

Elimination of the reaction rate “scale effect”: Application of the Lagrangian reactive particle-tracking method to simulate mixing-limited, field-scale biodegradation at the Schoolcraft (MI, USA) site.

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This article has been accepted for publication and undergone full peer review but has not been through the copyediting, typesetting, pagination and proofreading process which may lead to differences between this version and the Version of Record. Please cite this article as an 'Accepted Article', doi: 10.1002/2017WR021103

Abstract

Measured (or empirically fitted) reaction rates at groundwater remediation sites are typically much lower than those found in the same material at the batch- or laboratory-scale. The reduced rates are commonly attributed to poorer mixing at the larger scales. A variety of methods have been proposed to account for this scaling effect in reactive transport. In this study, we use the Lagrangian particle tracking and reaction (PTR) method to simulate a field bioremediation experiment at the Schoolcraft, Michigan site. A denitrifying bacterium, *Pseudomonas Stutzeri* strain KC (KC), was injected to the aquifer, along with sufficient substrate, to degrade the contaminant, Carbon Tetrachloride (CT), under anaerobic conditions. The PTR method simulates chemical reactions through probabilistic rules of particle collisions, interactions, and transformations to address the scale effect (lower apparent reaction rates for each level of upscaling, from batch- to column- to field-scale). In contrast to a prior Eulerian reaction model, the PTR method is able to match the field-scale experiment using the rate coefficients obtained from batch experiments.

1 Introduction

Bioremediation is an important technology to remove contaminant mass, especially organic pollutants, from aquifers. Application of an effective and efficient remediation system depends in large part on prediction of the time scale of contaminant degradation and/or removal. Thus, accurate characterization of multiple reactive transport processes is critical for field-scale bioremediation design [Steefel *et al.*, 2005; Hesse *et al.*, 2009; Scheibe *et al.*, 2009].

Numerous modeling efforts have focused on developing mathematical equations to incorporate chemical reaction kinetics to the transport processes. The most common model is the advection-dispersion equation with the reaction as a source or sink term (ADRE) (e.g., [Hesse *et al.*, 2009; Yabusaki *et al.*, 2011; Porta *et al.*, 2012a; Ding *et al.*, 2013]). However, a variety of studies [Chapelle and Lovley, 1990; Scholl, 2000; Phanikumar *et al.*, 2005; Meile and Tuncay, 2006] indicated that the ADRE models using reaction parameters derived from laboratory experiments overestimated the field-scale reaction rates—sometimes by orders-of-magnitude. One major reason is the “scale effect” for chemical reactions [Lohse *et al.*, 2009; Rubin *et al.*, 2012]; for instance, Rubin *et al.* [2012] suggested three possible scaling reasons that batch parameters may not be applicable to transport problems: 1) different timescales to reach chemical equilibrium; 2) different transfer rates due to the degree of mixing at different scales; and 3) different mass ratios of chemical saturation at different scales. Of these, poorer mixing of reactants induced by the increased heterogeneity of the transport media at larger scales may result in significantly lower reaction rates [Dentz *et al.*, 2011; Bolster *et al.*, 2012].

Because parameters from laboratory-scale experiments have limited applicability to field-scale studies, effective reaction rates are usually used. The

effective reaction rates vary from site to site and may change with time. The estimated parameters are also model dependent and are not directly related to any measurable property of the system [Pedretti *et al.*, 2013]. Because of the lack of model predictive ability, an accurate assessment of field-scale parameters would appear to require field-scale (in space and time) tests, severely limiting the advantage of model simulations.

The limited predictive capability and uncertainty associated with the ADRE model in practice has prompted the development of other models to incorporate the effects of poor mixing. One of these is the Lagrangian particle tracking and reaction (PTR) algorithm, which simulates the reactive transport via Monte Carlo simulation of particle collision and interaction through probabilistic rules [Waite, 1957; Gillespie, 1976; Benson and Meerschaert, 2008; Paster *et al.*, 2014]. Benson and Meerschaert [2008] proposed a PTR method to simulate diffusion-controlled bimolecular reaction under incomplete mixing conditions. Their method showed that self-organized patterns of chemical heterogeneity engendered poor mixing and explained the slowed reaction at late times. The method was extended to moving flows, and the degree of mixing was linked to the number of particles used in a simulation, which represents the non-uniform distribution of initial concentrations (chemical heterogeneity) [Paster *et al.*, 2014]. The PTR method also successfully reproduced the results of two benchmark laboratory-scale column experiments that showed the “scale effect” of poor mixing relative to beaker-scale reactions [Ding *et al.*, 2013].

Moving toward the goal of simulating realistic field-scale experiments, Ding and Benson [2015] extended the PTR method to the Monod-type reactions and applied the method to a column experiment of Carbon Tetrachloride (CT) biodegradation. The authors found that various mechanisms that may contribute to slower biochemical reactions (e.g., crowding, enzyme de-activation) all manifest as diffusion-limited mixing. Therefore, the intricacies of bioremediation can be handled by the PTR method. In this study, we focus on the application of the PTR method to accurately simulate reactive transport associated with a bioremediation experiment at the Schoolcraft site in Michigan, USA. Previous studies (e.g., [Dybas *et al.*, 2002; Phanikumar *et al.*, 2002, 2005]) noted the scale effect when moving from flask- to column- to field-scale biodegradation of CT. Our hypothesis is that the PTR method will explain the scale-dependent chemical reactions through simulation of the poor-mixing effect, rather than by adjusting the empirical kinetic rates as in previous studies. In other words, we will test if the PTR method using batch-scale derived reaction rates is able to simulate the field-scale behavior.

2 Background

In the 1990s, a comprehensive field-scale bioremediation campaign was launched at the Schoolcraft site in Michigan (MI), USA [Dybas *et al.*, 1998; Hyndman *et al.*, 2000; Phanikumar *et al.*, 2005]. Numerous wells were installed, including many with continuous coring, to allow high-resolution mea-

surements of hydraulic conductivities and chemical conditions. Cores were taken across the site to allow detailed characterization of aquifer properties. A line of wells (D1–D15) were installed in a manner that allowed injection and/or withdrawal of bacteria and nutrients to stimulate bioremediation (Figure ??c). Finally, a series of multi-level wells was installed to monitor the progress of the remediation experiment.

2.1 The basis of bioremediation: Laboratory work

Prior to conducting the field-scale bioremediation at the site, laboratory studies [Dybas *et al.*, 1995; Mayotte *et al.*, 1996] revealed that a denitrifying bacterium, *Pseudomonas Stutzeri* strain KC (KC), in the presence of sufficient substrate, can rapidly degrade CT to carbon dioxide, formate, and dechlorinated non-volatile byproducts under anaerobic conditions without producing chloroform, a more persistent contaminant. With this finding, biodegradation of CT by KC was tested in the laboratory and field. Tests included batch (flask) experiments [Criddle *et al.*, 1990; Tatara *et al.*, 1993; Dybas *et al.*, 1995], column experiments [Witt *et al.*, 1999; Phanikumar *et al.*, 2002], and pilot field studies [Dybas *et al.*, 1998]. One level of upscaling was achieved in the laboratory when a no-flow column experiment was conducted by Witt *et al.* [1999]. In this experiment, a 100 cm-long column was filled with sediments and groundwater extracted from site borings. The groundwater was supplemented with initial concentrations of CT and nitrate at 0.1 and 25 milligrams per liter (mg/L), respectively. The column was inoculated with KC, acetate, and base (to mediate the pH) at the center of the column (between 44.4 and 59.6 cm) and was maintained as a static incubation to understand the role of chemotaxis. The inoculation had KC at $1.2 \pm 0.1 \times 10^8$ colony-forming units per milliliter (CFU/mL) and an acetate concentration of 1,533 mg/L. One CFU/mL is approximately equal to 1.67×10^{-7} ppm for strain KC [Phanikumar *et al.*, 2002]. The column had 10 sampling ports spaced at 7.6-cm intervals to monitor the concentrations of dissolved species and biomass. Over the course of a month, a significant fraction of CT was degraded, demonstrating the viability of the technology.

2.2 Site Information

The unconfined aquifer at the Schoolcraft site is composed of glaciofluvial sediments overlying a thick clay unit, which acts as an aquitard [Kehew *et al.*, 1996; Phanikumar *et al.*, 2005]. The top of the aquitard was found at approximately 27.3 meters (m) below ground surface (bgs), while the water table was around 4.5 m bgs [Hyndman *et al.*, 2000]. The natural hydraulic gradient at the site was roughly 0.001, with a general groundwater flow direction from northwest to southeast (Figure ??).

As part of the installation of the bioremediation delivery and monitoring wells, 346 soil core samples were taken from 11 borings, repacked and placed in constant-head permeameters. The repacked samples were shown to provide reasonable estimates of the horizontal hydraulic conductivity (K) values according to a model verification of tracer tests against observed concentra-

tion profiles [Biteman *et al.*, 2004]. The K analysis and core logging revealed 4 stratigraphic zones, with mean $\ln(K)$ (cm/s) of -1.26 , -1.81 , -1.49 , and -1.86 , from deepest to shallowest. In general, the highest K zone exists at the bottom of the aquifer. The large number of samples allowed an estimate of the anisotropic variograms in each zone. In general, the variogram ranges in the horizontal directions were estimated to be from 3 to 18 m, and vertical ranges were from 0.35 to 1.62 m. The overall variance of $\ln(K)$ is 0.634. Flow and bromide tracer transport modeling (discussed in more detail below) showed that the K -field generated from zonal (non-stationary) kriging was superior to non-zonal kriging [Biteman *et al.*, 2004]. We will use this K field (Figure ??), as did Phanikumar *et al.* [2005], to simulate flow conditions during the bromide tracer test and the bioremediation experiment.

There were several contaminant plumes reported in the aquifer [Hyndman *et al.*, 2000; Dybas *et al.*, 2002]. The field remediation experiment was conducted within a plume, designated Plume A, which was contaminated with carbon tetrachloride (CT) [Hyndman *et al.*, 2000; Phanikumar *et al.*, 2005]. The CT contamination within plume A was 1,600 m long and 160 m wide [Phanikumar *et al.*, 2005]. Concentrations from 221 locations indicated that higher CT concentrations were in the deeper, high-conductivity part of the aquifer, as illustrated in Figure ??.

2.3 Bioremediation Method

The field remediation system at the Schoolcraft site was designed to inoculate non-native microbes and recirculate the groundwater through a series of injection and extraction wells aligned perpendicular to the natural gradient flow (Figure ??). These pumping wells were screened from 9.1 to 24.4 m bgs using 0.025 cm slotted screen [Hyndman *et al.*, 2000]. A total of 134 piezometers, each with 0.33 m-long screens across the vertical extent of the plume, composed the monitoring array to record the concentrations. Prior to the bioremediation, a bromide tracer test was conducted under the approximate cyclic injection/withdrawal cycle for 20 days to assess transport rates within the contaminated heterogeneous aquifer unit [Phanikumar *et al.*, 2005].

To initiate the bioremediation process, a single inoculation was conducted using 18,900 L of KC-laden groundwater through the fifteen (15) delivery wells, which were 1 m apart. The locations of these wells (names start with D) are shown in Figure ??.

Groundwater was recirculated for 6 hours every week through pumping and injection. The recirculation consisted of: 1) extracting from every other well (e.g., even numbered wells: D02, D04, . . . , D14) and re-injecting into intervening wells (e.g., odd numbered wells: D01, D03, . . . , D15) after addition of constituents (acetate, bromide, pH amendment, etc.) for 5 hours; 2) reversing the pumping/injection (e.g., pumping from odd numbered wells and injecting back to the even numbered wells) for 1 hour; and 3) keeping natural flow condition for the rest of the week. The extraction/injection orders (even or odd numbered) on wells in the first two stages were switched in the following weeks. The details are described by

Phanikumar *et al.* [2005]. The circulation and monitoring were conducted for 165 days.

3 Methods and Models

3.1 The Simple Form of Enzymatic Reaction

Biodegradation occurs as microorganisms metabolize accessible nutrients (substrates) to grow. The substrates, including organic contaminants, are degraded to inorganic compounds or smaller molecules by biomass [Alexander, 1999; King *et al.*, 2010]. A simple biodegradation (1) following this mechanism under certain conditions can be characterized by the Monod equation [Monod, 1949]:



where k_f , k_r , and k_c are forward, reverse and conversion (transform) rate constants. The substrate S and the biomass or enzyme E form the intermediate enzyme/substrate complex ES through the initial bimolecular reaction with a rate constant k_f [$M^{-1}T^{-1}$]. The ES complex can dissociate to E and S , with a rate constant k_r [T^{-1}], or proceed to form the product P , with a rate constant k_c [T^{-1}].

Under perfectly-mixed conditions, the rates of concentration change are quantified through the law of mass action:

$$d[S]/dt = -k_f[E][S] + k_r[ES] \quad (2a)$$

$$d[E]/dt = -k_f[E][S] + k_r[ES] + k_c[ES] \quad (2b)$$

$$d[ES]/dt = k_f[E][S] - k_r[ES] - k_c[ES] \quad (2c)$$

$$d[P]/dt = k_c[ES] \quad (2d)$$

Michaelis and Menten [1913] originally proposed a simplified solution of (2) by assuming that 1) only a vanishingly small fraction of substrate is bound by enzyme, 2) the complex is very labile and decays to free enzyme, 3) the substrate is in instantaneous chemical equilibrium with the complex, and 4) the conversion rate is directly proportional to the concentration of enzyme. Under these conditions, Eqs. (2) reduce to

$$\frac{d[P]}{dt} = v_{max} \frac{[S]}{K_S + [S]} = k_c[E]_0 \frac{[S]}{K_S + [S]}, \quad (3)$$

where the conversion rate $v_{max} \equiv k_c[E]_0$, $[E]_0$ is the initial enzyme concentration, and K_S is the half saturation coefficient, or Michaelis constant, defined by $(k_r + k_c)/k_f$.

