

# Torsion as an Integral

IDEA: WANT TO write  $\tau(X)$  - torsion as manifold or CW complex

a field theory-type integral on  $X$  as a "space-time"  $\int_{\mathbb{F}_X} e \in S_X(\varphi) D\varphi$  field on  $X$

Algebraic story:

$$C^0 \xrightarrow{d} C^1 \xrightarrow{d} \dots \xrightarrow{d} C^n \quad \text{cochain complex of vector spaces}/\mathbb{R} \text{ with inner product } (\cdot, \cdot) \text{ on } C^j$$

Assume  $(C^\bullet, d)$  acyclic. Choose a chain contraction  $K$ -section  $C^j \xrightarrow{d} C^{j+1}$  exact

Thus:  $C^j = \underbrace{C^j_{\text{ex}}}_{\text{im}(d)} \oplus \underbrace{C^j_{\text{coex}}}_{\text{im}(K)}$



Torsion:  $\tau(C^\bullet, d, (\cdot, \cdot)) = \text{"Sdet}(d) = \prod_{j=0}^{n-1} (\det_{C^j_{\text{coex}} \rightarrow C^{j+1}_{\text{ex}}}(d))^{(-1)^j}$   
 $= \prod_{j=0}^{n-1} (\det_{C^j} \Delta)^{-\frac{(-1)^j j}{2}} \in \mathbb{R}_{>0}$  where  $\Delta = d d^T + d^T d : C^j \rightarrow C^j$

$\tau$  does not depend on  $K$

Non-acyclic case:

choose an  $(i, p, K)$  triple



$$\begin{aligned} p \circ i &= \text{id} \\ d \circ i &= p \circ d = 0 \\ \boxed{dK + Kd} &= \text{id} - i \circ p \\ K^2 &= 0 \\ Ki &= pK = 0 \end{aligned}$$

Hodge decomposition:

$$C^j = \underbrace{C^j_{\text{harm}}}_{\text{im}(i)} \oplus \underbrace{C^j_{\text{ex}}}_{\text{im}(d)} \oplus \underbrace{C^j_{\text{coex}}}_{\text{im}(K)}$$



$$\tilde{\tau} = \prod_{j=0}^{n-1} (\det_{C^j_{\text{coex}} \rightarrow C^{j+1}_{\text{ex}}}(d))^{(-1)^j} = \prod_{j=0}^{n-1} (\det'_{C^j} \Delta)^{-\frac{(-1)^j j}{2}} \in \mathbb{R}_{>0}$$



determinant lines  $\text{Det } C^0 = \bigotimes_{j=0}^n \mathbb{D}(\Lambda^{\text{top}} C^j)^{(-1)^j}$  notation:  $L^{-1} := L^*$  Torsion  
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differential  $d$  induces an isomorphism of determinant lines

$$\mathbb{T}: \text{Det } C^0 \xrightarrow{\sim} \text{Det } H^0$$

from  $0 \rightarrow C_{cl}^j \rightarrow C^j \xrightarrow{d} C_{ex}^{j+1} \rightarrow 0$   
 $\Rightarrow \text{Det } C^j \cong \text{Det } C_{cl}^j \otimes (\text{Det } C_{ex}^{j+1})^{-1}$   
 $0 \rightarrow C_{ex}^j \rightarrow C_{cl}^j \rightarrow H^j \rightarrow 0$   
 $\Rightarrow \text{Det } C_{cl}^j \cong \text{Det } C_{ex}^j \otimes \text{Det } H^j$

inner product  $(,)$  induces a ~~canonical~~ preferred element  $\mu \in \text{Det } C^0$

and an element  $\mu_H \in \text{Det } H^0$

- by wedging any o/n basis in  $C^0$

then the torsion is:  $\tau(C^0, d, (,)) = \mathbb{T}(\mu) = \tilde{\tau} \cdot \mu_H \in \text{Det } H^0 / \{\pm 1\}$

Integral formula: view  $C^*$  as a  $\mathbb{Z}$ -graded vector space

$$\mathbb{F} := C^* \mathbb{K} \oplus (C^*)^* [-2] \quad \text{"space of (BV) fields"}$$

	$C^0$	$C^1$	$C^2$	$C^3$	...	$C^n$
$(C^*)^*$	$(C^3)^*$	$(C^2)^*$	$(C^1)^*$	$(C^0)^*$		
$\mathbb{F}^{2n}$	$\mathbb{F}^{-1}$	$\mathbb{F}^0$	$\mathbb{F}^1$	$\mathbb{F}^2$	...	$\mathbb{F}^{n-1}$

fields:  $(A, B) \in \mathbb{F}$   
 $\uparrow$  cochain  $\uparrow$  chain

$A = A^{(0)} + \dots + A^{(n)}$  ← cochain degree  
 $B = B^{(0)} + \dots + B^{(n)}$  ← chain degree

