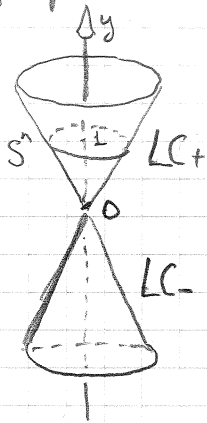


- Consider the case of  $\mathbb{R}^n$ , i.e.  $p=n, q=0$

$SO(n+1, 1)$  acts on  $\mathbb{R}^{n+1, 1}$  by linear isometries and preserves the light-cone

$$LC = \{ (x^0, \dots, x^n, y) \in \mathbb{R}^{n+1, 1} \mid (x^0)^2 + \dots + (x^n)^2 - y^2 = 0 \} \subset \mathbb{R}^{n+1, 1}$$

"orthochronous Lorentz group"  $SO_+(n+1, 1)$  preserves the positive light-cone



$$LC_+ = \{ (x^0, \dots, x^n, y) \in \mathbb{R}^{n+1, 1} \mid (x^0)^2 + \dots + (x^n)^2 - y^2 = 0, y > 0 \} \subset LC$$

We have two commuting actions:

$$SO_+(n+1, 1) \curvearrowright LC_+ \supset \mathbb{R}_+ \quad \leftarrow \text{dilations}$$

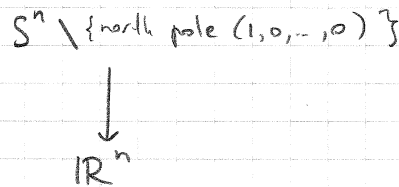
i.e.  $SO_+(n+1, 1)$  acts on  $LC_+/\mathbb{R}_+ \sim S^n$  - section of  $\mathbb{R}_+^+$  action on  $LC_+$  singled out by  $y=1$

$LC_+$  inherits degenerate metric from  $\mathbb{R}^{n+1, 1}$  (with 1-dim. kernel)

$\leadsto LC_+/\mathbb{R}_+$  inherits non-deg. conformal structure (the kernel is exactly killed by quotienting over  $\mathbb{R}_+$ )

so,  $SO_+(n+1, 1)$  acts on  $S^n$  by conformal maps

then we use stereographic projection



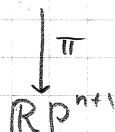
Along the way, we identified  $S^n$  as the conformal compactification of  $\mathbb{R}^n$ :

conf. vect. fields extend to  $S^n$ , finite conf. maps are everywhere well-defined on  $S^n$

General case ( $\mathbb{R}^{p, q}$ )

$SO(p+1, q+1)$  acts on  $\mathbb{R}^{p+1, q+1}$  and preserves  $LC = \{ (x^0, \dots, x^p, y^0, \dots, y^q) \in \mathbb{R}^{p+1, q+1} \mid \sum_{i=0}^p (x^i)^2 - \sum_{j=0}^q (y^j)^2 = 0 \}$

$$SO(p+1, q+1) \curvearrowright LC \setminus \{0\} \supset \mathbb{R}_+$$



Denote the image of the light-cone in  $\mathbb{R}P^{n+1}$  by  $\mathcal{N}^{p, q}$ .  $\mathcal{N}^{p, q}$  inherits conformal structure

from  $\mathbb{R}^{p+1, q+1}$  and  $SO(p+1, q+1)$  acts on  $\mathcal{N}^{p, q}$  by conformal maps,

$$L: \mathbb{R}^{p, q} \rightarrow \mathcal{N}^{p, q} \quad \text{-injective, } \text{im}(L) \text{ is } q \text{ open-dense in } \mathcal{N}^{p, q}$$

$$(x^0, \dots, x^p, y^1, \dots, y^q) \mapsto \left( \frac{1 - \sum_{i=0}^p (x^i)^2 + \sum_{j=0}^q (y^j)^2}{2} : x^0 : \dots : x^p : \frac{1 + \sum_{i=0}^p (x^i)^2 - \sum_{j=0}^q (y^j)^2}{2} : y^1 : \dots : y^q \right)$$

•  $\mathcal{N}^{p, q}$  is the conf. compactification of  $\mathbb{R}^{p, q}$

•  $S^p \times S^q \xrightarrow{\pi} \mathcal{N}^{p, q}$  is the double cover

# Conformal symmetry of $\mathbb{R}^2$

equation for a conf. vect. field:  $\partial_i \epsilon_j + \partial_j \epsilon_i = \omega \delta_{ij} \iff \begin{cases} \partial_x \epsilon_x = \partial_y \epsilon_y = \frac{1}{2} \omega \\ \partial_x \epsilon_y = -\partial_y \epsilon_x \end{cases}$