3.2 ADRE-Based Model

Employing the Monod/Michaelis-Menten (hereafter called M-M) kinetics, Phanikumar *et al.* [2005] developed a reactive transport model (Eqs. 4) specifically for CT bioremediation to account for microbial-mediated reac-

tions, advection, dispersion, attachment, and detachment of reactants:

$$\frac{\partial E}{\partial t} = \mathcal{L}_E(E) + \left[\mu_{max} \frac{S}{K_S + S} \frac{A}{K_A + A} - k_{decay} \left(1 - \frac{A}{K_A + A} \right) - k_{att} \right] E + k_{det} \left(1 - \frac{A}{K_A + A} \right) X + Q^s E^s \quad (4a)$$

$$\frac{\partial X}{\partial t} = \left[\mu_{max} \frac{S}{K_S + S} \frac{A}{K_A + A} - (k_{decay} + k_{det}) \left(1 - \frac{A}{K_A + A} \right) \right] X + k_{att} E \quad (4b)$$

$$\frac{\partial S}{\partial t} = \mathcal{L}_S(S) - \left(\frac{\mu_{max}}{Y_n} \frac{S}{K_S + S} \frac{A}{K_A + A} + \gamma \frac{S}{K_S + S} \right) (E + X) - \frac{k_{decay}}{Y_{nd}} \left(1 - \frac{A}{K_A + A} \right) (E + X) + Q^s S^s \quad (4c)$$

$$\frac{\partial A}{\partial t} = \mathcal{L}_A(A) - \frac{\mu_{max}}{Y_a} \frac{S}{K_S + S} \frac{A}{K_A + A} (E + X) + Q^s A^s \quad (4d)$$

$$\left(1 + \frac{\rho f K_d}{\theta} \right) \frac{\partial c}{\partial t} = \mathcal{L}_c(c) - k' c \left(E + X \frac{\rho f K_d}{\theta} \right) - \frac{\rho k_{des} c}{\theta} [(1 - f) K_d c - c_S] + Q^s c^s \quad (4e)$$

$$\frac{\partial c_S}{\partial t} = k_{des} [(1 - f) c K_d - c_S] - k' c_S X \quad (4f)$$

where we dropped the square brackets when denoting concentration, E is the concentration of mobile bacteria; X is the amount of bacteria attached to solids; S is the substrate, nitrate; A is the concentration of acetate; c is the concentration of CT, and c_S is the concentration of CT adsorbed to the solids. The concentrations have units of mg/L , including the mobile and immobile bacteria, which have the units converted from CFU/mL [Phanikumar *et al.*, 2005]. For each mobile species, there is a linear advection/dispersion operator $\mathcal{L}(f) = -\nabla \cdot (\mathbf{v}f - \mathbf{D}\nabla f)$ that includes the effects of spatio-temporally variable velocity \mathbf{v} and species-dependent diffusion/dispersion tensor \mathbf{D} . Q^s is the flow of source/sink term, and the s superscript denotes the concentration of each constituent in the source/sink term. K_S and K_A are half saturation constants for nitrate and acetate, respectively, μ_{max} is the maximum conversion rate, k_{decay} is biomass decay rate, k_{att} is the attachment coefficient of biomass, k_{det} is the detachment coefficient of biomass, k' is the degradation rate for CT, and k_{des} is the desorption rate of CT. Y_n , Y_a , and Y_{nd} are the cell yields for nitrate, acetate, and biomass consuming nitrate, respectively. The factor f is the fraction of exchange sites at equilibrium, K_d is the CT distribution coefficient, ρ is the bulk density of soil, and γ is the nitrate utilization rate by indigenous microflora or endogenous respiration. The population of indigenous microflora is assumed proportional to the KC bacteria and its reactions have the same form as those of KC [Phanikumar *et al.*, 2002].

In this model, a correction factor $[1 - A/(K_A + A)]$ was added to the bacteria decay term to account for the increase of decay rate at low nutrient concentration [Beefink *et al.*, 1990; Phanikumar *et al.*, 2005]. However, recent models (i.e., [Tan *et al.*, 1994; Tufenkji, 2007; Ding, 2010]) assumed that the decay rate is independent of the concentration of nutrient. Thus, we mod-

ified the models as Eqs. 5 by ignoring the acetate concentration-dependency terms. A comparison of RT3D simulations using Eqs. 4 and Eqs. 5, provided in Appendix B: , indicates that the difference of model results is negligible.

$$\frac{\partial E}{\partial t} = \mathcal{L}_E(E) + \left(\mu_{max} \frac{S}{K_S + S} - k_{decay} - k_{att} \right) E + k_{det} X + Q^s E^s \quad (5a)$$

$$\frac{\partial X}{\partial t} = \left(\mu_{max} \frac{S}{K_S + S} - k_{decay} - k_{det} \right) X + k_{att} E \quad (5b)$$

$$\frac{\partial S}{\partial t} = \mathcal{L}_S(S) - \left[\left(\frac{\mu_{max}}{Y_n} + \gamma \right) \frac{S}{K_S + S} + \frac{k_{decay}}{Y_{nd}} \right] (E + X) + Q^s S^s \quad (5c)$$

$$R \frac{\partial c}{\partial t} = \mathcal{L}_c(c) - k'(E + X)c + Q^s c^s \quad (5d)$$

where $R = 1 + \rho f K_d / \theta$ is a modified retardation factor for CT.

This ADRE-type model was applied to simulate a series of tests of CT biodegradation, from column-scale experiments [Witt *et al.*, 1999; Phanikumar *et al.*, 2002] to field-scale pilot studies [Dybas *et al.*, 1998; Phanikumar *et al.*, 2005]. This model is also used in this study for comparison with simulations using the PTR method, which simulates the reactions as a series of elementary steps.

3.3 Particle Tracking Method

The PTR method used here simulates chemical reactions through probabilistic rules of particle collisions, interactions, and transformations. For a bimolecular reaction, the potential reaction between any two particles is based on an explicit calculation of co-location probability multiplied by independent thermodynamic probability that two particles react upon co-location [Benson and Meerschaert, 2008; Ding *et al.*, 2013; Paster *et al.*, 2014].

Using the PTR method, the biodegradation or enzymatic reactions (1) can be simulated as a series of chemical reactions or elementary steps (2). The initial bimolecular reaction that transforms the substrate to the enzyme-complex (i.e., the first part of the reaction: $S + E \rightarrow ES$) is characterized by a second-order kinetics: $d[S]/dt = -k_f[E][S]$. Assuming each E and S particle carries the same amount of mass $m_p = \Omega[S]_0/N_S(t=0)$, where Ω [L^d] is the domain size in d -dimensions, $[S]_0$ is the average initial concentration of S [M], and $N_S(t=0)$ is the initial number of S particles, the probability comprises a co-location density function $v(s)$ and the thermodynamic probability function [Benson and Meerschaert, 2008]:

$$P(react) = k_f \Delta t m_p v(s) \quad (6)$$

where Δt is the numerical time step size and s is the separation of any pair of S and E particles.

The co-location probability density function is the convolution of the individual motion densities of two reactant particles (S and E) over a short time period: $v(s) = \int f_S(x) f_E(s+x) dx$, where $f_S(x)$ and $f_E(x)$ denote the motion densities of S and E particles away from their current positions

through diffusion. Each is a Gaussian density if particles diffuse under Brownian motion (see details in *Benson and Meerschaert* [2008]; *Benson et al.* [2013]). The reaction probability $P(\text{react})$ is compared with a random number that is uniformly distributed between 0 and 1. If the probability of the reaction is larger than the random number, the two particles are converted to an intermediate ES complex particle. This reaction calculation requires that $k_f \Delta t m_p v(s = 0) < 1$ [*Benson and Meerschaert*, 2008]. Other forms of bimolecular reaction, such as $A + B \rightarrow 0$ and $A + B \rightarrow C + D$, can be simulated similarly.

For the monomolecular reactions with first-order kinetics of the general form $dC/dt = -kC$, including the reverse dissociation reaction ($ES \rightarrow E + S$) and transform reaction ($ES \rightarrow P$), the density of particles N represents the local concentration C , thus the reactions can be expressed as $dN/dt = -kN$. For a small time step, Δt , the fraction change of numbers of particles is $\Delta N/N = -k\Delta t$. If the particle transitions are independent of each other, the left hand side is the probability that any particle will transform. In any time step, each particle is chosen and if $k\Delta t$ is greater than a uniform random variable $[0, 1]$, the particle is converted. This first-order kinetics simulation requires that $k\Delta t < 0.1$ for suitable accuracy.

The series of reactions (2a)-(2d), which characterize the M-M type of reaction that bacteria consume substrate and nutrients, are simulated as follows. For every time step, each E particle is selected sequentially to find nearby S particles, and the probability of co-location for each pair of S and E particles is calculated. If one reaction occurs, an intermediate particle ES is placed randomly between the pair of reactant particles, which are removed. The intermediate particle ES either transforms to a product particle, or reverses to the initial S and E particles, or stays intact. These three processes are independent and are characterized by first-order kinetics. One random number is generated to check the probability for each of these reaction processes at every time step. The impact on the reaction from the locations of released S and E particles was found to be minor [*Ding and Benson*, 2015]. Thus, we assume here that the released reactant particles are randomly distributed around an intermediate ES particle within a diffusion distance $\sqrt{2D\Delta t}$.

3.4 Particle Transport Model

Our goal is to assess the differences in the transport and reaction algorithms, not to re-create the underlying hydraulics at the site. To that end, we use the exact 3-dimensional velocity fields that were generated (using MODFLOW) in the initial study [*Phanikumar et al.*, 2005]. Between each reaction step, each particle is moved based on its specific location and the flow field around it using the numerical random walk particle tracking code RW3D [*Fernández-García et al.*, 2005].

RW3D simulates solute transport by partitioning the solute mass into a large number of representative particles. The evolution of a particle's location is driven by a drift term that includes the advective movement, and a

superposed Brownian motion responsible for dispersion. The displacement of a particle is modified from the Itô-Taylor integration scheme by substituting the drift vector with a modified velocity vector that includes the effects of a gradient of the dispersion tensor components [Salamon *et al.*, 2006]:

$$\mathbf{X}_p(t + \Delta t) = \mathbf{X}_p(t) + \Delta t[\mathbf{v}(\mathbf{X}_p(t)) + \nabla \cdot \mathbf{D}(\mathbf{X}_p(t))] + \sqrt{2\mathbf{D}(\mathbf{X}_p(t))\Delta t} \cdot \xi, \quad (7)$$

where Δt is the time step, $\mathbf{X}_p(t)$ is the position of a particle at time t , \mathbf{v} is the velocity vector, \mathbf{D} is the dispersion coefficient tensor made diagonal in the direction of transport, and ξ is a vector of independent standard normal random variables. The random walk code uses a hybrid scheme for the velocity interpolation that provides divergence-free velocity fields and a continuous dispersion tensor field that enforces mass balance at grid interfaces of adjacent cells with any degree of hydraulic conductivity contrast [Salamon *et al.*, 2006].

3.5 Schematic of Modeling Procedure

A schematic of calculation algorithm of the PTR simulations, with the developed particle tracking algorithm of reactions incorporated into the flow code, is shown in Figure ???. At any time step, the simulation follows the model procedures: i) the CT and bacteria KC experience attachment and detachment processes, which are assumed to follow a linear isotherm (see e.g. [Benson and Bolster, 2016]). These processes follow the first-order kinetics; ii) The biomass particles are looped over to find all potential nitrate particles that may bind together into the intermediate complex in the presence of sufficient acetate as an electron donor, acetate, as described in Section 3.3; iii) The complex either transforms to the product, reverses back to the reactants, or stays intact as the complex. If the intermediate complex particle transforms to a product, the bacteria particle is released; at the same time, the bacteria grow by randomly adding bacteria particles according to their growth yield. If the reverse reaction occurs, a substrate and a microbe particle are regenerated. iv) Concurrently with reactions between biomass and substrate (ii and iii), the degradation of CT by bacteria is simulated as a bimolecular reaction. v) The biomass also experiences decay, which is simulated first-order kinetics related to mass/concentration of bacteria. vi) The mobile particles move via random walks after the elementary steps to the next time step.

The elementary steps and model parameters, as quantified in Eqs. 5, are listed in Table 1. Reactions also occur between mobile and immobile particles (including the attached KC and adsorbed CT), similar to the steps shown in Figure ??. However, the probability function, particularly the co-location density, for the bimolecular reaction is modified to account for the immobility of attached particles, as described by Ding and Benson [2015].

4 Results and Discussion

4.1 Kinetic Parameters

The PTR model uses chemical reaction parameters from batch experiments directly in the field-scale simulation. As introduced in Section 2.1, a

Table 1. Elementary steps of the reactions

Step	Equation	Parameter
Adsorption of CT	$c \rightarrow c_S$	k_{ads}^1
Desorption of CT	$c_S \rightarrow c$	k_{des}
Attachment of mobile bacteria KC ²	$E \rightarrow X$	k_{att}
Detachment of immobile bacteria KC ²	$X \rightarrow E$	k_{det}
Nitrate binds to mobile bacteria	$S + E \rightarrow ES$	k_s^3
Nitrate binds to immobile bacteria	$S + X \rightarrow XS$	k_s^3
CT biodegradation by mobile bacteria	$c + E \rightarrow P^4$	k'
CT biodegradation by immobile bacteria	$c + X \rightarrow P^4$	k'
Bacteria biodegrade adsorbed CT	$c_S + E \rightarrow P^4$	k'
Intermediate ES reverts to nitrate and mobile KC	$ES \rightarrow S + E$	k_r
Intermediate XS reverts to nitrate and immobile KC	$XS \rightarrow S + X$	k_r
Transformation of ES and growth of mobile KC	$ES \rightarrow (1 + Y)E + p^5$	k_c
Transformation of XS and growth of immobile KC	$XS \rightarrow (1 + Y)X + p^5$	k_c
Decay of mobile bacteria KC	$E \rightarrow 0$	k_{dec}
Decay of immobile bacteria KC	$X \rightarrow 0$	k_{dec}

¹ The rate is calculated based on the fraction of exchange sites and distribution coefficient. Linear isotherm is assumed.

