + B \left[(1 - \text{Hct})^C - 1 \right], \quad (11)$$

216 where μ_{plasma} is the viscosity of plasma and values of the
217 experimentally determined constants B and C are given in
218 Table 2 of Pries et al. (1992). Their data for relation (11)
219 spans the range of Hct presented in this study, which is char-
220 acteristic of the variability of Hct in the healthy population.
221 Applicability of relation (11) is indicated to extend to patients
222 with polycythemia, whose Hct can exceed 50%.

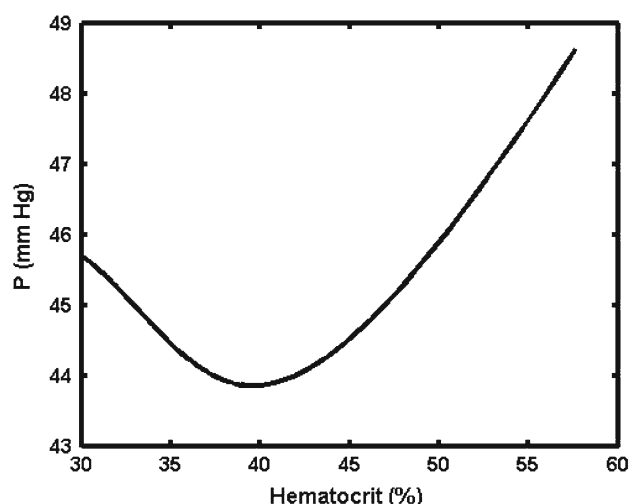


Fig. 1 Relationship between pressure (MAP) and Hct in a model of an elastic arteriolar microvessel, whose behavior is dictated by blood viscosity, a function of Hct. Transmural pressure is with reference to the outer pressure in the tissue, which is assumed to be atmospheric

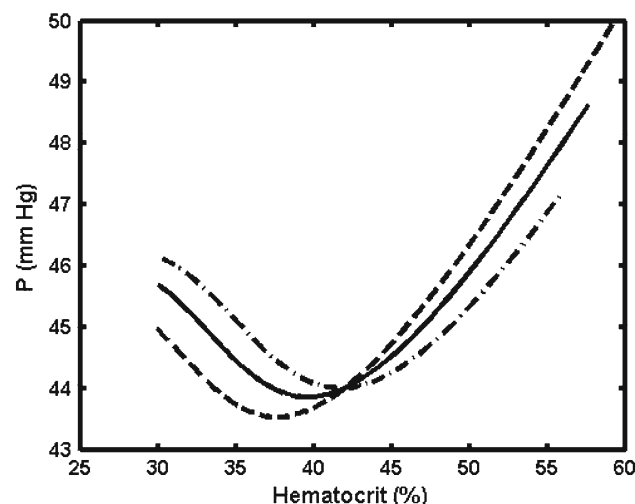


Fig. 2 Sensitivity to C_{passive} , the parameter that characterizes the tension at base conditions, i.e., the tension when $r = r_0$. The *solid line*, corresponding to $C_{\text{passive}} = 280$ dyn/cm, is the base state in Fig. 1. The *dash-dot* and *dashed lines* correspond to $C_{\text{passive}} = 240$ dyn/cm and $C_{\text{passive}} = 320$ dyn/cm, respectively

The MAP versus Hct relationship $P(\text{Hct})$ is found by combining Eqs. (8), (10), and (11). The resulting autoregulatory elastic compliance model provides a nonlinear implicit relationship between blood pressure P and vessel radius r . We used Matlab's nonlinear solver *fsolve* to numerically invert this implicit relation. As demonstrated in Fig. 1, this relation exhibits the "anomalous" U-shaped behavior observed by Salazar-Vázquez et al. (2006) in type 2 diabetic patients.

To understand why such behavior is observed in some patient groups but not others, we note that for it to occur, the function $P(\text{Hct})$ must reach its minimum within the physiological range of Hct. The function $P(\text{Hct})$ achieves its minimum at a point for which

$$\frac{dP}{d\text{Hct}} = \frac{dP}{dr} \frac{dr}{d\mu} \frac{d\mu}{d\text{Hct}} = 0. \quad (12)$$

Since both $r = r(\mu)$ in Eq. (10) and $\mu = \mu(\text{Hct})$ in Eq. (11) are monotonically increasing functions, it follows from Eq. (12) that for the relation $P(\text{Hct})$ to exhibit a U-shaped behavior (i.e., to have a minimum within the physiological range of Hct) requires that

$$\frac{dP}{dr} = 0. \quad (13)$$

Nonlinear expression (8) suggests that this condition is, in general, satisfied. Whether $P(\text{Hct})$ exhibits the U-shaped behavior within the physiological range of Hct is determined by the various modeling constants, whose values depend on the blood vessels' elastic properties and autoregulatory capacity. Consequently, we hypothesize that populations that

exhibit the anomalous MAP–Hct relationship tend to have blood vessels with abnormal elastic and autoregulatory properties. This is likely to be the case for diabetic, aged, etc., populations.