(notation:  $x = x^1, y = x^2$ )  $\iff \epsilon_x + i \epsilon_y$  satisfies Cauchy-Riemann equations

$\iff$  the vector field  $\epsilon_i \partial_i$  is of the form

$$\epsilon(z) \frac{\partial}{\partial z} + \bar{\epsilon}(\bar{z}) \frac{\partial}{\partial \bar{z}}$$

↑ holom. vect. field      ↑ conjugate anti-hol. v.f.

notation:  $z = x + iy, \bar{z} = x - iy$

$$\partial = \partial_z = \frac{1}{2}(\partial_x - i \partial_y), \bar{\partial} = \partial_{\bar{z}} = \frac{1}{2}(\partial_x + i \partial_y)$$

conf. factor:  $\omega = \partial \epsilon + \bar{\partial} \bar{\epsilon}$

Thus:  $\left\{ \begin{array}{l} \text{conformal vect. fields} \\ \text{on } \mathbb{R}^2 \simeq \mathbb{C} \\ \epsilon_x \partial_x + \epsilon_y \partial_y \\ (\text{Re } \epsilon(z)) \partial_x + (\text{Im } \epsilon(z)) \partial_y \end{array} \right\} \xrightarrow{\cong} \left\{ \begin{array}{l} \text{holom. vect.} \\ \text{fields on } \mathbb{C} \\ (\epsilon_x + i \epsilon_y) \cdot \partial_z \\ \epsilon(z) \partial_z \end{array} \right\}$  Lie algebra homomorphism

## Finite version:

A diffeo  $\varphi: \mathbb{D} \xrightarrow{\sim} \mathbb{D}'$  is conformal iff  $\varphi$  is either holomorphic or anti-holomorphic

Proof:  $\varphi^* g = \frac{\partial \varphi^i}{\partial x^j} \frac{\partial \varphi^i}{\partial x^k} dx^j dx^k = \frac{\partial \varphi}{\partial z} \frac{\partial \bar{\varphi}}{\partial \bar{z}} (dz)^2 + \frac{\partial \varphi}{\partial z} \frac{\partial \bar{\varphi}}{\partial \bar{z}} dz d\bar{z} + \frac{\partial \varphi}{\partial \bar{z}} \frac{\partial \bar{\varphi}}{\partial z} d\bar{z} dz + \frac{\partial \varphi}{\partial \bar{z}} \frac{\partial \bar{\varphi}}{\partial \bar{z}} (d\bar{z})^2$   
 $dx^i dx^i = dz d\bar{z}$       setting  $\varphi = \varphi^1 + i \varphi^2$   
 $\bar{\varphi} = \varphi^1 - i \varphi^2$

therefore  $\varphi^* g = \Omega g \iff$  either  $\bar{\partial} \varphi = 0$ , then  $\Omega = |\partial \varphi|^2$   
 or  $\partial \varphi = 0$ , then  $\Omega = |\bar{\partial} \varphi|^2$        $\square$

•  $\text{conf}(\mathbb{C} \setminus \{0\}) = \left\{ \begin{array}{l} \text{real parts of merom. vect. f. on } \mathbb{C} \\ \text{with pole at } 0 \text{ allowed} \end{array} \right\}$

Introduce the Witt algebra  $\mathcal{A} = \left\{ \sum_{n=-\infty}^{\infty} c_n l_n \mid c_n \in \mathbb{C}, l_n = -z^{n+1} \frac{\partial}{\partial z} \right\} = \left\{ \begin{array}{l} \text{merom. v.f. on } \mathbb{C} \\ \text{with pole at } 0 \text{ allowed} \end{array} \right\}$

Generators  $\{l_n\}$  satisfy commutation relations

$$[l_m, l_n] = (m-n) l_{m+n}$$

$$l_n = -z^{n+1} \partial_z, \bar{l}_n = -\bar{z}^{n+1} \partial_{\bar{z}}$$

$\text{conf}(\mathbb{C} \setminus \{0\}) \xrightarrow{\cong} \underbrace{\mathcal{A}}_{\text{Span}_{\mathbb{C}} \{l_n\}} \oplus \underbrace{\bar{\mathcal{A}}}_{\text{Span}_{\mathbb{C}} \{\bar{l}_n\}}$

$$\begin{cases} [l_m, l_n] = (m-n) l_{m+n} \\ [\bar{l}_m, \bar{l}_n] = (m-n) \bar{l}_{m+n} \\ [l_m, \bar{l}_n] = 0 \end{cases}$$