² Indigenous microflora are assumed to have the same steps as KC.

³ The reaction rates involving indigenous microflora are calculated based on the ratio of γ and μ_{max} in Eq. (5).

⁴ P represents the product of CT biodegradation.

⁵ Y is the growth yield of biomass, and p is the products of nitrate transformation.

series of batch experiments under different conditions were conducted to estimate the reaction rates prior to the column- and field-scale studies. The batch parameters used for the simulation are tabulated in Table 2. For instance, in evaluating the role of trace metals on CT degradation rate, *Tatara et al.* [1993] found that the second-order rate coefficient decreased as culture age increased from 48 to 72 hours, which were the times for the culture to grow prior to the inoculation [Dybas *et al.*, 1995; Del C. Sepulveda-Torres *et al.*, 1999]. *Phanikumar et al.* [2002, 2005] reported the reaction rate as $2.70 \text{ L mg}^{-1} \text{ day}^{-1}$ by taking the reaction rate for cultures aged 72 hours and grown under iron-limiting conditions without the precipitate in [Tatara *et al.*, 1993]. The attachment and detachment coefficients were derived from a column experiment on the transport of KC. In addition, the microbial decay rate, which is the only parameter not measured directly, was from literature, however, the value was shown to be applicable in the simulation of CT biodegradation [Phanikumar *et al.*, 2002].

4.2 Simulation of the Column Experiment in Witt *et al.* [1999]

The capability of the PTR method for biodegradation reactions was tested first on a column-scale experiment. We incorporated the PTR simulation of reactions into the RW3D code [Salamon *et al.*, 2006] to simulate the no-flow column experiment conducted by Witt *et al.* [1999], as introduced in Section

2.1, for verification. The simulation used the procedures introduced in Section 3.5. The mobile particles move purely by diffusion via random walks with different diffusion coefficients for the solutes and biomass. The diffusion coefficients for solutes (e.g., nitrate) and the bacteria, estimated from the concentration profiles and the detection time for each component to reach different ports, were 6×10^{-6} and 7.8×10^{-5} m²/hour, respectively [Ding and Benson, 2015]. Possibly due to the heterogeneities, whether physical, biological, or chemical, the measured initial concentrations (ICs) at different sampling ports at day 0 were not uniform through the column length [Witt *et al.*, 1999]. To represent the non-uniform initial condition, the particles were assigned individually in the 12 sections, which are separated by the 10 ports, based on the concentrations measured at adjacent ports. Figures ??a and ??b show the simulations of PTR model for CT and nitrate, respectively, at days 2 and 26. This heterogeneous IC is reflected in the asymmetric concentrations at later time. We ran 150 simulations and obtained the smoothed concentration profiles by simple binning of particle numbers to account for the stochastic nature of the simulations. The plots reflect the ensemble mean values plus or minus one standard deviation. With the total domain initial number of particles of 3,300, 2,640, and 13 (proportional to the initial concentrations) assigned to nitrate, biomass, and CT, the PTR model in RW3D when populated with the batch rate parameters showed good matches of measured concentrations in the column experiment. Due to the lack of data on the covariance structure of initial concentration perturbation, we tested the simulations via trial-and-error to achieve a reasonable match of the observations to obtain the numbers of particles. The simulation was consistent with [Ding and Benson, 2015], in which the same column experiment was simulated with the PTR model in a Matlab code.

In contrast, the ADRE type of model (Eq. (5)) needed to adjust the effective kinetic parameters to match the column measurements [Ding and Benson, 2015]. In particular, because of incomplete mixing and lower apparent transport rates, the fitted CT reaction rate k' was reduced more than an order-of-magnitude, from 2.70 to 0.189 L mg⁻¹day⁻¹. Additionally, the decay rate of microbes was increased from 0.10 to 0.221 day⁻¹, and the detachment coefficient was changed from 0.018 to 0.043 day⁻¹. The adjustment of effective parameters is consistent with the simulation of a later column experiment under flowing conditions [Phanikumar *et al.*, 2002]. Phanikumar and Hyn-dman [2003] further improved the model simulation by including a degradation term in the sorbed phase and solved with a code based on a fourth-order accurate compact Hermitian scheme, the degradation rate of CT was lowered further to 0.121 L mg⁻¹day⁻¹. The rate was later used for the field-scale bioremediation simulation [Phanikumar *et al.*, 2005]. The comparison of kinetic parameters values used in the various previous grid-based codes is listed in Table 2.

Table 2. Laboratory Measured and ADRE Model Fitted Parameters for CT Biodegradation.

Parameter	Symbol	Units	Batch value	Column-ADRE ^a	Field-ADRE ^b	Field-PTR ^c
Biodegradation rate	k'	L mg ⁻¹ d ⁻¹	2.70	0.189	0.121	2.70
Maximum specific growth rate	μ_{max}	d ⁻¹	3.11	3.11	3.11	3.11
Nitrate utilization by microflora	γ	d ⁻¹	0.0	18.89	18.89	0.0
Microbial decay rate	k_{decay}	d ⁻¹	0.1 ^d	0.221	0.13 ^e	0.1
Attachment coefficient	k_{att}	d ⁻¹	0.9 ^f	0.9	0.9 / 9 ^g	0.9
Detachment coefficient	k_{det}	d ⁻¹	0.018 ^f	0.018	0.04	0.018
Growth yield for nitrate	Y_n	–	0.25	0.25	0.25	0.25
Growth yield for biomass	Y_{nd}	–	0.46	0.46	0.46	0.46
Half saturation coefficient of nitrate	K_m	mg/L	12.0	12.0	12.0	12.0
Binding rate constant	k_f	L mg ⁻¹ d ⁻¹	0.36 ^h	–	–	0.36
Longitudinal dispersivity	D	m	–	–	0.01 ⁱ	0.03 ⁱ
Diffusion coefficient of solute	Dm_s	m ² /hour	–	6×10^{-6} ^j	–	6×10^{-6}
Diffusion coefficient of KC	Dm_E	m ² /hour	–	7.8×10^{-5} ^j	–	7.8×10^{-5}

^a Fitted from the simulation of the no-flow column experiment in [Ding and Benson, 2015].

^b Effective parameters used in the field-scale ADRE-based RT3D model [Phanikumar et al., 2005].

^c Values used in the current PTR simulation of the field-scale bioremediation.

^d The value is from literature, as noted in [Phanikumar et al., 2002].

^e The decay rate was 0.00016 after conversion from Eq. 4 to first-order rate in Eq. 5 by multiplying the acetate correction factor.

^f The attachment and detachment rates were estimated from a column experiment of KC transport.

^g 10 times higher attachment coefficient was used during the inoculation period ([Phanikumar et al., 2005]).

^h This rate, calculated from [Tatara et al., 1993], is used only in elementary reaction steps but not in the Monod equation.

ⁱ The dispersivity was estimated from tracer test prior to the bioremediation using PTR method and ADRE method, respectively.

^j The diffusion coefficients were estimated from the concentration profiles and the detection time for each component to reach different ports in [Witt et al., 1999].

4.3 Simulation of the Field-Scale Non-Reactive Tracer Test

As shown in section 4.2 and in [Ding and Benson, 2015], the PTR method is able to simulate the relatively small degree of upscaling from batch to column scales without adjusting reaction rates. The reduced degree of mixing was achieved by fitting the particle numbers. The particle numbers code the “smoothness” of the initial concentrations and are determined by the concentration autocovariance function(s), if they are estimated with sufficient accuracy [Paster *et al.*, 2014]. This data was not available in the column studies, but is available for the field site, so we hypothesize that the particle method can accurately simulate the field experiment without adjusting any rate parameters from their thermodynamics, batch-scale values, as long as the velocities are well represented by the particles. This hypothesis follows from an analysis of the subgrid velocity and concentration fluctuation terms in the ADRE that need to be accounted for to numerically track imperfect mixing (Appendix D:).

Prior to bioremediation at the Schoolcraft site, a non-reactive tracer test using bromide was conducted for 20 days [Phanikumar *et al.*, 2005]. For the first five hours, groundwater was pumped out of the odd numbered wells (D01, D03, . . . , D15) at a total rate of approximately 9.085 m³/hr. The extracted water, with the addition of Br^- at different concentrations (from 14 to 18 mg/L), was injected into the even numbered wells (D02, D04, . . . , D14); see well locations in Figure ???. Then approximately 9.085 m³ groundwater was pumped out of the even-numbered wells for one hour and injected back into the odd-numbered wells after Br^- was added at the concentration of 23.5 mg/L. After the pumping-injection cycle, the natural flow condition was maintained until day 20. The breakthrough curves of Br^- were recorded at five monitoring wells (9, 10, 11, 12, and 13, as shown in Figure ??) each with five slotted intervals of 0.609 m at depths of 10.7 m, 13.7 m, 16.8 m, 19.84 m, and 22.9 m bgs, respectively [Hyndman *et al.*, 2000]. These depths correspond to approximately 35, 45, 55, 65, and 75 feet bgs, which was how the five intervals were named (e.g., Fig. 1).

Phanikumar *et al.* [2005] used MODFLOW on the grid shown in Figure ??b to calculate heads and discharges. On the same grid, they applied the RT3D model, a mixed Lagrangian and Eulerian finite-difference (FD) implementation of the ADRE, to simulate the transport of the tracer. The advection was (mostly) performed by particles in the hybrid method of characteristics (HMOC), but the dispersion and reaction operations were performed by averaging particle concentrations back to a grid for standard FD calculations. Through calibration, they found that the RT3D model with a longitudinal dispersivity value of 0.01 m and an effective porosity of 0.3 matched the field measurements. The relatively small dispersion coefficient implied that the variations of velocity were captured with the heterogeneous and nonstationary kriged hydraulic conductivity field. Moreover, the relatively rapid breakthrough of tracer (and higher mass recovery) in the deeper region, and slow and low concentration breakthrough in the shallow region, reflected the different hydraulic conductivity zones.

Using the exact same velocities from the MODFLOW model, we simulated the bromide tracer transport using the RW3D model. The re-circulation process was simulated by extracting particles within a radius of 0.1 m of pumping wells and transferring them to the injection wells. The injected particles were distributed randomly within the screened interval of injection wells with probability based on the flux rates at different depths. The PTR method simulated the injection, re-circulation, and transport of 94,100 particles representing the total mass of 94.1 grams of Br^- in the system. We chose the number for the balance of simulation variations and the computation time for a single run, because the numbers of particles do not affect the average of simulated results in the conservative tracer simulation. A small number of particles would lead to a high variation of the simulations, but less computation time for each run. Through model tests, the number of particles used (94,100) based on the assumption that each particle carries 1 mg mass was sufficient to obtain a smooth curve of simulation.

The mean breakthrough curves (normalized by a concentration of 30 mg/L) from an ensemble of 50 PT (RW3D) simulations match somewhat better than those of RT3D model (Fig. 1). In particular, the total mass recovery is better for the PT method in 16 of the 25 observation locations, and the RMSE is lower in 24 of the 25 locations (Fig. 1). The means of the ensemble of PT models are used in the comparison. Mass recovery is calculated using the Matlab function *trapz*, which calculates the area under a set of concentration data by breaking the region into trapezoids. The RMSE is the square root of the sum of square differences between simulations and measurements. When these values are not coincident in time, the simulation values are interpolated to the measurement times using Matlab *interp1* function. It is important to stress that we seek to compare RW3D and RT3D when reactions are included, so that we have not tried to make the new model fit the Br^- data any better. The better fits by RW3D are simply a result of zero numerical dispersion — this feature tends to keep the Br^- more separated in layers than the FD model was able to simulate for the chosen grid resolution.

Through a limited trial-and-error effort, we found that RW3D performed well enough with a longitudinal dispersivity value of 0.03 m, which is larger than that of RT3D model (0.01 m). The parameters for the RT3D model were taken directly from *Phanikumar et al.* [2005]. The difference is due to either numerical dispersion generated from discretization in the FD scheme and/or recirculation well concentration calculation methods. Regarding the first point, finer mesh or sub-scale grid models (e.g., regriding the RT3D model) might allow the dispersivities to match, but that effort is irrelevant to this study. Regarding the second point, the RT3D concentrations of groundwater pumped out of wells were weighted by the hydraulic conductivity of model cells that the pumping wells penetrate, rather than transmissivity, which overestimated the contribution from the layers with small thicknesses and underestimated the contribution from layers with large thicknesses. However, the two models match the measurements from the complex tracer test remarkably well, so that the RW3D model can be applied to the bioremediation ex-

Figure 1. Measured and simulated breakthrough curves of Bromide. The symbols are measured concentrations; the blue lines are simulations using RT3D, and the green lines with error bars are simulated results from RW3D model. The error bars are plus/minus one standard deviation from 50 realizations. Bold, dashed numbers denote well and sampling depth (ft bgs). The numbers on the left are the RMSE from RT3D model and RW3D model compared to the measurements, respectively. The right side labels are the mass recovery from measurements, RT3D, and RW3D, respectively. The sub-figures in which the PT simulations match better than, or equal to, those of ADRE-type model are highlighted in yellow.

periment to assess the effect of maintaining batch-scale reaction rates in the field-scale model.