3 Results

We used our model to evaluate the equilibrium intraluminal average blood pressure in an elastic autoregulated arteriole-like blood vessel. The reliance on a single blood vessel is based on the assumption that the systemic result is the aggregate of effects at the different levels of branching in the arteriolar microcirculation. While our model allows for both pseudo-elastic descriptions of a blood vessel's wall and an arbitrary functional dependence between blood viscosity and Hct, the results presented in Figs. 1, 2, 3, 4, 5 and 6 correspond to an exponential elastic model (4) and a power-law viscosity–Hct relation (11). The use of nonlinear pseudo-elastic relationships and alternative viscosity–Hct constitutive models does not alter the predicted behavior. The MAP–Hct relationship presented in Fig. 1 corresponds to $r = 50\mu\text{m}$ and $Q = 10\text{nl/s}$, both of which are representative of the arterioles of awake hamsters; the rest of the parameters are given in Table 1. These values are deemed representative and similar to those in Carlson et al. (2008). The sensitivity of the proposed model to each of the main fitting parameters is illustrated in Figs. 2, 3, 4 and 5. It is important to realize that, according to expression (8), the results reported in these figures depend largely on the vessel stretching-ratio, r/r_0 , rather than on the vessel radius, r , alone.

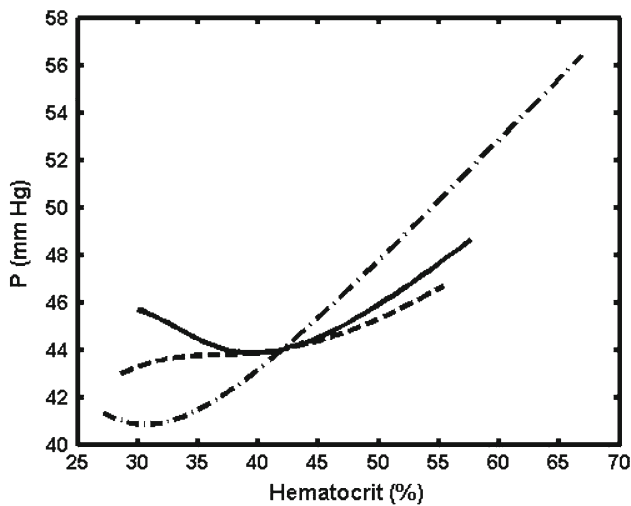


Fig. 3 Sensitivity to C'_{passive} , the parameter that accounts for the passive elasticity of the vessel, i.e. the resistance to deformation as the tension changes. The *solid line*, corresponding to $C'_{\text{passive}} = 11.5$, is the base state in Fig. 1. The *dash-dot* and *dashed lines* correspond to $C'_{\text{passive}} = 10.5$ and $C'_{\text{passive}} = 12.5$, respectively

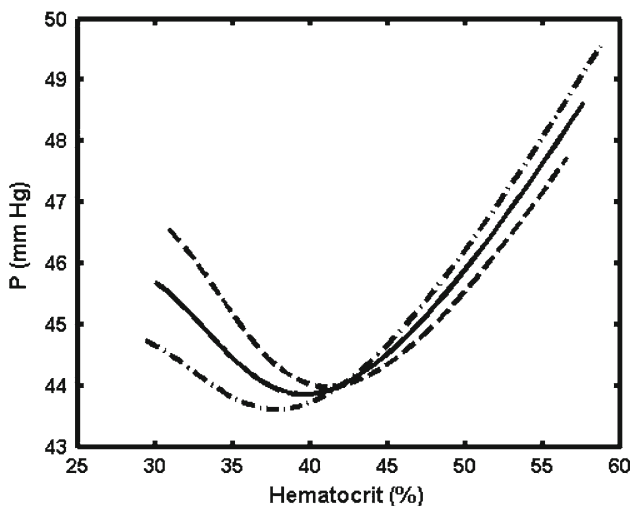


Fig. 4 Sensitivity to C_{active} , the parameter that characterizes the maximum possible value of the active component of tension. The *solid line*, corresponding to $C_{\text{active}} = 275$ dyn/cm, is the base state in Fig. 1. The *dash-dot* and *dashed lines* correspond to $C_{\text{active}} = 300$ dyn/cm and $C_{\text{active}} = 250$ dyn/cm, respectively