$\parallel$   
 $\text{Span}_{\mathbb{R}} \{l_n + \bar{l}_n, i(l_n - \bar{l}_n)\}_{n=-\infty}^{\infty}$

Exercise: show that  $\mathcal{A}$  is the complexification of the Lie algebra of vector fields on a circle

- $\text{conf}(\mathbb{C}) = \text{Span}_{\mathbb{R}} \{l_n + \bar{l}_n, i(l_n - \bar{l}_n)\}_{n \geq -1}$
- $\left\{ \begin{array}{l} \text{conf. v.f. on } \mathbb{C} \\ \text{vanishing at } 0 \end{array} \right\} = \text{Span}_{\mathbb{R}} \{ \dots \}_{n \geq 0}$
- $\text{conf}(\mathbb{C}P^1) = \text{Span} \{l_n + \bar{l}_n, i(l_n - \bar{l}_n)\}_{n=-1,0,1} \cong \mathfrak{sl}_2(\mathbb{C}) \cong \mathfrak{so}(3,1)$   
 $\stackrel{\parallel}{\mathbb{C}}$

$$\begin{array}{l}
 -(l_{-1} + \bar{l}_{-1}) = \partial_x \\
 -i(l_{-1} - \bar{l}_{-1}) = \partial_y
 \end{array}
 \left. \vphantom{\begin{array}{l} \\ \\ \end{array}} \right\} \text{translations}
 \quad
 \begin{array}{l}
 -(l_0 + \bar{l}_0) = x\partial_x + y\partial_y \\
 -i(l_0 - \bar{l}_0) = -y\partial_x + x\partial_y
 \end{array}
 \left. \vphantom{\begin{array}{l} \\ \\ \end{array}} \right\} \begin{array}{l} \text{dilatation} \\ \text{rotation} \end{array}$$

$$\begin{array}{l}
 -(l_1 + \bar{l}_1) = (x^2 - y^2)\partial_x + 2xy\partial_y \\
 -i(l_1 - \bar{l}_1) = -2xy\partial_x + (x^2 - y^2)\partial_y
 \end{array}
 \left. \vphantom{\begin{array}{l} \\ \\ \end{array}} \right\} \begin{array}{l} \text{special conformal} \\ \text{transformations} \end{array}$$

- $\text{Conf}(\mathbb{C}P^1) = \text{PSL}_2(\mathbb{C}) \cong \text{SO}_+(3,1)$   
Möbius transformations

- Despite the fact that  $\text{conf}(\mathbb{C})$  is  $\infty$ -dimensional, global conf. automorphisms of  $\mathbb{C}$  comprise only a finite-dimensional group.
- $\mathbb{C}$  does not have a conf. compactification (in the sense of a compact mfd containing  $\mathbb{C}$ , to which any conf. vect. field on  $\mathbb{C}$  can be extended)
- Riemann's mapping thm: any two simply-connected domains  $D, D' \subset \mathbb{C}$  are conformally equivalent

Naïvely,  $\mathbb{C} \setminus \{0\}$ , punctured disc  $D \setminus \{0\} = \{z \in \mathbb{C} \mid 0 < |z| < 1\}$  and annulus  $\text{Ann}_{r,R} = \{z \in \mathbb{C} \mid r < |z| < R\}$  have the same Lie algebras of conf. vect. fields:

$$\text{conf}(\mathbb{C} \setminus \{0\}) \cong \text{conf}(D \setminus \{0\}) \cong \text{conf}(\text{Ann}_{r,R}) \cong \mathcal{A} = \text{Span}_{\mathbb{R}} \{l_n\}_{-\infty}^{\infty}$$

But there is a subtlety: different convergence restrictions for coefficients  $\{c_n\}$  in Laurent expansion  $\mathcal{E}(z)\partial_z = \sum_{n=-\infty}^{\infty} c_n l_n$  for these domains

Conformal symmetry of  $\mathbb{R}^1$  (trivial case)