4.4 Simulation of the Field-Scale Bioremediation

Our goal here is to compare the Eulerian grid-based (RT3D) and PTR (RW3D) methods, so we duplicate as closely as possible the modeling efforts of *Phanikumar et al.* [2005]. We incorporated the reactions listed in Eqs. (5) into the RW3D code to simulate the field-scale CT biodegradation. The initial condition and boundary conditions were consistent with those in the RT3D model [*Phanikumar et al.*, 2005]. As described in Section 2.2, the aquifer had a plume of CT at concentrations from 1.23 to 42.9 $\mu\text{g/L}$ and nitrate concentrations from 21.62 mg/L to 44.25 mg/L from 10.6 m bgs to the top of the aquitard (27.3 m bgs). The RW3D model simulates the transport of CT and nitrate without any reaction for the first 67 days. At this point, the inoculation medium (with KC and acetate) was added, the bacteria then consume nitrate and acetate to grow and biodegrade CT. Throughout the bioremediation, the pumping-injection recirculation scheme was conducted as described in Section 2.3 and 4.3 (see details in [*Phanikumar et al.*, 2005]).

Regarding the initial conditions for the particle simulations, *Paster et al.* [2014] showed that the number of particles is directly related to the “smoothness” of the initial concentrations, as given by the autocovariance functions of the concentration fluctuations. In other words, the particles represent concentration fluctuations as well as the mean, so the number is important for accurate reactant interaction probabilities (see Appendix A:). They equated the effective correlation function for the Dirac-delta particles and the covariance function of measured concentration data C to find that the particle density (in d -dimensions) should follow $\rho \approx \bar{C}_0^2 / (\sigma_C^2 l^d)$, where \bar{C}_0 is the mean concentration, σ_C^2 is concentration variance, and l^d is the autocorrelation volume, or the integral of the correlation function in d -dimensions. Ideally, the CT concentrations from groundwater samples would be used to estimate the autocovariance function. We only have the CT concentrations that were kriged from the original data and used in the RT3D model. We calculated the autocovariance function from these initial conditions separately in the horizontal and vertical directions (Appendix A:). In the vertical direction, we estimated an average particle density of approximately 2 particles per meter. In the horizontal direction, we estimated a much lower density (because of greater correlation lengths in the horizontal space) of approximately 0.1 to 0.3 particles per square meter. To save computation time, only initial concentrations within the well field area were considered. The appropriate well field area was determined by MODFLOW capture zone analysis (traced by backward tracking of inert particles), which suggested that only the area $0 < x < 42$ m, $15 < y < 41$ m, and $2 < z < 20$ m are inside the influence of the well field for the duration of this test. So the volume of aquifer in which we simulate transport and reaction is 42 m \times 26 m in area \times 18 m thick and must contain an initial distribution of 4,000 to 12,000 CT particles based on the CT spatial statistics (Appendix A:).

One main objective using the PTR method is to evaluate if the observed overall reduced reaction rates in the field scale can be attributed to the incomplete mixing. Therefore, the PTR model within RW3D used all prior laboratory (batch) parameters. This is different from the RT3D model, which over-predicted degradation significantly using the batch CT reaction rate k' .

Because the concentration of injected acetate (electron donor) was more than 20 times higher (800 versus 30 mg/L) than that of nitrate (electron acceptor), the concentration profile of acetate was reported to resemble that of non-reactive tracer Br^- , even though a small amount of acetate is consumed during the reactive transport [Witt *et al.*, 1999; Phanikumar *et al.*, 2005]. Hence, for the sake of brevity, we only show the comparison of simulated and observed concentrations of CT and nitrate at monitoring wells. The simulation results include those from the RT3D model from [Phanikumar *et al.*, 2005] based on the fitted reaction rate [Phanikumar and Hyndman, 2003] and the PTR method within RW3D. Concentrations of CT and nitrate were measured at wells 9, 10, 11, 12 and 13 at five observation depths, 10.7 m, 13.7 m, 16.8 m, 19.84 m, and 22.9 m bgs. The breakthrough curves of CT and nitrate were normalized with concentrations of 0.032 and 42 mg/L, respectively [Phanikumar *et al.*, 2005]. The measured and simulated breakthrough curves of CT and nitrate are plotted in Figures ?? and 3, respectively.

Given the estimated range of initial number of CT particles from autocovariance analysis, we ran simulation tests by varying the mass each particle carries (m_p) and found that 4,612 initial particles — on the lower-end of the range of 4,000 to 12,000 — provided a good match of concentration profiles. The low end was derived based on ignoring the hole effect when integrating CT autocovariance, which may be a numerical artifact at large variogram lags. In other words, the particle number is more closely associated with the estimation of positive correlation. The CT initial particle number dictates the mass of every particle in a simulation, hence the numbers of nitrate particles, for example, is fixed by the initial mass in the aquifer. The total number of initial sorbed CT particles was calculated as 27,460 based on distribution coefficients at different layers [Dybas *et al.*, 2002]. The number of nitrate particles within the influence area of the well field was calculated as 2,867,400. The distribution of the initial particles was calculated from individual concentrations at each MODFLOW model cell and the mass each particle carries (see details in Appendix C.1). During the inoculation, 471 KC particles were added. The number of KC particles grew rapidly, especially in the attached phase, so that 100 days after inoculation approximately 1,760 detached and 76,000 attached KC particles were present in the model domain. In addition, the consumption of nitrate by the native flora was assumed to occur where the nutrient (acetate) and nitrate were both available. We also assumed, as did Phanikumar *et al.* [2002], that the population of native flora is proportional to that of KC. The calculated number of microbe particles representing the native bacteria is described in Appendix C.2. This is different from the simulation of the column experiment, where we assumed that the impact of native microbes was negligible because the column was flushed 4 weeks to achieve

a denitrifying condition [Witt *et al.*, 1999]. However, measurements of nitrate in the field suggested that the consumption rate of nitrate was beyond the capability of the limited amount of KC injected (see also [Phanikumar *et al.*, 2005]).

To account for the stochastic nature of the PTR method, we ran 50 simulations to obtain ensemble statistics for simulated concentrations, and plotted the ensemble means \pm one standard deviation (Figs. ?? and 3). We found that the relevant statistics of the simulations at most wells converged at around 30 to 40 realizations, as shown in Appendix E: . Similar to the Br^- breakthrough curves, good matches between measured and PTR simulated nitrate and CT are found in all the monitoring well locations. Simulated concentrations from both models in the upper low K zone were generally lower than those of observations, especially at the depth of 13.7 m (45 ft), similar to the breakthrough curves of bromide, as shown in Figure 1. This implies that the preferential flow was not fully captured in the MODFLOW flow field, particularly in the low K zone. This under-prediction might also be due to the kriging method interpolating hydraulic conductivities, which smoothed the variability of K . Fractal interpolation methods, such as used by Dogan *et al.* [2014], would likely improve simulation results.

As also shown in Figures ?? and 3, the standard deviation of the simulated results in some zones was relatively large. This is because the fast moving or easy reacting particles may or may not be captured in the small counting bins (capture zone) of individual wells in different model runs. The randomness of the numbers of particles reflects the imperfect mixing condition. If an infinite number of particles, which represents a complete mixing condition, were used for the simulation, the variance would be close to zero and we would expect results similar to the Eulerian model. In other words, the finite number of particles accounts for the degree of mixing in the site (Appendix D:), which explains why the apparent reaction rate was more than one order of magnitude lower in the field than obtained from the batch experiment.

The over-prediction of CT reaction rates by the RT3D model using laboratory-optimized rates was thought to be linked to the availability of electron acceptor and limitation of microbial growth at the field scale [Phanikumar *et al.*, 2005]. These factors contribute to the overall process of reactants mixing at a range of scales. To match the field measurements, Phanikumar *et al.* [2005] increased the kinetic attachment value and lowered the CT degradation rate. In contrast, the PTR model did not adjust the kinetic parameters; instead, the number of particles, which represents the mass of solutes and biomass, as well as the variability of concentrations within a fixed volume, were estimated to account for the incomplete mixing [Benson *et al.*, 2013].

Moreover, attachment/detachment process combined with the difference of degradation capability between the mobile and immobile microbes were thought to lead to the increase in CT observed in the high conductivity layers for some wells after the post-inoculation decline (e.g., well 10-75 at 22.9 m (75 ft) depth) [Phanikumar *et al.*, 2005]. Because only limited in-

formation is available for the difference of reaction rate constants between mobile and immobile bacteria, the reaction rates are assumed to be the same for both phases, as used in the RT3D model [Phanikumar *et al.*, 2005]. During the inoculation period (2 hours), the attachment coefficient for bacteria was increased by one order of magnitude by Phanikumar *et al.* [2005]. In the PTR method, the attachment coefficient is kept constant and equal to the laboratory-measured batch values. We assumed that 90% of the microbes are attached on the aquifer material during injection. This is consistent with previous studies on bacteria transport and field observations [Ding, 2010; Dybas *et al.*, 2002].

As shown in Figures ?? and 3, the biodegradation of CT and consumption of nitrate during the field-scale bioremediation are well-simulated using the PTR method with batch-scale parameters. The RMSE of the simulations from the two numerical models were calculated for both CT and nitrate. By this measure, the PTR method better predicted the CT concentration breakthrough curves in 23 of 25 wells (Fig. ??). On the other hand, the PTR method predicted a slower decline, or consumption rate, of nitrate. This is most likely because we assigned the numbers of particles based on the autocovariance of initial CT concentrations. To maintain stoichiometry, a very large number of nitrate particles were needed, which may or may not represent the spatial heterogeneity of the nitrate initial condition. The large number implies that nitrate consumption is not limited by mixing due to its high concentration and smoothness. This smoothness is reflected in the gradual overall breakthrough of nitrate in many wells in the PTR simulations. The PTR model results also show more high-frequency variability in the BTC, which most likely represents the impact of the re-circulation (pump/inject) process on the concentrations. On the other hand, the RT3D model provides smooth curves that could be the result of numerical dispersion (especially vertical mixing). The better nitrate RMSE fit is evenly split (12 to 13) between the two models.

5 Discussion and Conclusions

This study presents a series of novel developments, including the first implementation of complex reaction kinetics at the field-scale using a purely Lagrangian particle transport and reaction (PTR) code. The reasons to implement such a code are primarily: 1) to avoid the spurious mixing that grid-based Eulerian algorithms can impart; and 2) represent subgrid velocity and concentration perturbations. The difficulty that grid-based codes have in accurately simulating the degree of mixing between chemical species is accurately handled by the particle methods [Benson *et al.*, 2017; Herrera *et al.*, 2017].

The column experiment of CT biodegradation that was performed in support of the Schoolcraft field-scale experiment was simulated using the

Figure 3. Measured and simulated breakthrough curves of nitrate. The circles are the measured concentrations in five wells at five depths, the blue lines are simulations using RT3D, and the green lines with error bars are the means plus/minus one standard deviation from 50 simulations of the PTR method in RW3D. Bold, dashed numbers denote well and sampling depth (ft bgs). Numbers reflect RMSE differences between modeled and measured normalized concentrations, and wells with better PTR simulations are given a yellow background.

PTR method within RW3D. Kinetics parameters from batch experiments were directly used in this method. The results are consistent with those from the PTR simulation using a Matlab code, as presented by *Ding and Benson* [2015]. Observed concentration profiles at 10 sampling ports at both days 2 and 26 were closely matched with most measurements within one standard deviation of the ensemble mean. This contrasts with Eulerian simulations of the columns, which required reductions of the CT degradation rate parameter from 2.70 to 0.189 L mg⁻¹ day⁻¹ [*Ding and Benson*, 2015]. The column experiment simulation suggests that the PTR method within RW3D can simulate CT biodegradation, which involves processes of first-order, second-order, and Monod-type reactions, as well as attachment/detachment, growth, and decay of biomass. The upscaling of mixing that accompanied moving to the column scale was handled by the particle method through the calibration of particle numbers. These numbers are dictated by the chemical autocovariance functions that were not measured (because of the paucity of sample ports) at the beginning of the column test.

On the other hand, the statistics of the CT initial condition were measured at the Schoolcraft field site. The input files to RT3D from the study of *Phanikumar et al.* [2005] gave us an estimate of the covariance functions, and we calculated the initial particle numbers prior to PTR simulations of bioremediation. Because the PTR method explicitly accounts for subgrid mixing via the particle numbers' representation of the concentration autocovariance, the overall rates of reaction will depend on an accurate assessment of the initial concentration statistical structure (Appendix D:). Our goal here was to test the hypothesis that the PTR method is capable of differentiating the mixing effects from the empirical reaction rate reduction (or scale effect). We did not systematically assess the sensitivity on the initial condition in this paper: we leave that for a future study of model sensitivities.

Before running those simulations, we simulated the transport of bromide tracer test at the Schoolcraft site using particle tracking (RW3D) and the same velocities as an RT3D model. The RW3D simulations matched Br⁻ measurements with a longitudinal dispersivity value of 0.03 m, which is about 3 times larger than that used in the RT3D model. Due to the lack of numerical dispersion that arises from transferring back and forth from Lagrangian and Eulerian schemes, the RW3D model better matches the breakthrough curves in most observation wells.

Finally, we applied the PTR model to simulate the site bioremediation. The simulation involved the processes of solute and bacteria transport, attachment/detachment, growth and decay of biomass, as well as the reactions among CT, bacteria KC, electron donor (acetate), and electron acceptor (nitrate). The comparison between simulated and measured breakthrough curves at 25 monitoring well locations, as well as the comparison between RT3D and RW3D simulations, indicate that the PTR method can accurately simulate the field experiment without adjusting any parameters from the batch- to field- scales, particularly the CT biodegradation rate, which needed to be reduced by a factor of 22 in the RT3D model [*Phanikumar et al.*, 2005]. How-

ever, the success of the PTR method requires accurate velocity fields and an accurate assessment of the spatial autocovariance of the reactant initial condition, because these factors are the primary controls of potential mixing and dictate the number of particles used in the domain.

