## INTERMEDIATE GEOMETRY AND TOPOLOGY EXERCISES 10, 11/12/2021. HAMILTONIAN VECTOR FIELDS.

1. (Harmonic oscillator.) Consider the plane $\mathbb{R}^{2}$ with coordinates $q, p$ and standard symplectic form $\omega=d q \wedge d p$. Let $H=\frac{1}{2}\left(p^{2}+q^{2}\right)$. Find the corresponding Hamiltonian vector field $X_{H}$ and its flow in time $t \in \mathbb{R}$.
2. (Pendulum.) Consider the symplectic manifold $M=T^{*} S^{1}$ with coordinate $\theta \in \mathbb{R} / 2 \pi \mathbb{Z}$ on the base and $p$ the coordinate on the cotangent fiber, and with the canonical symplectic form of the cotangent bundle, $\omega=d \theta \wedge d p$. Let

$$
H=\frac{1}{2} p^{2}-\cos \theta
$$

Find the corresponding Hamiltonian vector field $X_{H}$. Analyze qualitatively the flow $\rho_{t}$ of $X_{H}$ :
(a) Draw a sketch of the integral curves of $X_{H}$ (orbits of $\rho_{t}$ ).
(b) Does $\rho_{t}$ have constant orbits? Which of them are stable (perturbing the initial condition yields an orbit that stays near the constant one), which are unstable?
(c) Does $\rho_{t}$ have non-closed orbits?
(d) Do closed orbits of $\rho_{t}$ all have the same period? If not, how does the period behave depending on the orbit?
(e) Find a closed (integral) formula for the period of a periodic orbit.
3. (Angular momentum.) Consider $M=T^{*} \mathbb{R}^{3}$; denote $q \in \mathbb{R}^{3}$ a point in the base and $p \in\left(\mathbb{R}^{3}\right)^{*} \simeq \mathbb{R}^{3}$ a point in the cotangent fiber. For any vector $\nu \in \mathbb{R}^{3}$, define a Hamiltonian function $J_{\nu}=(\nu, q \times p)$ where (, ) is the interior product of vectors in $\mathbb{R}^{3}$ and $\times$ the exterior product.
(a) Describe the flow of the Hamiltonian vector field $X_{J_{\nu}}$.
(b) Show that the Poisson brackets ${ }^{1}$ are $\left\{J_{\nu}, J_{\mu}\right\}=J_{\nu \times \mu}$.
(c) Show that $J$ defines a homomorphism of Lie algebras $\mathfrak{s o}(3) \xrightarrow{J} C^{\infty}\left(T^{*} \mathbb{R}^{3}\right)$ which fits into a sequence of homomorphisms

$$
\mathfrak{s o}(3) \xrightarrow{J} C^{\infty}\left(T^{*} \mathbb{R}^{3}\right) \xrightarrow{X \ldots} \mathfrak{X}\left(\mathbb{R}^{3}\right)
$$

Here the Lie algebra structure on $C^{\infty}$ is given by the Poisson bracket and Lie algebra structure on vector fields is the usual Lie bracket of vector fields.

## 4. (Integrals of motion.)

(a) Let $(M, \omega)$ be a symplectic manifold, $H$ a Hamiltonian function on $M$ and $I$ another function such that $\{H, I\}=0 .^{2}$ Show that if $\gamma$ is an integral curve of the Hamiltonian vector field $X_{H}$, then $I$ is constant along $\gamma$ (i.e. $\left.\frac{d}{d t} I(\gamma(t))=0\right)$.

[^0](b) Let $M=\prod_{i=1}^{n} T^{*} \mathbb{R}$. A point in $M$ is given by $q_{1}, \ldots, q_{n} \in \mathbb{R}$ and $p_{1}, \ldots, p_{n} \in$ $\mathbb{R}^{*}$; the symplectic form is $\sum_{i} d q_{i} \wedge d p_{i}$. Let $H=\sum_{i} \frac{1}{2} p_{i}^{2}+\sum_{i \neq j} V\left(q_{i}-q_{j}\right)$ with $V \in C^{\infty}(\mathbb{R})$ some fixed function. ${ }^{3}$ Show that the "total momentum" $P=\sum_{i} p_{i}$ is an integral of motion. Describe the flow of the Hamiltonian vector field $X_{P}$ generated by $P$.
(c) Let $M=T^{*} \mathbb{R}^{3}$ with canonical symplectic structure, let $H=\frac{(p, p)}{2}+V(q)$ with $V(q)=f(\|q\|)$ a function on $\mathbb{R}^{3}$ depending only on the norm of $q$. Show that in this system, the angular momentum $J_{\nu}=(\nu, q \times p)$ for any vector $\nu \in \mathbb{R}^{3}$ is an integral of motion.
5. Let $(M, \omega)$ be a symplectic manifold of dimension $2 n$ and let $H_{1}, \ldots, H_{k}$ be a collection of functions on $M$ satisfying $\left\{H_{i}, H_{j}\right\}=0$ for any $i, j$. Consider the $\operatorname{map} M \rightarrow \mathbb{R}^{m}$ given by $\mu(x)=\left(H_{1}(x), \ldots, H_{m}(x)\right)$. Assuming that $c \in \mathbb{R}^{m}$ is a regular value of $\mu$, show that $\mu^{-1}(c)$ is a coisotropic submanifold of $M$. Show that if additionally $m=n$, then $\mu^{-1}(c)$ is a Lagrangian submanifold of $M .{ }^{4}$

[^1]
[^0]:    ${ }^{1}$ For $f, g$ two functions on a symplectic manifold $(M, \omega)$, the Poisson bracket is defined as $\{f, g\}:=\mathcal{L}_{X_{f}} g-$ the Lie derivative of $g$ along the Hamiltonian vector field corresponding to f .
    ${ }^{2}$ Such $I$ is called an "integral of motion" - where the motion is understood as determined by $X_{H}$.

[^1]:    ${ }^{3}$ The physical interpretation of this system is: $n$ particles on a real line (with positions $q_{i}$ and momenta $p_{i}$ ), of mass 1 , with a pairwise interaction via a force potential $V$ depending only on the distance.
    ${ }^{4}$ Such a situation - a maximal collection of Poisson-commuting Hamiltonians on a symplectic manifold is called an (Liouville-) integrable system.