Riem. metric  $g = (dx)^2$        $\left\{ \begin{array}{l} \text{conformal diffeo} \\ \varphi: \mathbb{R}^1 \rightarrow \mathbb{R}^1 \end{array} \right\} = \left\{ \begin{array}{l} \text{all diffeo} \\ \varphi: \mathbb{R}^1 \rightarrow \mathbb{R}^1 \end{array} \right\}$       conf. factor:  $\Omega = \left(\frac{\partial \varphi}{\partial x}\right)^2$

$\text{conf}(\mathbb{R}^1) = \text{Vect}(\mathbb{R}^1)$ ; vect. field  $\mathcal{E}(x)\partial_x$  has conf. factor  $\omega = 2\partial_x \mathcal{E}(x)$

$\text{Conf}(\mathbb{S}^1) = \text{Diff}(\mathbb{S}^1) \supset \text{PSL}_2(\mathbb{R}) \cong \text{SO}_+(2,1)$  - "restricted conformal group of  $\mathbb{R}^1$ "  
 $\stackrel{\parallel}{\mathbb{R}^2}$       Möbius transf. of  $\mathbb{S}^1$

Again:  $\mathbb{S}^1$  is not the true conf. compactification of  $\mathbb{R}^1$ , it is rather the "imposed" compactification

# Conformal symmetry of Minkowski plane $\mathbb{R}^{1,1}$

Minkowski metric:  $g = (dx)^2 - (dy)^2 = h_{ij} dx^i dx^j$  where  $h_{ij} = \begin{pmatrix} 1 & \\ & -1 \end{pmatrix}$

Switch to light-cone coordinates:

$$\begin{cases} x_+ = x + y \\ x_- = x - y \end{cases}$$

useful f-lae:  $\begin{cases} x = \frac{x_+ + x_-}{2} \\ y = \frac{x_+ - x_-}{2} \end{cases}$

$$\partial_+ := \frac{\partial}{\partial x_+} = \frac{\partial_x + \partial_y}{2}$$

$$\partial_x = \partial_+ + \partial_-$$

$$\partial_- := \frac{\partial}{\partial x_-} = \frac{\partial_x - \partial_y}{2}$$

$$\partial_y = \partial_+ - \partial_-$$

$$g = dx_+ dx_- , \quad h_{ij} = \begin{pmatrix} & 1/2 \\ 1/2 & \end{pmatrix}$$

Equation  $\partial_i \varepsilon_j + \partial_j \varepsilon_i = \omega h_{ij}$  for

components of a conf. vector field  $\varepsilon$

$$\varepsilon^i \partial_i = \varepsilon_+(x_+, x_-) \partial_+ + \varepsilon_-(x_+, x_-) \partial_- \quad \text{reads:}$$

$$\begin{cases} \partial_- \varepsilon_+ = 0 \\ \partial_+ \varepsilon_- = 0 \\ \partial_+ \varepsilon_+ + \partial_- \varepsilon_- = \omega \end{cases}$$

Therefore: a generic conf. v.f. on  $\mathbb{R}^{1,1}$  is a vect. field

of the form  $\varepsilon_+(x_+) \partial_+ + \varepsilon_-(x_-) \partial_-$

it has conf. factor  $\omega = \partial_+ \varepsilon_+ + \partial_- \varepsilon_-$

Thus  $\text{conf}(\mathbb{R}^{1,1}) \cong \underbrace{\text{Vect}(\mathbb{R})}_{\varepsilon_+ \partial_+} \oplus \underbrace{\text{Vect}(\mathbb{R})}_{\varepsilon_- \partial_-}$

terminology:

functions of  $x_+$  - "right-movers"  
functions of  $x_-$  - "left-movers"

Now consider conf. maps

$$\varphi: \mathbb{R}^{1,1} \rightarrow \mathbb{R}^{1,1}$$

$$(x_+, x_-) \mapsto (\varphi_+, \varphi_-)$$

$$\begin{aligned} \text{We have: } \varphi^* g &= \varphi^*(dx_+ dx_-) = d\varphi_+ d\varphi_- = \\ &= \frac{\partial \varphi_+}{\partial x_+} \frac{\partial \varphi_-}{\partial x_+} (dx_+)^2 + \frac{\partial \varphi_+}{\partial x_+} \frac{\partial \varphi_-}{\partial x_-} dx_+ dx_- + \frac{\partial \varphi_+}{\partial x_-} \frac{\partial \varphi_-}{\partial x_+} dx_- dx_+ + \\ &\quad + \frac{\partial \varphi_+}{\partial x_-} \frac{\partial \varphi_-}{\partial x_-} (dx_-)^2 \end{aligned}$$