A large number of sites have shown the scale effect of reaction rates. One source is chemical heterogeneity, especially subgrid or unrepresented fluctuations. Another source has received more recent attention: the ADRE contains only one term that must simultaneously account for both spreading and mixing of solutes (e.g., *Kapoor et al.* [1998]; *Battiato et al.* [2009]; *Le Borgne et al.* [2010]; *Dentz et al.* [2011]; *Le Borgne et al.* [2013]; *de Anna et al.* [2014]; *Porta et al.* [2016]). Only at the very smallest scales are these two quantities of similar magnitudes. As solutes encounter more heterogeneous Darcy velocities, the spreading grows faster than the local mixing. For the ADRE to accurately describe spreading, it must overpredict mixing and vice-versa: accurate representation of mixing will under-disperse solutes and place reactants in the wrong places. A corollary is that perfectly homogeneous sites (i.e., $\text{VAR}(\ln(K)) \rightarrow 0$) would not suffer from this particular effect. A notable example of a reactive transport experiment in relatively homogeneous material is the petroleum hydrocarbon injection/biodegradation experiment in the Borden aquifer [*Schirmer et al.*, 2000]. With $\text{VAR}(\ln(K)) = 0.244$, *Schirmer et al.* [2000] used laboratory-estimated M-M parameters in a finely-discretized Eulerian field-scale model to accurately simulate aerobic degradation of injected contaminants (under natural gradient conditions). For comparison, the Schoolcraft aquifer's overall $\text{VAR}(\ln(K)) = 0.634$, about 2.6 times greater than Borden's. Because 2nd-order (including M-M) or higher reactions introduce a nonlinear amplification into any transport errors [*Benson et al.*, 2017], we conclude that the scale effect due to velocity fluctuations will manifest at all but the most homogeneous sites. Going from $\text{VAR}(\ln(K)) = 0.244$ to 0.634 appears to have made a significant difference, although there were other differences in the two experiments that may have contributed, including aerobic versus anaerobic conditions, small and relatively homogeneous injected contaminant volumes at the Borden site, and the natural-flow versus forced-recirculation conditions.

In this study, we used the original PTR method from *Benson and Meerschaert* [2008], which requires that all reactant particles carry the same amount of mass. Because of the large difference in concentrations of CT and nitrate, a very large number of nitrate particles (≈ 3 million) were assigned in the simulation and thus it requires a large computational effort relative to the prior RT3D model (approximately 22 versus 4 hours on a 3.4 GHz i7-3770 processor with 24 Gb RAM). However, new PTR methods address the problem of large particle numbers and concentration discrepancies, by either allowing particles to have variable mass [*Bolster et al.*, 2016; *Benson et al.*, 2017], allowing particles to carry multiple species [*Benson and Bolster*, 2016], or larger “footprints” by using kernels with optimal particle influence instead of the current Dirac-delta functions [*Fernández-García and Sánchez-Vila*, 2011; *Rahbaralam et al.*, 2015; *Schmidt et al.*, 2017]. Much shorter com-

Figure A.1. Estimated horizontal autocovariance functions in each of 39 non-zero layers from the initial CT concentrations from the RT3D file [Phanikumar *et al.*, 2005]. The thickness weighted average of the layers is shown with a thick black line.

Figure A.2. Estimated vertical autocovariance functions for CT in the entire model domain using the initial CT concentrations from the RT3D file [Phanikumar *et al.*, 2005].

putation times should be expected with these methods and a more rigorous benchmarking of the current study, including parameter uncertainty, could be performed.

In summary, the PTR method with RW3D is capable of simulating field-scale bioremediation with equal or better accuracy than traditional methods. Furthermore, the reaction parameters transfer from the smallest scale, separating the scale-dependence of reaction rates from the underlying source of reduced reaction: poor mixing at larger scales.

A: Estimation of initial CT particle numbers (density)

Estimation of the CT autocovariance function is performed on the input files for RT3D, which has 39 non-zero layers. In the horizontal direction, the autocovariance is calculated individually in each layer using standard methods and assuming isotropy with respect to lag separation. Data pairs were grouped in lag intervals (0 0.5), (0.5 1.5), (1.5 2.5), ... (23.5 24.5). A plot of each layer's estimated autocovariance function versus radial lags is shown in figure A.1. Also plotted is the layer-thickness weighted average autocovariance, which has a summed correlation function (which includes the "hole effect" of negatively correlated values) of $l = 3.2$ m. Ingoing the negative values gives a visual estimate of the correlation length on the order of 5 m. Extending to 2- d , it is safe to say that the 2- d correlation volume is on the order of 10 to 30 m². The total CT mean and variance within the non-zero layers in the RT3D input file are 0.0127 and 5.3×10^{-5} , respectively, so that the average initial particle density (see Paster *et al.* [2014] for a derivation) in the horizontal is $\rho = \bar{C}^2/(\sigma^2 l^d) \approx 3/l^d \approx 0.3$ to 0.1 particle per square meter. In the vertical, more noise was resolved, and the average autocovariance function has a 1- d correlation length of about 1.5 m (Fig. A.2), so that the average particle density in the vertical direction is about 2 particles per meter.

B: Modification on the ADRE-based model and differences in the simulation

In the ADRE-based model in [Phanikumar *et al.*, 2005], as listed in Eqs. 4, a correction factor $[1 - A/(K_A + A)]$ was added to the bacteria decay term to account for the increase of decay rate at low nutrient concentration [Beefink *et al.*, 1990; Phanikumar *et al.*, 2005]. However, Beefink *et al.* [1990] proposed this correction term because they considered the growth and

decay of biomass together (or net growth) in their study. Moreover, during the bioremediation in the Schoolcraft field, the concentration of acetate was nearly three orders of magnitude higher than the half saturation constant (800 mg/L versus. 1 mg/L) and acetate has been continuously added to the system, the correction term was always close to zero in the well field. This results in nearly no decay in the equation.

To assess the effect of the modification from Eqs. 4 to Eqs. 5, we ran the RT3D model with both equations in parallel. As shown in Figures B.1 and B.2, the differences using these two equations were minor, especially at locations with high concentrations (deeper part, e.g., wells at 65 and 75 feet bgs), RT3D simulations using the two equations were nearly overlapped. This is because the concentration of injected acetate (electron donor) was more than 20 times higher (800 versus 30 mg/L) than that of nitrate (electron acceptor). Moreover, it is common in field bioremediation systems that more electron donor (e.g., acetate) than needed is added to promote the reactions [Alexander, 1999; Dybas *et al.*, 1998; Finneran *et al.*, 2002; Anderson *et al.*, 2003; Williams *et al.*, 2011]. As reported by Witt *et al.* [1999] and Phanikumar *et al.* [2005], the concentration profile of acetate resembles that of Br^- , even though a small amount of acetate is consumed during the reactive transport. Similarly, the correction factor applied to the detachment term has negligible effect on the simulations.

C: Correlation of particle numbers with initial concentrations and injections

C.1 Initial concentrations

Initially, CT and nitrate were present in the groundwater system. In the PTR simulation, initial numbers of CT and nitrate particles were calculated based on the concentrations from the RT3D model and groundwater volumes.

C.1.1 CT in groundwater

The observed CT values were divided into six layers (28–15.5 m, 15.5–11.5 m, 11.5–8 m, 8–5 m, 5–2 m, and 2–0 m bgs, respectively) and kriged as separate zones [Phanikumar *et al.*, 2005].

The RT3D model has the kriged CT initial concentrations (c_i), which were used directly to calculate the number of particles at each MODFLOW/RT3D model cell. The total number of particles is based on the total mass of CT in groundwater, M_{CT} :

$$M_{CT} = \sum_{i=1}^{r \times n \times l} c_i \cdot V_i \cdot \theta \quad (C.1)$$

where V_i = finite-difference cell volume; θ = porosity; r = number of rows in the model; n = number of columns in the model; and l = number of layers in the model. The total number of CT particles is pre-determined by the autocovariance (Appendix A:), i.e., N_{CT} is within the range of 4000 to 12,000. Model simulations suggested that simulations with $N_{CT} = 4,612$ provided a reasonable match of measurements, so the mass of each particle of each

Figure B.1. RT3D-simulated breakthrough curves of CT using Eqs.4 and 5. Symbols are the measured concentrations; the black solid lines are RT3D simulations using Eqs.4, and the magenta dashed lines are simulated results from RT3D model using Eqs. 5. Subplots for individual wells are the same locations as in Figures ?? and 3.

Figure B.2. RT3D-simulated breakthrough curves of nitrate using Eqs.4 and 5. Symbols are the measured concentrations; the black solid lines are RT3D simulations using Eqs.4, and the magenta dashed lines are simulated results from RT3D model using Eqs. 5. Subplots for individual wells are the same locations as in Figures ?? and 3.

species is given by

$$m_p = \frac{M_{CT}}{N_{CT}} \quad (C.2)$$

The concentration of CT was from 0 to 0.00429 mg/L. The total mass was calculated as 371.04 g. The mass within the influence area of the well field was about 73.3 g and the mass each particle carries is 0.016 g.

C.1.2 Sorbed CT

The initial sorbed CT is assumed to be in equilibrium of aqueous CT. The distribution coefficients were reported vary with the depth, from 0.145 to 0.353 L/kg [Dybas *et al.*, 2002; Phanikumar *et al.*, 2005]. The numbers of sorbed CT particles are calculated from the aqueous CT concentration and distribution coefficients at different depth.

$$M_{SCT} = \frac{c \cdot K_d \cdot \rho_b}{\theta} = \sum_{i=1}^{r \times n \times l} \frac{c_i \cdot K_{dl} \cdot \rho_b}{\theta} \quad (C.3)$$

The number of sorbed CT particles would be Eq. C.4.

$$N_{SCT} = \frac{M_{SCT}}{m_p} \quad (C.4)$$

C.1.3 Nitrate

Initial nitrate concentrations were fairly constant across the region, the layer averaged concentrations were used for the current simulation. The number of particles used for initial nitrate in the system is calculated similarly as that of CT.

$$N_{Nitrate} = \frac{M_{Nitrate}}{m_p} = \frac{S \cdot V \cdot \theta}{m_p} \quad (C.5)$$

The concentration of nitrate (S) at each layer are constant, thus, the calculation of mass is conducted on layers, instead of model cells. The concentrations were found ranging from 21.62 mg/L to 44.25 mg/L from 10.6 m bgs to 27.4 m bgs. Linear interpolation is used to assign nitrate concentration to different layers. The mass of initial nitrate was calculated as 116,514 g within a smaller influence zone of the well field with length of 41.43 m, width of 14.4 m, and the effective porosity of 0.3. Model tests indicated that a smaller zone for nitrate did not affect the simulation results due the uniform distribution of nitrate, but it saved the computational time.

C.2 Injected mass

During inoculation, at day 67, the concentrations of KC and acetate injected to the biocurtain were 10^6 CFU/mL and 800 ppm, respectively [Phanikumar *et al.*, 2005]. Certain numbers of particles were simulated to be injected based on the fluxes of injection and the addition of constituents.

The injection contained the concentrations of KC at 10^6 CFU/mL, while 1 CFU/ mL is approximately equal to 1.67×10^{-7} ppm for strain KC [Phaniku-

mar *et al.*, 2002], thus, the concentration of KC in the injection was 0.167 mg/L.

The total number of KC particles added to the injection wells is given by:

$$N_{KC} = \frac{M_{KC}}{m_p} = \frac{E \cdot Q_{inj} \cdot t}{m_p} = \frac{\sum_{i=1}^{15} \sum_{j=1}^2 E_{i,j} \cdot Q_{inj,i,j} \cdot t_j}{m_p}, \quad (C.6)$$

where M_{KC} = mass of KC injected through the inoculation, $Q_{inj,j}$ = volume of groundwater injected to well i at period j , and t_j is the duration of injection at period j . The stress periods 45 and 47 had durations of 0.05555 and 0.04514 day, respectively, with injection volumes of approximately 12 m³ and 9.75 m³, respectively. The number of acetate particles is proportional to the nitrate particles based on the ratio of concentrations in the injection of re-circulation processes, which were 800 mg/L for acetate and 30 mg/L for nitrate. After inoculation, acetate concentrations injected were one order-of-magnitude lower, 80 mg/L.

C.3 Consumption of Nitrate by indigenous microflora

The consumption of nitrate by the native flora was assumed to occur where acetate was available. The consumption is represented with parameter γ . Even though endogenous respiration is the process by which microbes consume cell reserves in the absence of an electron donor (acetate) and continue to use an electron acceptor (nitrate), we use the same assumption that *Phanikumar et al.* [2002] made, which states that the population of native flora is proportional to that of KC.

Phanikumar and Hyndman [2003] estimated the γ term as 18.89 day⁻¹. Based on the model of nitrate consumption in *Phanikumar et al.* [2005], the proportion of native flora over KC is related to the ratio of γ over μ_{max}/Y_n as shown in (4)c. Given the laboratory-obtained specific growth rate (nitrate utilization rate), $\mu_{max} = 3.11 \text{ day}^{-1}$, and the yield for nitrate, $Y_n = 0.25 \text{ mg cells/mg substrate}$. The population of native microflora would be 4.554 times greater than KC.