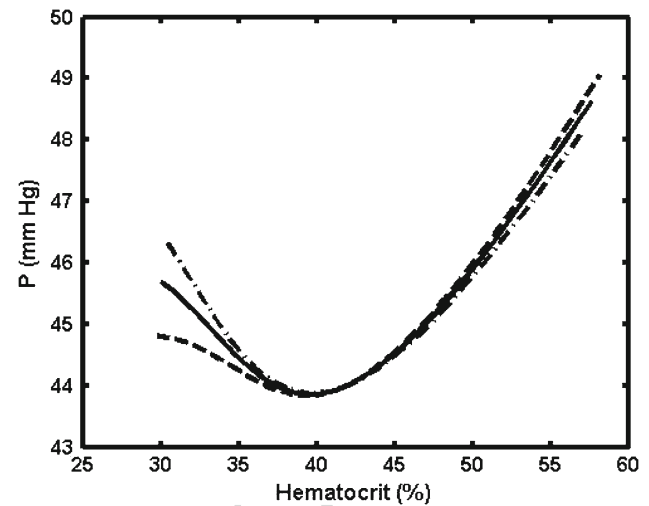


Fig. 5 Sensitivity to C'_{active} , the parameter that determines the radius at which the maximum active tension occurs. The *solid line*, corresponding to $C'_{\text{active}} = 0.75$, is the base state in Fig. 1. The *dash-dot* and *dashed lines* correspond to $C'_{\text{active}} = 0.45$ and $C'_{\text{active}} = 1.05$, respectively

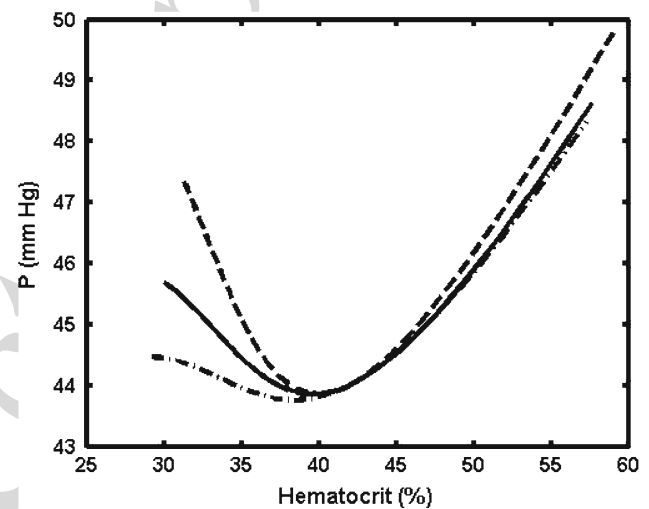


Fig. 6 Sensitivity to C''_{active} , the parameter that represents the range of radii where the active component of tension can act (similar to a variance on a normal distribution). The *solid line*, corresponding to $C''_{\text{active}} = 0.354$, is the base state in Fig. 1. The *dash-dot* and *dashed lines* correspond to $C''_{\text{active}} = 0.4$ and $C''_{\text{active}} = 0.3$, respectively

in vessels account for their elasticity by reference to vascular compliance (Keener and Sneyd 2004; Loscalzo and Schafer 2002). The proposed model describes how blood pressure changes lead to vessel diameter changes determined by vascular compliance and autoregulation.

Our analysis simplifies the problem by assuming steady flow and limiting the calculation to blood pressure in a single, straight (nontapering) tube. Our mathematical analysis shows that depending on Hct (or blood viscosity) the average blood pressure of an arteriolar vessel in the microvascular network may vary by several mmHg. Within a network this effect is

4 Discussion

The principal finding of our study is that the anomalous U-shaped dependence of MAP on Hct found in type 2 diabetic patients (Natali et al. 2005; Salazar-Vázquez et al. 2006) can be modeled by varying Hct (and thus, viscosity) of blood flowing in an elastic tube.