$A^{(j)}$  is (even) for  $j$  (odd)

$B^{(j)}$  is (odd) for  $j$  (even)

$\mathbb{F}$  has a canonical  $-1$ -symplectic structure  $\omega: \mathbb{F}^j \otimes \mathbb{F}^{1-j} \rightarrow \mathbb{R}$

let  $(C^*, d)$  be acyclic,  $K \rightarrow$  a chain contraction

$L = C_{coex}^* [1] \oplus (C_{coex}^*)^* [-2] \subset \mathbb{F}$  a Lagrangian subspace - "gauge-fixing"

$S = \langle B, dA \rangle = \sum_{j=0}^{n-1} \langle B^{(j+1)}, dA^{(j)} \rangle$  - "action functional"

$\text{Dens}^{1/2} \mathbb{F} \cong \text{Dens } \lambda$

Thus:  $\text{Dens}^{1/2} \mathbb{F} \cong \text{Dens } C^* [1] \cong \text{Det } C^0 / \{\pm 1\}$   $\exists \mu$  preferred element

Theorem: for  $C^*$  acyclic, any Lagrangian  $\lambda \subset \mathbb{F}$

$\int_{L \subset \mathbb{F}} e^{\frac{i}{\hbar} S(A,B)} \mu^\dagger = \tau(C^*, d, (,))$

where  $\mu^\dagger = \sum_{\mathbb{Z}} \mu$  and  $\sum_{\mathbb{Z}} = \prod_{j=0}^n \left( \underbrace{\left( \frac{2\pi i \hbar}{\epsilon} \right)^{\frac{1}{4} + \frac{1}{2} j (-1)^j}}_{\sum_{\mathbb{Z}}^j} \left( \frac{2\pi i \hbar}{\epsilon} \right)^{\frac{1}{4} + \frac{1}{2} j (-1)^{j-1}} \right)^{\dim C^j}$



Model Gaussian (Fresnel) integrals:

pair of even variables:  $\int_{V \oplus V^*} da db e^{\frac{i}{\hbar} \langle b, Ma \rangle} = (2\pi\hbar)^{\dim V} (\det M)^{-1}$

$M \in \text{End } V$

pair of odd variables:  $\int_{\Pi V \oplus \Pi V^*} D\alpha D\beta e^{\frac{i}{\hbar} \langle \beta, M\alpha \rangle} = \left(\frac{i}{\hbar}\right)^{\dim V} \cdot \det M$

non-cyclic case:

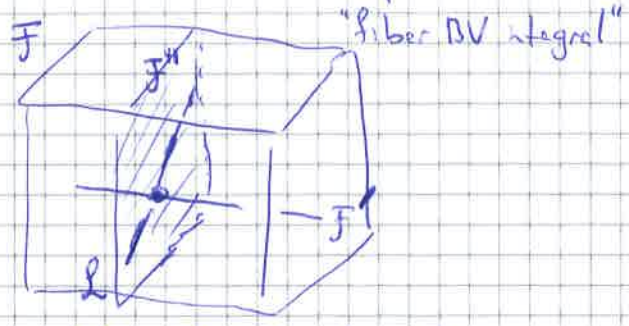
choose an  $(i, p, K)$ -triple  $C^* \xrightarrow{(i, p, K)} H^*$

Kodge decomp.  $C^* = \underbrace{C_{\text{charm}}^*}_{\cong H^*} \oplus C_{\text{ex}}^* \oplus C_{\text{coes}}^*$  induces symplectic splitting

$\mathcal{F} = \mathcal{F}' \oplus \mathcal{F}''$  and a Lagrangian  $\mathcal{L} \subset \mathcal{F}''$   
 $H^* \oplus (H^*)^* [2]$   $\leftarrow$   $\ker p [1] \oplus (\ker p)^* [-2]$   
 "slow fields" "fast fields"

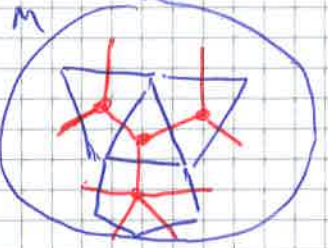
Theorem 1:

$\int_{\mathcal{L} \subset \mathcal{F}''} e^{\frac{i}{\hbar} S(C, B)} \mu^{\hbar} = \sum_{H^*} \tau(C, d, (i)) \in \mathbb{C} \otimes \text{Det } H^* / \{\pm 1\}$   
 with  $\sum_{H^*}^{\hbar} = \prod_{j=0}^n \left(\sum_j^{\hbar}\right)^{\dim H^j} \in \mathbb{C}$



Geometrical setting (cellular complexes)

$M$  - a <sup>closed</sup> manifold,  $X$  - CW decomposition  $X^V$ -dual CW decomposition



$E$  -  $\mathbb{Q}(m)$ -local system on  $M$

$C^0 = C^0(X, E)$  - cellular cochains

$\mathcal{F}_X = C^0(X, E)[1] \oplus C_{-1}(X^V, E^*)[-2]$   $\leftarrow$  model I

or  $C^0(X, E)[1] \oplus C^0(X^V, E^*)[n-2]$   $\leftarrow$  model II

(uses Poincaré duality)

fields

$A = \sum_{e \in CX} e^* (A_e)$   
 $\mathbb{R}^m$ ,  $gh = 1 - |e|$

$B = \sum_{e^V \in CX^V} (e^V)^* (B_{e^V})$   
 $(\mathbb{R}^m)^*$ ,  $gh = n - 2 - |e^V|$

$S_X = \langle B, dA \rangle_X = \sum_{\substack{e, e' \in CX \\ e' \in \partial e}} B_{e^V} A_{e'}$   
 a boundary cell

$\mu_X \in \text{Dens}^{1/2} \mathcal{F} \cong \text{Det } C^0(X, E)$  - standard cellular half-density.



(Torsion)  
←

Partition Function:

from  $C^*(X, E) \xrightarrow{(i, p_0, K)} H^*(M, E)$ ,  
one refers to splitting  $\mathbb{F} = \mathbb{F}' \oplus \mathbb{F}''$  and a Lagrangian  $\mathbb{L} \subset \mathbb{F}''$

$$Z(X, E) = \int_{\mathbb{L} \subset \mathbb{F}''} e^{\frac{i}{\hbar} S_X(A, B)} \mu_X^{\hbar} = \sum_{\mathbb{H}^*} \underbrace{\tau(X, E)}_{\text{Reidemeister torsion}} \in \mathbb{C} \otimes \text{Det } H^*(M, E) / \langle \pm 1 \rangle$$

- does not depend on  $X$ ! (only on  $M$ ).

(Normalized) two-point Correlation Function:

$$\langle\langle A e^i B e^v \rangle\rangle = \frac{1}{Z} \int A e^i B e^v e^{\frac{i}{\hbar} S_X} \mu_X^{\hbar} = K e^i e^v \quad \in E_e^i \otimes E_e^{v*}$$

- matrix element of  $K: C^*(X, E) \rightarrow C^{-1}(X, E)$

- "Green's function"

satisfies a difference equation corresp to  $dK + Kd = id - iop$

Field theory setting (on manifolds)

AS Schwarz '78:  $M$ -closed manifold,  $E$ -acyclic local system

$$\int_{\Omega^1(M, E) \otimes \Omega^{n-2}(M, E^*)} e^{\frac{i}{\hbar} \int_M \langle B, d_A A \rangle} \xrightarrow{\substack{\text{abelian BF theory} \\ \text{DADBS} \rightsquigarrow \text{gauge fixing}}} \int_{(U, N) \in \Omega^1(M, E) [1] \otimes \Omega^1(M, E^*) [n-2]} e^{\frac{i}{\hbar} \int_M \langle B, d_A A \rangle} dU d_B$$

$\rightsquigarrow$   $\infty$ -dimensional Gaussian-integral;  $\xi$ -regularization

$$\rightsquigarrow Z = \prod_{j=0}^n \left( \det_{\Omega^j} \Delta_E \right)^{\frac{-(-1)^j j}{2}} \quad \text{such that } \begin{cases} d_A^* u = 0 \\ d_E^* B = 0 \end{cases} \text{ - "Lorentz gauge"}$$

= Ray-Singer torsion of  $(M, E)$ .