D: Perturbation analysis

We adopt the methodology of *deAnna et al.* [2011], *Tartakovsky et al.* [2012], and *Paster et al.* [2014] to examine the components of the ADRE that contribute to reduced effective reaction rates and to assess whether the Lagrangian method is an appropriate tool to simulate these components. Assume that the ADRE with bimolecular reaction has random components v , C_A , and C_B with means denoted by overbars and zero-mean fluctuations denoted by primes. For simplicity we assume that the local dispersion is relatively constant:

$$\frac{\partial(\bar{C}_i + C'_i)}{\partial t} = -\nabla \cdot [(\bar{v} + v')(\bar{C}_i + C'_i) + D\nabla(\bar{C}_i + C'_i)] - k(\bar{C}_A + C'_A)(\bar{C}_B + C'_B). \quad (D.1)$$

Taking the ensemble mean,

$$\frac{\partial \bar{C}_i}{\partial t} = -\nabla \cdot [\bar{v}\bar{C}_i - D\nabla\bar{C}_i + \overline{v'C'_i}] - k(\bar{C}_A\bar{C}_B + \overline{C'_A C'_B}). \quad (\text{D.2})$$

So to first order, the new terms relative to the ADRE are a macrodispersion term and a modification of the macroscopic reaction rate by the concentration cross-covariance as in the case of purely diffusive transport [Tartakovsky *et al.*, 2012; Paster *et al.*, 2014]. It is worth discussing each of the terms on the right hand side of (D.2) with respect to “subgrid” quantities. The mean advection and local dispersion of the of the mean concentration (the first and second terms) as well as the reaction of the mean concentrations (the fourth term) are the only terms solved at a grid scale by typical Eulerian transport codes. Particle methods also represent these mean velocities and mean concentrations. However, the subgrid velocity perturbations are also solved by particles [Herrera *et al.*, 2017], i.e., velocities are interpolated between grid velocities to particles depending on their position within a cell [LaBolle *et al.*, 1996]. Therefore, shear, compression and dilation (all of which contribute to mixing and reaction [de Barros *et al.*, 2012; Engdahl *et al.*, 2014]) can be tracked by particles within cells. Furthermore, properly defined, the particles can also represent concentration perturbations at any scale as demonstrated by Paster *et al.* [2014] and Schmidt *et al.* [2017]. It is also worth noting that no study has examined the solution of (D.2) by particles in the most general cases, but several have looked at simpler systems where the velocity perturbations are known functions of space. In particular, Porta *et al.* [2012a] examined Poiseuille flow in which v' is parabolic across an aperture, and Porta *et al.* [2012b, 2013] did numerical volume averaging in an idealized unit cell. The former study showed that the particle method was an accurate simulator of the volume-averaged reacting system. The latter study showed that naively upscaled ADRE equations will follow the reaction-rate scaling that we seek to eliminate by representing subgrid fluctuations. To isolate the term that modifies the reaction rate in (D.2), first subtract the mean from the total equation:

$$\frac{\partial C'_i}{\partial t} = -\nabla \cdot [\bar{v}C'_i + v'\bar{C}_i + v'C'_i + \overline{v'C'_i} - D\nabla C'_i] + k\overline{C'_A C'_B} - k(\bar{C}_A C'_B + \bar{C}_B C'_A + C'_A C'_B). \quad (\text{D.3})$$

Now take (D.3) for $i = A$ multiplied by C'_B and add to (D.3) for $i = B$ multiplied by C'_A . Discarding third-order in perturbation terms and using fluid incompressibility yields

$$\begin{aligned} \frac{\partial C'_A C'_B}{\partial t} = & -\bar{v} \cdot \nabla C'_A C'_B + D\nabla^2 C'_A C'_B - C'_B v' \cdot \nabla \bar{C}_A - C'_A v' \cdot \nabla \bar{C}_B - 2D\nabla C'_B \cdot \nabla C'_A \\ & -k(\bar{C}_A C'^2_B + \bar{C}_B C'_A C'_B + \bar{C}_B C'^2_A + \bar{C}_A C'_A C'_B) \end{aligned} \quad (\text{D.4})$$

Define $g = \overline{C'_A C'_B}$, $f_A = \overline{C'^2_A}$ and $f_B = \overline{C'^2_B}$ and taking the ensemble average of (D.4) gives

$$\begin{aligned} \frac{\partial g}{\partial t} = & -\bar{v} \cdot \nabla g + D\nabla^2 g - \overline{C'_B v'} \cdot \nabla \bar{C}_A - \overline{C'_A v'} \cdot \nabla \bar{C}_B - 2D\overline{\nabla C'_B \cdot \nabla C'_A} \\ & -k(\bar{C}_A f_B + \bar{C}_B g + \bar{C}_B f_A + \bar{C}_A g). \end{aligned} \quad (\text{D.5})$$

Similarly for f_i ($i = A, B$); $j = (B, A)$

$$\begin{aligned} \frac{\partial f_i}{\partial t} = & -\bar{v} \cdot \nabla f_i + D \nabla^2 f_i - 2\overline{C'_i v'} \cdot \nabla \bar{C}_i - 2D \nabla \bar{C}'_i \cdot \nabla \bar{C}'_i \\ & - 2k(\bar{C}_i g + \bar{C}_j f_i). \end{aligned} \quad (\text{D.6})$$

It was this system of equations, with $\bar{v} = v' = 0$, that was solved by *Paster et al.* [2014] both analytically and with the particle method. They showed that the particle method was more accurate in that case because it does not throw out any higher-order terms (required for analytic closure). The interesting aspect of (D.5) and (D.6) is that the concentration auto- and cross-covariances follow an advection-diffusion equation with additional “macro-mixing” terms. Classic long-term closures for the terms $\overline{C'_i v'}$ for conservative tracers are often assumed to take the form $D_{macro} \nabla C_i$ [Taylor, 1953; Gelhar et al., 1979]. Inclusion of these types of terms would lead to additional terms of the form $D_{macro} \nabla C_i \cdot \nabla C_j$, which have been shown to represent the local mixing of constituents i and j [Le Borgne et al., 2010]. It is also worth noting that an attempt to analytically quantify the relative contributions of the various terms in (D.5) and (D.6), which dictate the evolution of g and deviations of the overall reaction rate (from (D.2)), will depend in complex and spatially variable ways according to local Peclet and Damkohler numbers as well as the initial conditions of g and f_i . A recent relevant discussion of the influence of the configuration of initial condition perturbations on early effective reaction rates in diffusion-only systems is given by *Ostvar and Wood* [2016]. They emphasize that the statistical chemical structure at the subgrid scale strongly influences small-time effective rates, but at late time (relative to mixing times) the subgrid fluctuations are of less importance.

To summarize, (D.5) and (D.6) show that concentration perturbations’ auto- and cross-covariance are advected, dispersed, macro-mixed, micro-mixed, and source/sinked in a coupled manner. The evolution of the cross-covariance, which is responsible for the decreased overall reaction rate in (D.2) is non-stationary and may be difficult to close accurately in an analytic sense. But prior (and separate) work has shown that the particle method can simulate all of the terms in these equations.

E: Ensemble Statistics

Because of the stochastic nature of the PTR simulation, a number of realizations were conducted to assess the sample mean concentrations and standard deviations. In this study, we ran 50 simulations, each with 25 monitoring wells and 40 timesteps. The moving averages $(C_m - C_{ave})/C_{ave}$ (where C_m is the moving average at realization m , and C_{ave} is the average for all 50 realizations) of nitrate concentration from all 25 wells and 50 realizations are plotted at three different times: before (day 30), during (day 72), and after (day 122) the inoculation process. The normalization is used because of the large differences in mean concentrations between wells and over time. Visually, the ensemble average of the simulations at most wells appears to converge at around 30 – 40 realizations (Fig. E.1). The ensemble sample

Figure E.1. The ratio of the difference between moving average (C_m) and final average (C_{ave}) over final average of nitrate concentrations at 25 well locations at days 30 (a), 72 (b), and 122 (c), which correspond to before, during, and after the inoculation process, respectively.

Figure E.2. a) Ensemble nitrate standard deviation plotted for each well as a function of simulation time. b) Ensemble coefficient of variation (standard deviation divided by ensemble mean concentration) at the same wells and timesteps.

Figure E.3. (a-c) Running coefficient of variation (CV) of the running mean nitrate concentrations at 25 well locations at days 30 (a), 72 (b), and 122 (c). (d-f) Running coefficient of variation (CV) of the running nitrate concentration variance at 25 well locations at days 30 (d), 72 (e), and 122 (f)

standard deviations may be plotted as functions of simulation time at each well (Fig. E.2a). The standard deviations at the simulation start are dictated by the initial particle numbers, among other factors such as the characteristics of the flow field. These standard deviations are on the order of 0.7 to 1 ppm. As the simulation progresses, nitrate is consumed, and the standard deviations drop by one to two orders of magnitude. To get a better idea of relative variability, we also plot the standard deviation normalized by the ensemble means - i.e., the coefficient of variation (Fig. E.2b). As the nitrate concentrations drop significantly by consumption, the concentration variabilities rise in about half the wells from about 1% to 10%. This rise implies that the reaction process is creating additional heterogeneity, in a manner consistent with older [Benson and Meerschaert, 2008] and more recent studies [Bolster *et al.*, 2016; Ostvar and Wood, 2016]. This feature might give the impression that the convergence of the ensemble statistics is adversely affected, but we also checked the rate of convergence of both sample means and variances by plotting the coefficients of variation (CV) of both the running means and running variances (\bar{C}_m and $VAR[C_m]$) over $m = 1$ to 50 realizations. As do Ballio and Guadagnini [2004], we approximate the CV by plotting the running root-mean-squared (RMS) deviations of the mean divided by the running mean (Fig. E.3a-c) and the RMS of the variance versus the running variance (Fig. E.3d-f). These measures clearly show that the ensemble statistics are converging at the expected rate of $m^{-1/2}$, consistent with the central limit theorem.

Acknowledgments

We thank associate editor Xavier Sanchez-Vila and the three other anonymous reviewers for extremely thoughtful and helpful reviews. D.D. and D.A.B acknowledge funding from NSF grant #1417145. Funding for D.F.-G. and C.V.H was partially provided by MINECO/FEDER (project INDEMNE, code

CGL2015-69768-R) and by MINECO and the UE (project WE-NEED, code PCIN-2015-248). All code and input files required to reproduce the results in this paper are located at <http://doi.org/10.5281/zenodo.848788>.