There is substantial literature on the analysis of flow relationships in elastic tubes, most of which is numerical (Quarneroni et al. 2000). Standard analytical models of blood flow

Table 1 Values of model parameters used in Fig. 1. The symbol in the value column refers to the source of this value

Constant	Value	Units
C_{passive}	280*	dyn/cm
C'_{passive}	11.5*	Unitless
C_{active}	275*	dyn/cm
C'_{active}	0.750*	Unitless
C''_{active}	0.354*	Unitless
C_{myog}	0.0359*	cm/dyn
C_{shear}	0.0258*	cm ² /dyn
C_{meta}	30*	1/($\mu\text{M cm}$)
C_{tone}	10.7*	Unitless
S_{CR}	0.45*	$\mu\text{M cm}$
C_{p}	1.6	M
B	2.15#	Unitless
C	0.875#	Unitless
μ_{plasma}	0.0013§	N s/m ²

* {Carlson et al. 2008 #1759}; # {Pries et al. 1992 #1536}; § {Koller et al. 1987 #1470}

cumulative, and the total pressure change for the whole circulatory system would depend on the number of segments or branches in which the effect occurs. As an example, Koller et al. (1987) found in the cat sartorius muscle an average of 4 segments (or order of branches) between arcading arterioles with diameter > 25 μm and capillaries.

There is no precise information on how many levels of arteriolar branching exist between the terminal arterioles and the small arteries with diameter $\sim 50 \mu\text{m}$ explored in our model, and how many vessel orders will exhibit the elastic behavior of our model as we ascend the arterial tree. The number is likely to be on the order of 20 arteriolar/small artery branches or more. Whether all of these vessels would show similar behavior is not known. However, according to Kassab (2006) “it has been shown that all vascular trees examined (coronary, pulmonary, vascular systems of various skeletal muscles, mesentery, omentum, and conjunctiva) in various species (rat, hamster, cat, rabbit, pig, and human) obey a set of design rules or scaling and scaling laws.” Thus, if some vessels present a specific behavior it may be assumed that, to a varying degree, this is a general feature for the circulation.

The constant flow assumption could be problematic since the circulation autoregulates flow in response to changes in the intrinsic oxygen-carrying capacity. We used this assumption solely to simplify the derivation and justify it by the very definition of autoregulation.

In our model, blood viscosity changes directly as a function of Hct. Under physiological conditions, factors such as a decrease in cardiac output (Shore and Tooke 1997) or vasoconstriction (Kassab 2006) may cause flow to decrease as Hct increases, an effect that occurs via the decrease in cardiac output (Richardson and Guyton 1959) or vasoconstriction

(Jefferson et al. 2002). This effect may be present in the diabetic population since MAP initially decreases as Hct increases. The subsequent increase in MAP could be due to the increase in blood viscosity which progressively increases resistance to flow and eventually MAP. However, in normal experimental animals subjected to small Hct changes, MAP decreased and cardiac output increased significantly as Hct was initially raised. This suggests that the oxygen autoregulatory control may not have a large gain for small changes in Hct (Martini et al. 2005).

While the U-shaped relationship between MAP and Hct is present in diabetic patients and appears to be a feature of a circulatory system that adapts to acute changes in blood viscosity, it is absent in the normal population. Thus, we hypothesize that populations with anomalous MAP–Hct relationships have blood vessels with abnormal elastic and autoregulatory properties. This is expected as healthy individuals have fully functioning autoregulation, the most important variable in determining peripheral resistance, where diabetics do not. Blood viscosity in the latter therefore exerts a proportionally greater influence.

The finding that small acute changes in viscosity in healthy animals anomalously affect MAP but is not evident in the healthy human population, explored cross-sectionally as a function of the individual’s Hct, indicates that there are significant differences between acute and chronic effects. In this context, effects in acute conditions were shown to be mediated by endothelial NO. Our findings indicate that in the presence of potentially diminished microvascular NO responsiveness these effects may also occur due to varying mechanical properties of blood vessels. Nevertheless, it is worthwhile emphasizing that our model, which accounts only for mechanical and fluid properties, is capable of modeling the different responses of healthy and diabetic populations.