Witten '89:  $M$ -closed oriented 3-manifold,  $G = SU(2)$  (can be generalized to other groups)

fields: connections (in the trivial  $G$ -bundle) on  $M$ ,  $A \in \text{Conn}_M \cong \Omega^1(M, \mathfrak{g})$

action:  $S_{CS}(A) = \int_M \text{tr} \left( \frac{i}{2} A \wedge dA + \frac{i}{3} A \wedge A \wedge A \right)$

$$Z = \int_{\text{Conn}_M} DA e^{\frac{i}{\hbar} S_{CS}(A)}$$

by stationary phase: crit. points of  $S_{CS}$  are flat connections

contribution of an acyclic flat connection  $A_0$ :

$$Z_{A_0}^{\text{flat}} = \int_{\Omega^1(M, \mathfrak{g})} Da e^{\frac{i}{\hbar} S_{CS}(A_0 + a)}$$

$S_{CS}(A_0) + \int_M \text{tr} \left( \frac{i}{2} a \wedge dA_0 + \frac{i}{3} a \wedge a \wedge a \right)$



Witten / Axelrod-Singer:

$$Z_{A_0}^{pert} \cong \int_{\Omega^1(M, \mathfrak{g})} Da e^{\frac{i}{h} S_{CS}(A_0 + a)} \xrightarrow{\text{gauge fixing}} \int D\tilde{a} e^{\frac{i}{h} S_{CS}(A_0 + \tilde{a})}$$

given by  $d_{A_0}^* \tilde{a} = 0$  - Lorentz gauge

$$= e^{\frac{i}{h} S_{CS}(A_0)} \cdot \tau(M, A_0)^{\frac{1}{2}} e^{\frac{\pi i}{2} \text{sign}(A_0, g)} \cdot e^{\frac{i}{24} \frac{dim G}{2\pi} S_{grav}(g, \varphi)} \cdot (1 + \alpha h^2 + \dots)$$

here:  $\tau(M, A_0)$  - Ray-Singer torsion

$$h(A_0, g) = \lim_{s \rightarrow 0} \sum \text{sign}(\lambda) \lambda^{-s}$$

- APS  $h$ -invariant of the operator  $L_- = *d_{A_0} + d_{A_0}^* G \Omega^{odd}(M, \mathfrak{g})$

$S_{grav}$  - Chern-Simons action for Levi-Civita connection

$\varphi$  - framing of  $M$

cf. Fresnel integral  $\int_V e^{\frac{i}{h} \langle a, Pa \rangle} da = (\pi i)^{\frac{dim V}{2}} (\det P)^{-\frac{1}{2}} \cdot e^{\frac{\pi i}{4} \text{sign } P}$

~~result~~  $Z_{A_0}^{pert}$  independent of  $g$  (but depends on  $\varphi$ )

- connections are given by configuration space integrals; integral-product of Green's functions  $\int_{\Omega^2(\text{Conf}_2(M), E \otimes E^*)}$
- integral kernel for  $K = d_{A_0}^* \Delta_{A_0}^{-1}$  - Hodge-de Rham chain contraction

Non-abelian BF theory on a <sup>closed</sup> surface  $\Sigma$  of genus  $g \geq 2$  (Witten, 1991)

$$\int_{\Omega^1(\Sigma, \mathfrak{g}) \otimes \Omega^2(\Sigma, \mathfrak{g}^*)} e^{\frac{i}{h} \int_{\Sigma} \text{tr } B \wedge (A + A \wedge A)} \xrightarrow{\text{gauge fixing, crit. points are } (A=A_0 \text{-flat connections, } B=0)} \int_{\mathcal{M}_{\Sigma}^{irred}} \tau(\Sigma, A_0) =: Z$$

$\mathcal{M}_{\Sigma}^{irred} \leftarrow \begin{matrix} \text{moduli space of (irreducible) flat connections} \\ \subset \text{Hom}(\pi_1(\Sigma), G) / G \end{matrix}$

$$\tau(\Sigma, A_0) \in \text{Det } H_{d_{A_0}}^* = \text{Det} \left( \Lambda^{top} H_{d_{A_0}}^* \right)^* \cong \Lambda^{top} T_{[A_0]}^* \mathcal{M}_{\Sigma}$$

$\Rightarrow \tau(\Sigma, A_0)$  is a volume form on  $\mathcal{M}_{\Sigma}$ . In fact (computer),  $\tau(\Sigma, A_0) = \frac{\omega^N}{N!}$  - Liouville volume form for Atiyah-Bott sym. structure on  $\mathcal{M}_{\Sigma}$ ,  $N = \frac{1}{2} \dim \mathcal{M}_{\Sigma} = (g-1) \dim G$ .

So:  $Z = \int_{\mathcal{M}_{\Sigma}^{irred}} \text{Symplectic volume of } \mathcal{M}_{\Sigma}^{irred} = \#Z(G) (\text{Vol } G)^{2g-2} \sum_{R \text{-irrep of } G} \frac{1}{(\dim R)^{2g-2}}$  for  $G = SU(2)$ ,  $Z \propto B_{2g-2}$  Bernoulli number