Accepted Article

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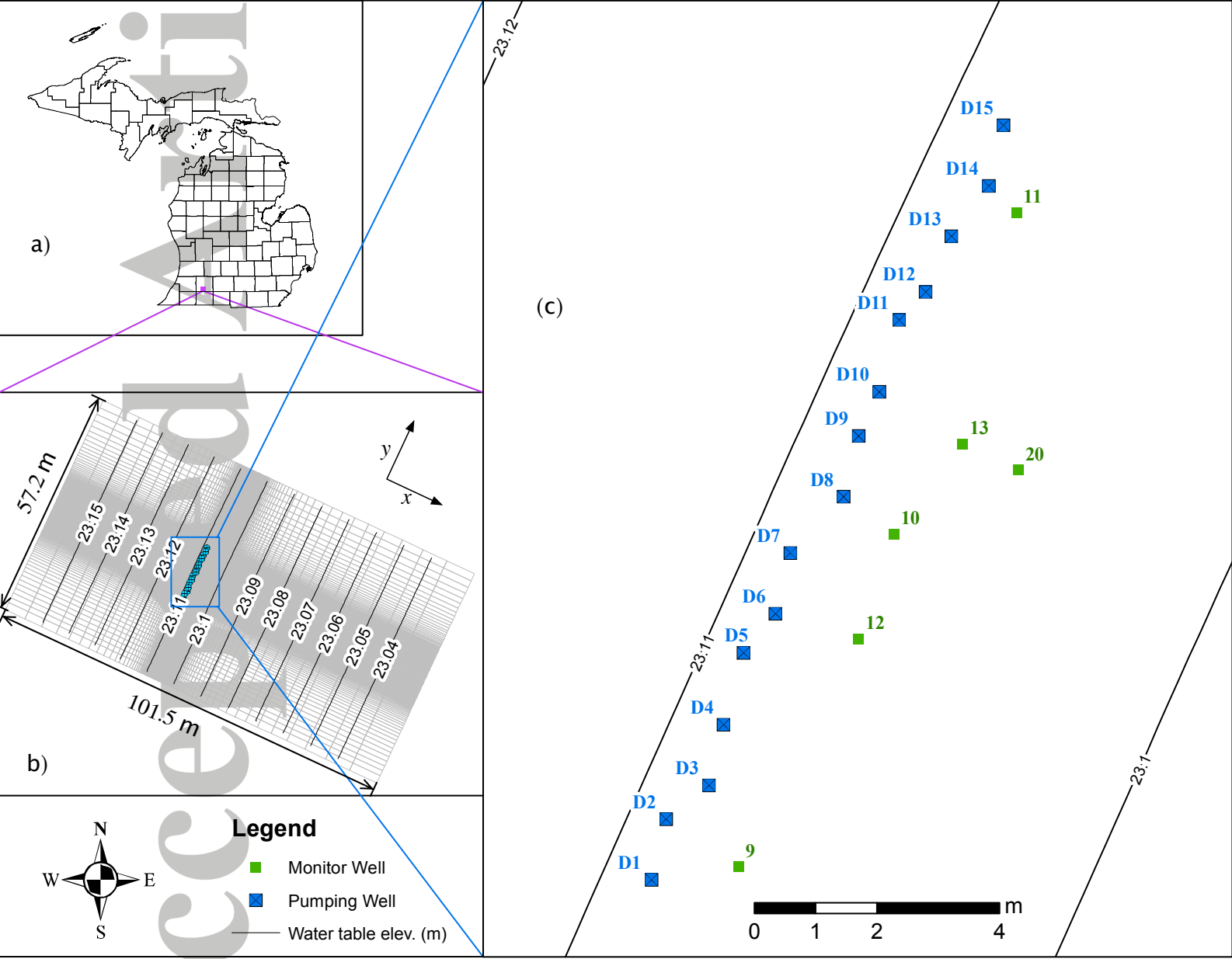
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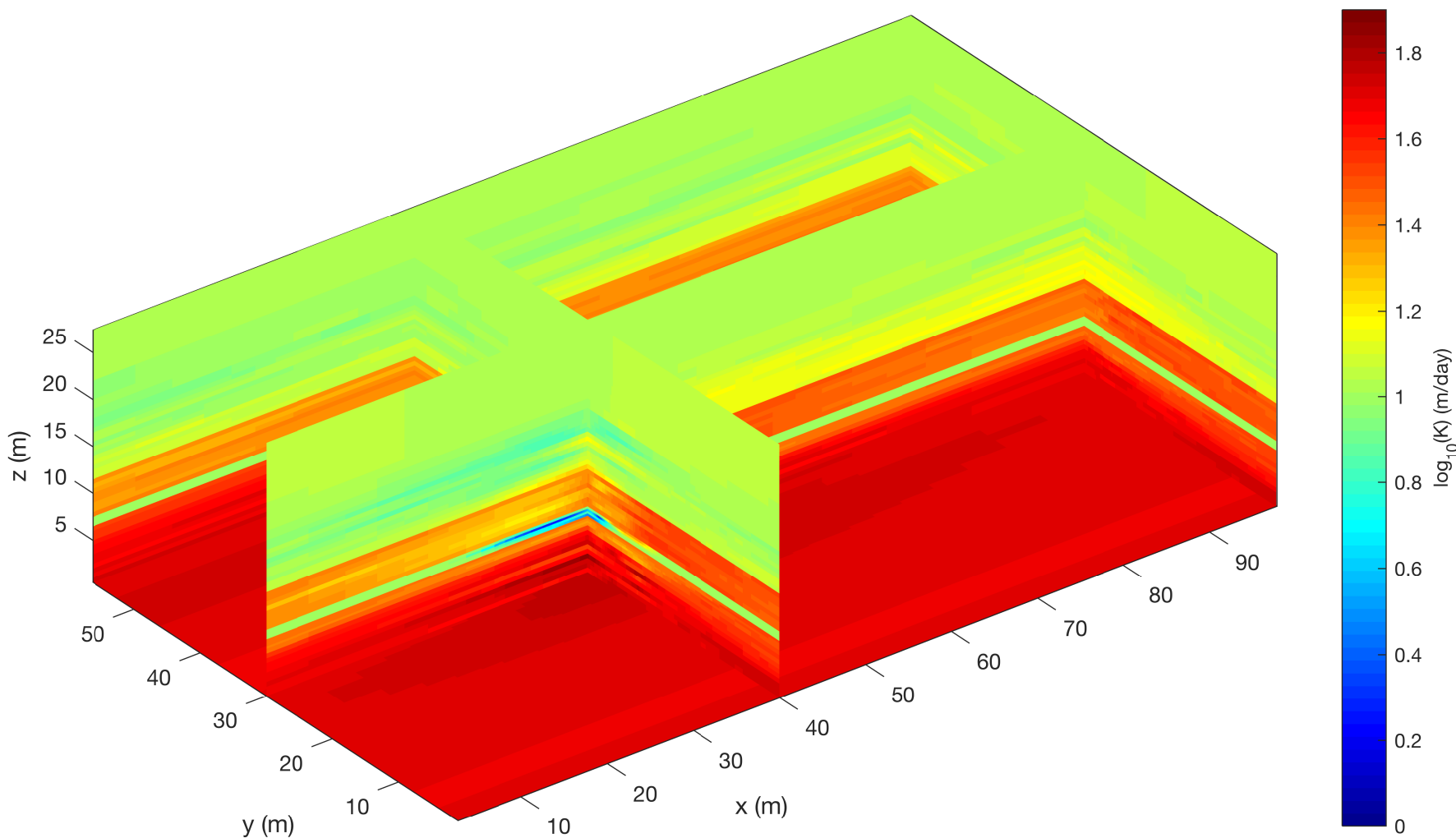
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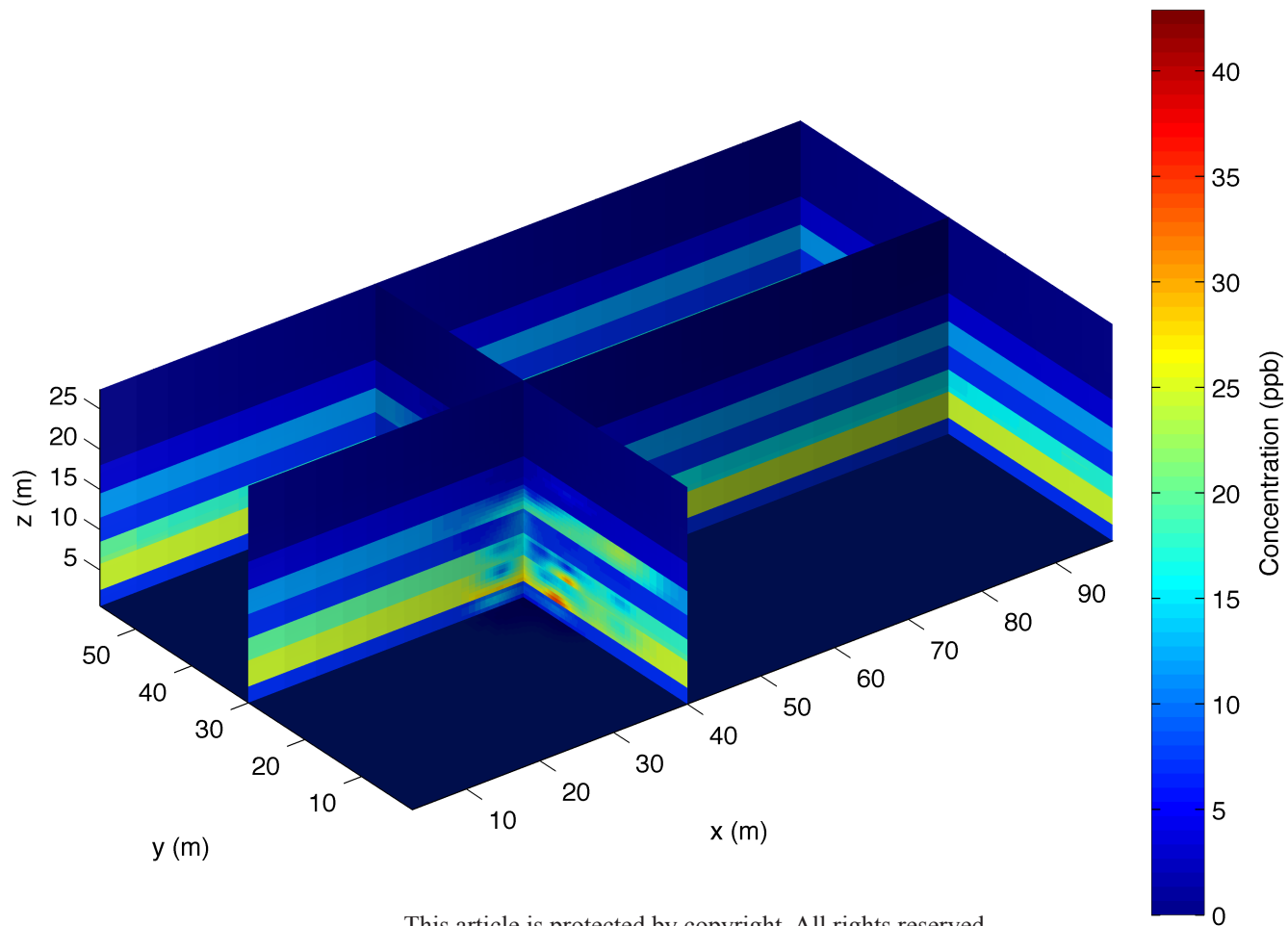
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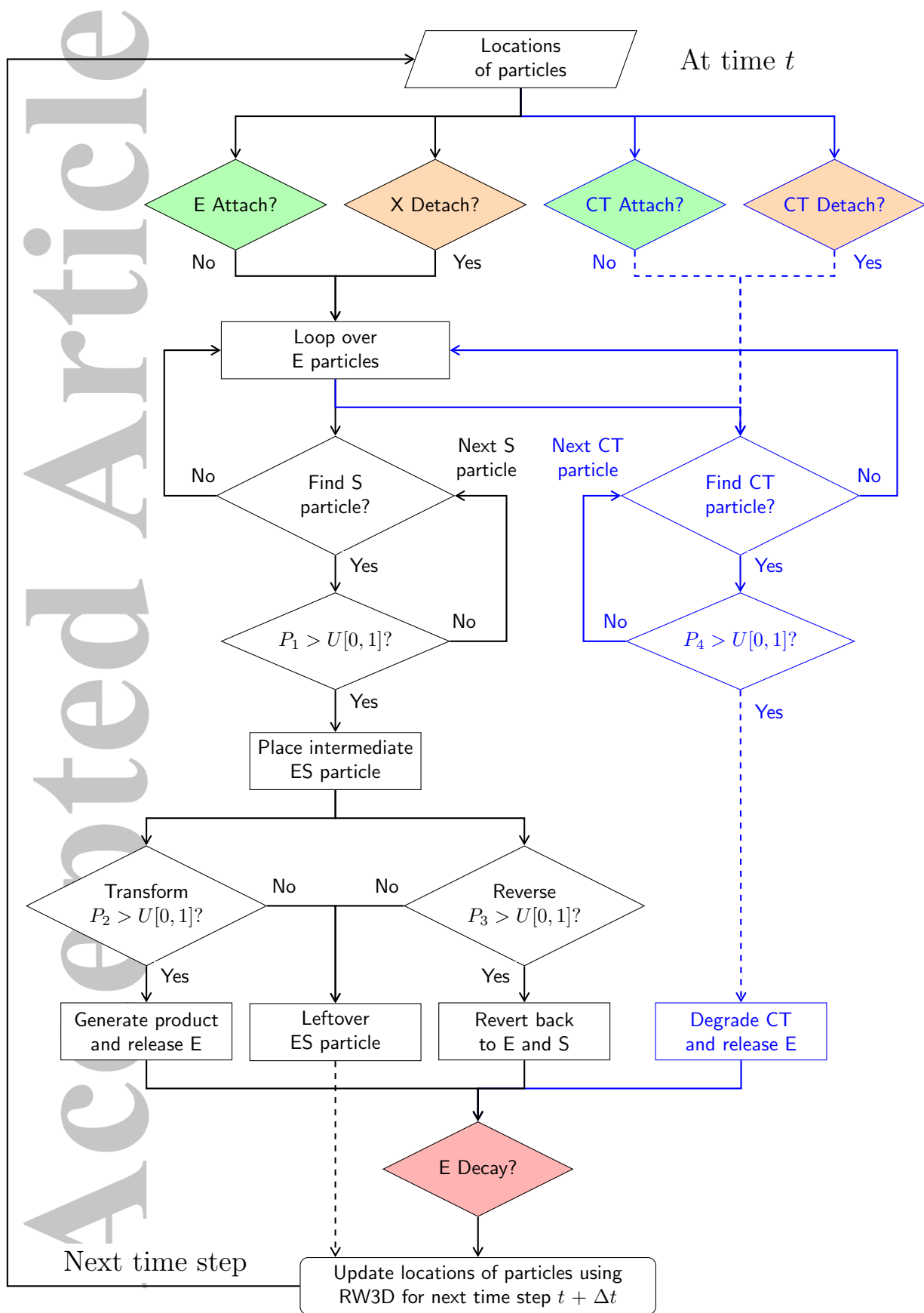