Hence:

$\varphi$  is conformal  $\Leftrightarrow$  (1) either  $\varphi_+ = \varphi_+(x_+), \varphi_- = \varphi_-(x_-)$   
i.e.  $\varphi \in \text{Diff}(\mathbb{R}) \times \text{Diff}(\mathbb{R})$  is a reparametrization of  $x_+$  and  $x_-$  (independently)

(2) Or  $\varphi_+ = \varphi_+(x_-), \varphi_- = \varphi_-(x_+)$   
i.e.  $\varphi = \left( \begin{smallmatrix} \text{reparam.} \\ \text{of } x_+, x_- \end{smallmatrix} \right) \circ \left( \begin{smallmatrix} \text{reflection} \\ (x, y) \mapsto (x, -y) \end{smallmatrix} \right)$

Conformal factors:

Case (1):  $\Omega = (\partial_+ \varphi_+) (\partial_- \varphi_-)$

Case (2):  $\Omega = (\partial_- \varphi_+) (\partial_+ \varphi_-)$

•  $\text{Conf}_0(\mathbb{R}^{1,1}) = \text{Diff}_+(\mathbb{R}) \times \text{Diff}_+(\mathbb{R})$  (the whole  $\text{Conf}(\mathbb{R}^{1,1})$  has 8 conn. components)

• "True" conf. compactification does not exist, but one may impose  $\overline{\mathbb{R}^{1,1}} = S^1 \times S^1$   
 $\uparrow \quad \quad \quad \uparrow$   
 $x_+ \quad \quad \quad x_-$

•  $\text{Conf}_0(\overline{\mathbb{R}^{1,1}}) = \text{Diff}_+(S^1) \times \text{Diff}_+(S^1) \supset \text{PSL}_2(\mathbb{R}) \times \text{PSL}_2(\mathbb{R}) \cong \text{SO}(2, 2) \leftarrow \begin{matrix} \text{Möbius} \times \text{Möbius} \\ \text{restricted} \\ \text{conf. group} \end{matrix}$

Exercise:  $\left. \begin{matrix} \text{identity} \\ \text{1 Lorentz boost} \\ \text{1 dilatation} \\ \text{2 SCTs} \end{matrix} \right\}$  as Möb $_+$   $\times$  Möb $_+$  -trans:  $(x_+, x_-) \mapsto \begin{pmatrix} a x_+ + b_+ \\ c x_+ + d_+ \end{pmatrix}, \begin{pmatrix} a x_- + b_- \\ c x_- + d_- \end{pmatrix}$

# Moduli of conformal structures (Some remarks)

def: A (pseudo-)Riemannian manifold  $(M, g)$  is called conformally flat if one can choose coord. charts on  $M$ , so that in each chart

$$g = \Omega(x) h_{ij} dx^i dx^j \quad \text{with } h_{ij} = \begin{pmatrix} 1 & & & \\ & 1 & & \\ & & -1 & \\ & & & -1 \end{pmatrix} \text{ and } \Omega(x) > 0$$

Rem: Being conformally flat is a local property

- For  $\dim(M) = 1, 2$  all (pseudo-)Riemannian manifolds are conformally flat (in case  $\dim(M) = 1$ , actually globally flat)

Exercise: prove this

- For  $\dim(M) > 2$  not every metric is conf. flat:

in case  $\dim M \geq 4$  metric  $g$  is conf. flat  $\Leftrightarrow$  Weyl curvature tensor  $W(g)$  vanishes

in case  $\dim M = 3$ :

metric  $g$  is conf. flat  $\Leftrightarrow$  Cotton tensor  $C(g)$  vanishes

(0,3) tensor, constructed out of cov. derivatives of Ricci tensor, scalar curvature and metric

(0,4) tensor constructed out of Riem. curvature, Ricci tensor, scalar curvature and metric

def Moduli space of conf. structures on a smooth mfd.  $M$  is

$$\mathcal{M}_M := \left\{ \begin{array}{l} \text{conf. structures} \\ \text{on } M \end{array} \right\} / \text{Diff}(M)$$