The predictive power of the proposed model rests on the ability to measure the fitting parameters of our model in real blood vessels. This is not a trivial task, especially in vivo in humans. We present a sensitivity analysis to determine the influence of parametric uncertainty and measurement errors. The results of this analysis are presented in Figs. 2, 3, 4, 5 and 6. In general the presence of the U-shape does not appear to be particularly sensitive to the parameters other than for C'_{passive} in Fig. 3. In particular, decreasing C'_{passive} leads to a much sharper increase in pressure with increasing hematocrit, almost eliminating the observable presence of the U-shape. Figure 2 reveals that changing C_{passive} results in a shift of the minimum pressure. However, the overall form of the curves does not change. As C_{passive} increases, the minimum pressure decreases and is associated with lower levels of hematocrit. A similar shift in minimum pressure to lower levels of hematocrit is observable in Fig. 4 for changes in C_{active} . The sensitivity of pressure to C'_{active} is illustrated in Fig. 5. At higher levels of hematocrit, the pressure seems fairly

383 insensitive to changes in C'_{active} , whereas a more significant
 384 divergence is observable at lower hematocrit where smaller
 385 values of C'_{active} predict larger pressures. Finally, Fig. 6 shows
 386 that C''_{active} appears to play a very similar role as C'_{active} with
 387 little difference in predicted pressures at larger values of
 388 hematocrit and more significant differences at lower hema-
 389 tocrits.

390 In conclusion, we show that blood pressure may be in part
 391 regulated by the interaction of blood viscosity and blood ves-
 392 sel mechanical properties. The changes in blood viscosity are
 393 small and related to the natural variability of this parameter
 394 in the population. The normal population does not exhibit a
 395 U-shaped MAP versus Hct relationship. The presence of this
 396 relationship in the type 2 diabetic individuals suggests that
 397 blood pressure in this population is in part determined by
 398 different elastic properties of the vessel wall and the body's
 399 lesser ability to autoregulate.

400 An extension of these studies is to analyze more realis-
 401 tic models of the vasculature, although it is not clear that
 402 increased modeling complexity will provide insights beyond
 403 the conclusions presented here. Finally, our findings suggest
 404 that in diabetic patients there may be an Hct range that cor-
 405 responds to a comparatively lowered central blood pressure.
 406 While there is agreement that increased Hct and blood vis-
 407 cosity can be detrimental, it is apparent that lowered blood
 408 viscosity is not necessarily beneficial.

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415 **Appendix A: Alternative closure based on the** 416 **pressure–shear hypothesis**

417 The pressure–shear hypothesis was proposed by Pries and
 418 Secomb (2000) who studied the relationship between wall
 419 shear stress and intravascular pressure in arterioles, capil-
 420 laries and venules. They reported a sigmoidal relationship
 421 between shear stress and intravascular pressure that can be
 422 written as

$$423 \tau = \frac{C_1}{1 + C_3 \exp[-C_2(P - P_{\text{ref}})]}. \quad (\text{A1})$$

424 This relationship was shown to hold over a wide population
 425 of different vessels as a result of structural adaptation. For
 426 Poiseuille's law (1), the wall shear stress is given by

$$427 \tau = \frac{4\mu Q}{\pi r^3}. \quad (\text{A2})$$

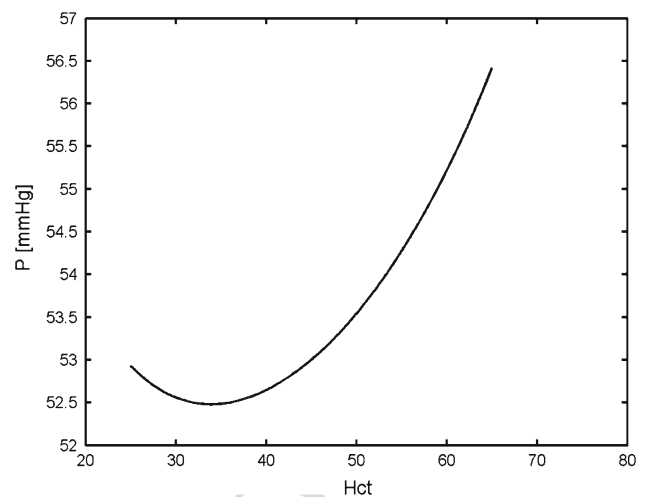


Fig. A1 Relationship between pressure (MAP) and Hct using the pressure–shear hypothesis, where the relationship between shear stress and intravascular pressure is sigmoidal. The model parameters are set to $C_1 = 100$, $C_2 = 2$, $C_3 = 1$, and $P_{\text{ref}} = 20$

Combining (A1) and (A2), using (2, 3, 4, 5, 6, 7) and (10), and noting that r is a function of P yields a relationship between pressure and hematocrit. A sample figure, displaying the U-shape obtained with this closure approximation is presented in Fig. A1.

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