- action  $\text{Diff}(M) \curvearrowright \left\{ \begin{array}{l} \text{conf. str.} \\ \text{on } M \end{array} \right\}$  is not free and has stabilizer  $\text{Conf}(M, \gamma)$

for a conf. structure  $\gamma = g/h$ , i.e. we have the following picture:

$$0 \rightarrow \text{Conf}(M, \gamma) \rightarrow \text{Diff}(M) \curvearrowright \left\{ \begin{array}{l} \text{conf. str. } \gamma \\ \text{on } M \end{array} \right\}$$

$$\downarrow$$

$$\mathcal{M}_M$$

- Or infinitesimally:

$$0 \rightarrow \text{conf}(M, \gamma) \rightarrow \text{Vect}(M) \xrightarrow{\text{infinitesimal action}} \left\{ \begin{array}{l} \text{conf. str. } \gamma \\ \text{on } M \end{array} \right\}$$

(i.e.  $\text{Vect}(M)$  determines an integrable distribution on  $\left\{ \begin{array}{l} \text{conf. str.} \\ \text{on } M \end{array} \right\}$  whose leaf space is  $\mathcal{M}_M$ )

$$\downarrow$$

$$\mathcal{M}_M$$

Case  $\dim(M) > 2$

if  $\dim(M) \geq 4$  Weyl curvature tensor  $W(g)$  is invar. under Weyl transf.  $g \mapsto g' = \Omega \cdot g$   
 $\Rightarrow W(g) = W(g/\sim)$

for  $\dim(M) = 3$ , Cotton tensor  $C(g)$  is invar. under Weyl transf.  $\Rightarrow C(g) = C(g/\sim)$

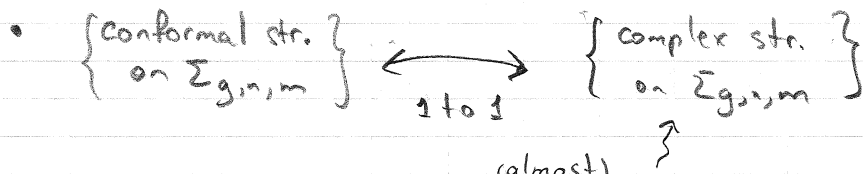
Therefore, for  $\dim(M) > 2$  conf. structures on  $M$  have local moduli:  $W(x), C(x)$

so  $\dim \mathcal{M}_M = \infty$

Case  $\dim(M) = 2$ , conf. structures of signature  $(2, 0)$

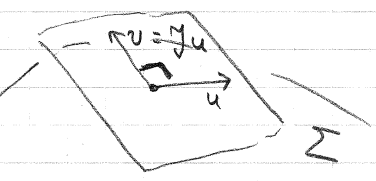
notation:  $\Sigma_{g,n,m}$  - 2-dim smooth oriented mfd ;  $\Sigma_{g,n} := \Sigma_{g,n,0}$   
 ↑ genus     ↑ # of punctures     ← # of boundary circles

Rem: we only consider cx structures consistent with chosen orientation; and only  $(2,0)$ -conf. str.



Reminder: <sup>(almost)</sup> cx str. on  $\Sigma$  is a section  $\gamma \in \Gamma(\Sigma, \text{End}(T\Sigma))$  with  $\gamma_x^2 = -1$  for all  $x \in \Sigma$   
 cx str. = almost cx. + an integrability condition which holds automatically for  $\dim \Sigma = 2$

$\rightarrow$ : given a conf. str.  $g$  on  $\Sigma$ , we construct  $J_x: T_x \Sigma \rightarrow T_x \Sigma$   
 $\begin{matrix} \underbrace{g} & & \underbrace{g/\sim} \\ \underbrace{u} & \mapsto & \underbrace{v} \end{matrix}$



where  $v$  is orthogonal to  $u$  and of the same length (w.r.t.  $g$ ) and such that the pair  $(u, v)$  is positively oriented.

$\leftarrow$ : given a cx. str.  $\gamma$  on  $\Sigma$ , choose arbitrary Riem. metric  $\tilde{g}$  on  $\Sigma$  and deform it to

$g_x(u, v) := \tilde{g}(u, v')$  where  $v'$  is the projection of  $v$  along  $J_x u$  to the line in  $T_x \Sigma$  containing  $u$

i.e.  $v' = v - \alpha J_x u = \beta u$  for some  $\alpha, \beta \in \mathbb{R}$

then take the conf. class  $g/\sim$

- Terminology:  $\Sigma_{g,n,m}$  endowed with conf. (or cx.) str. is called a Riemann surface (not a Riemannian mfd!)