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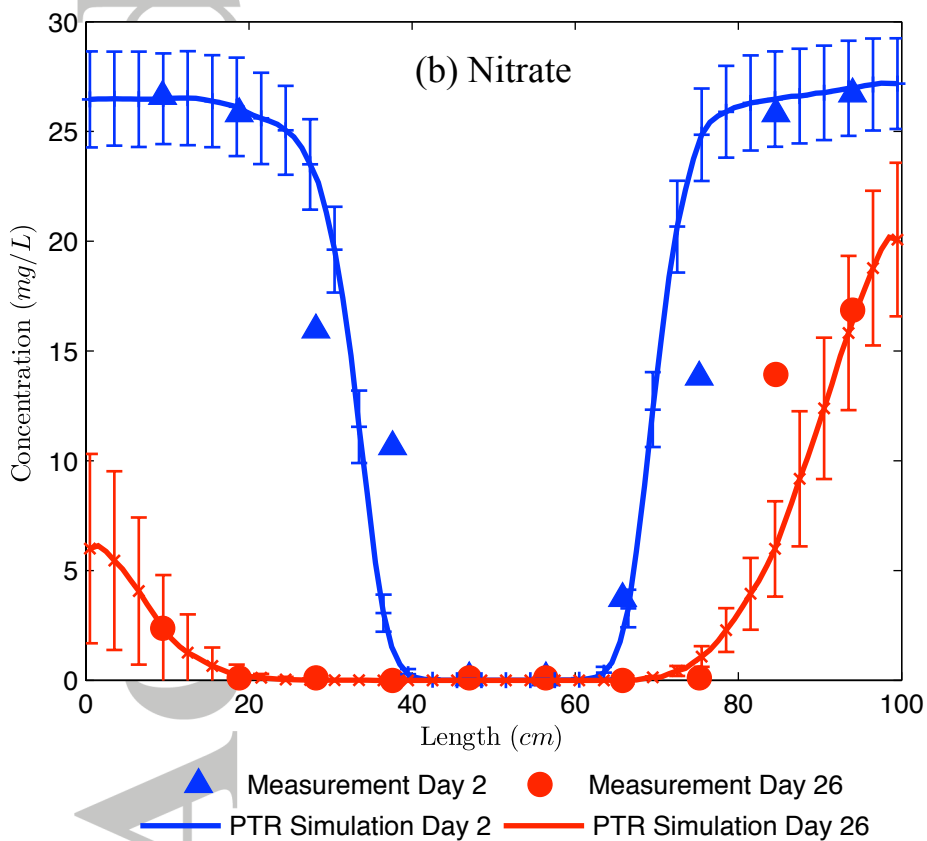
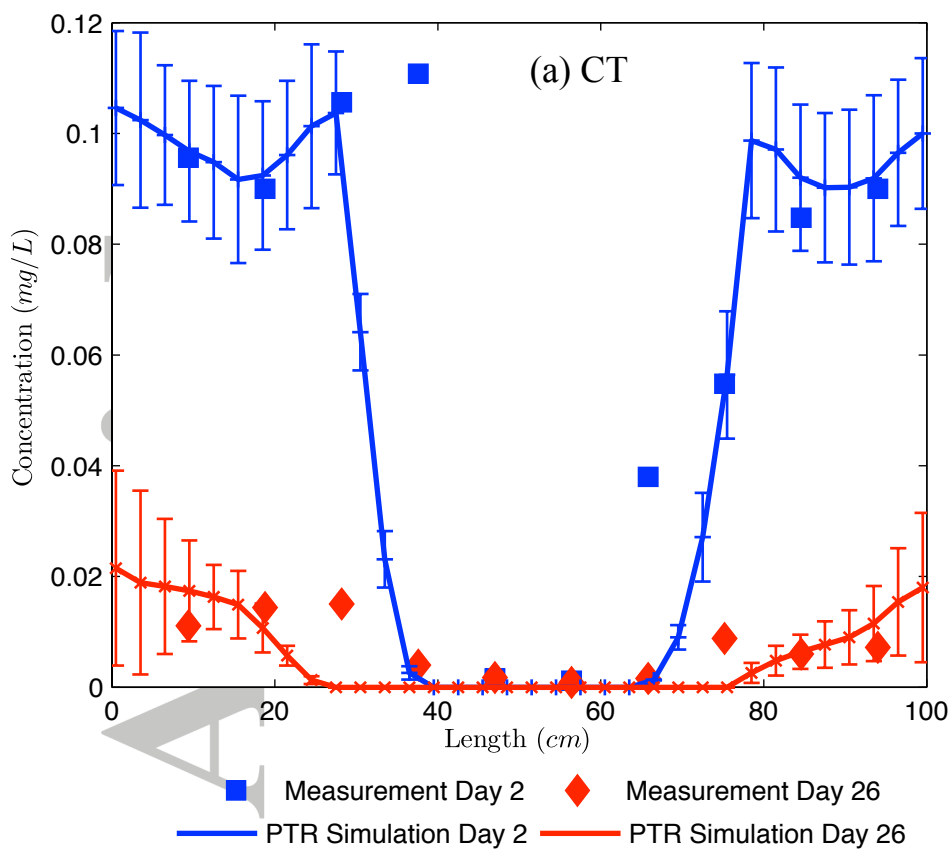
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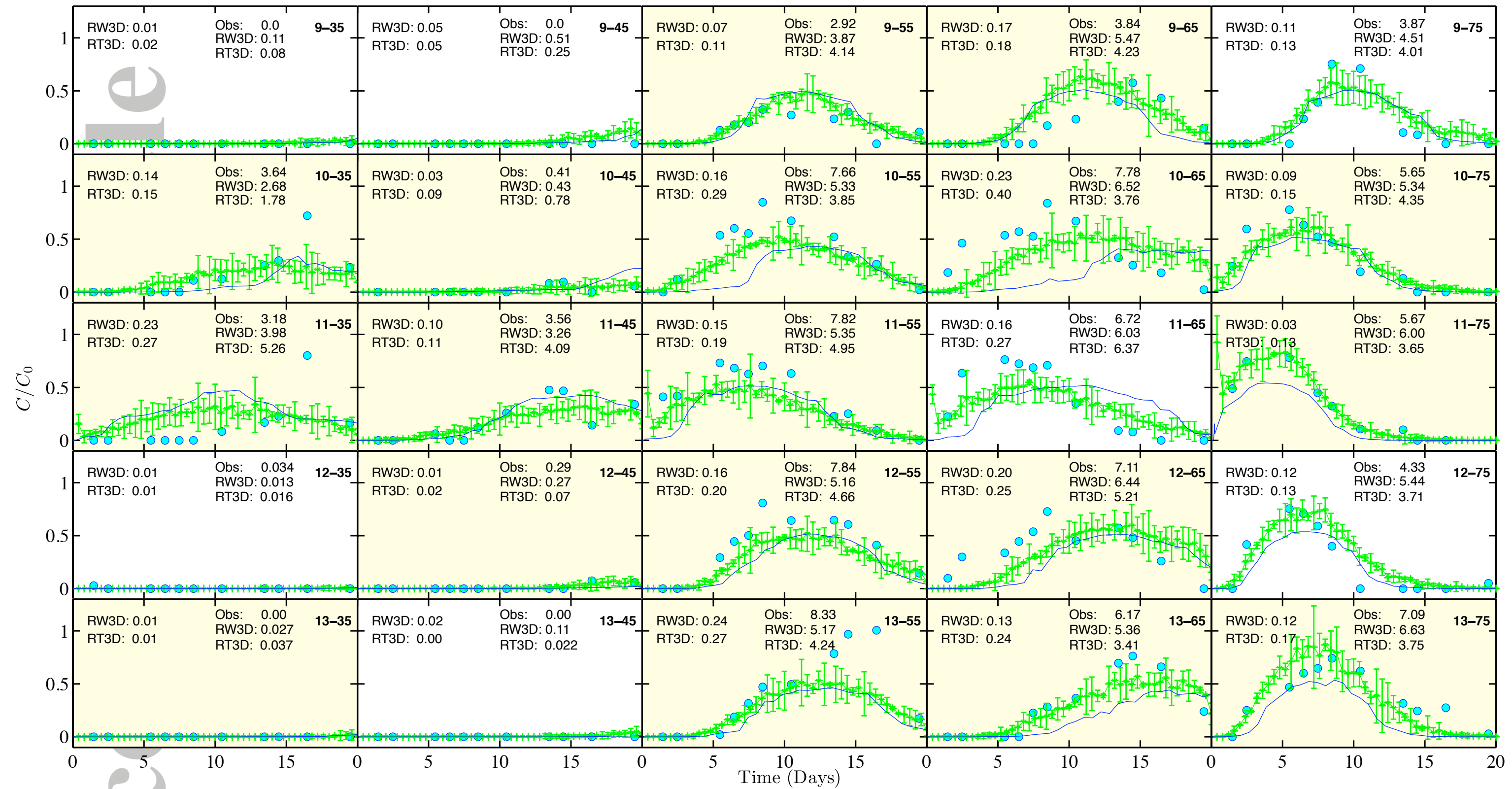
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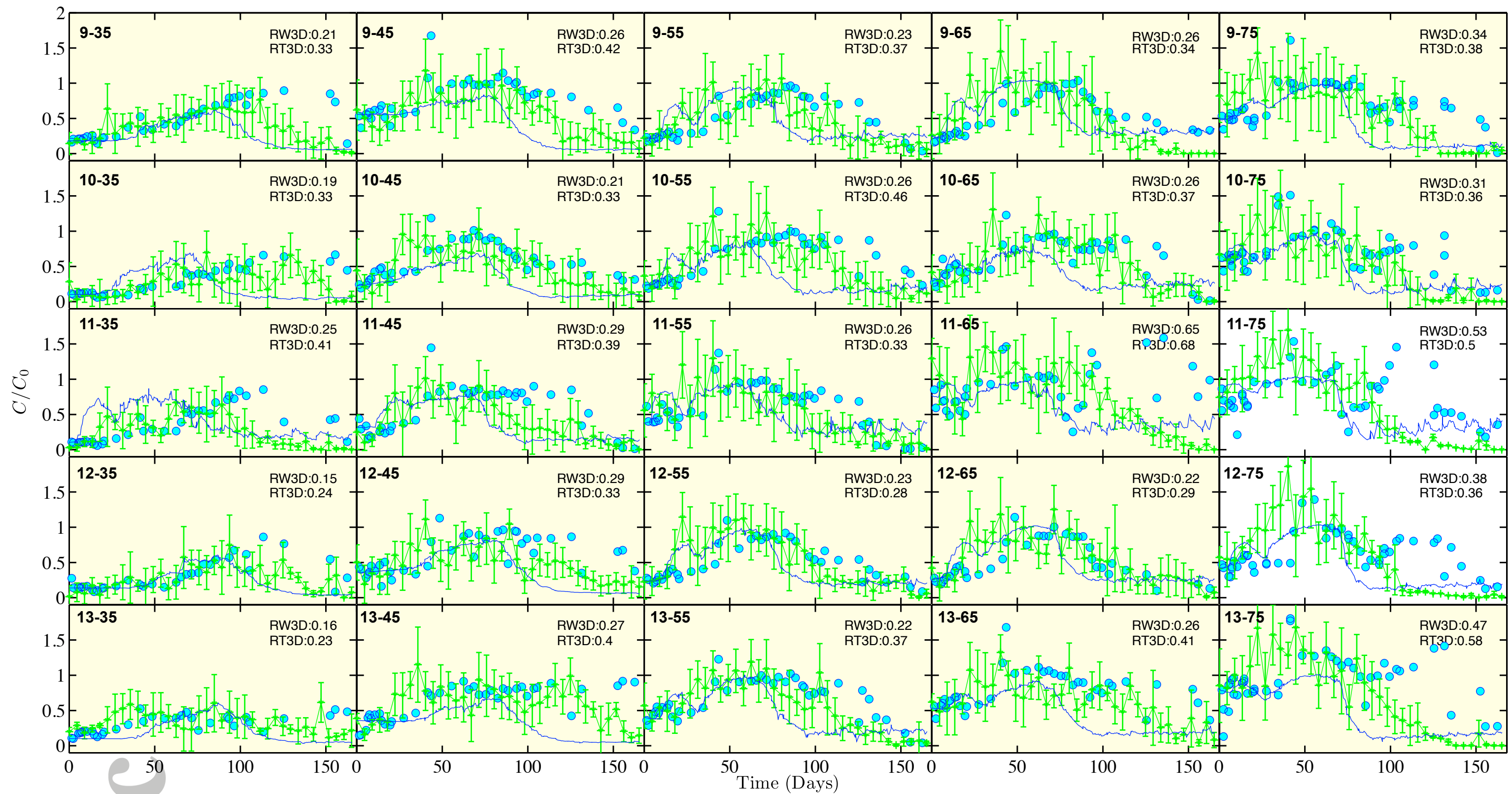
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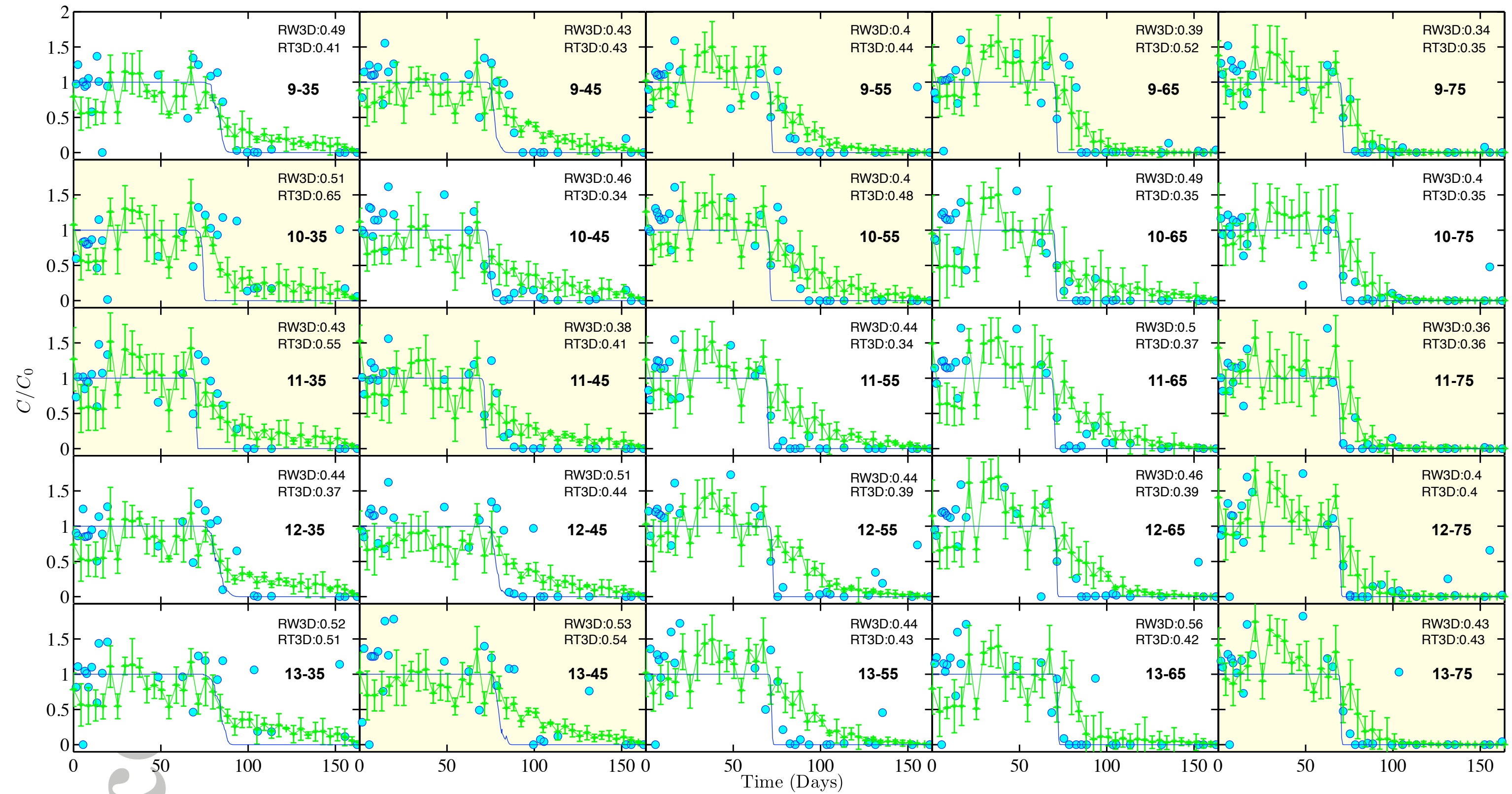
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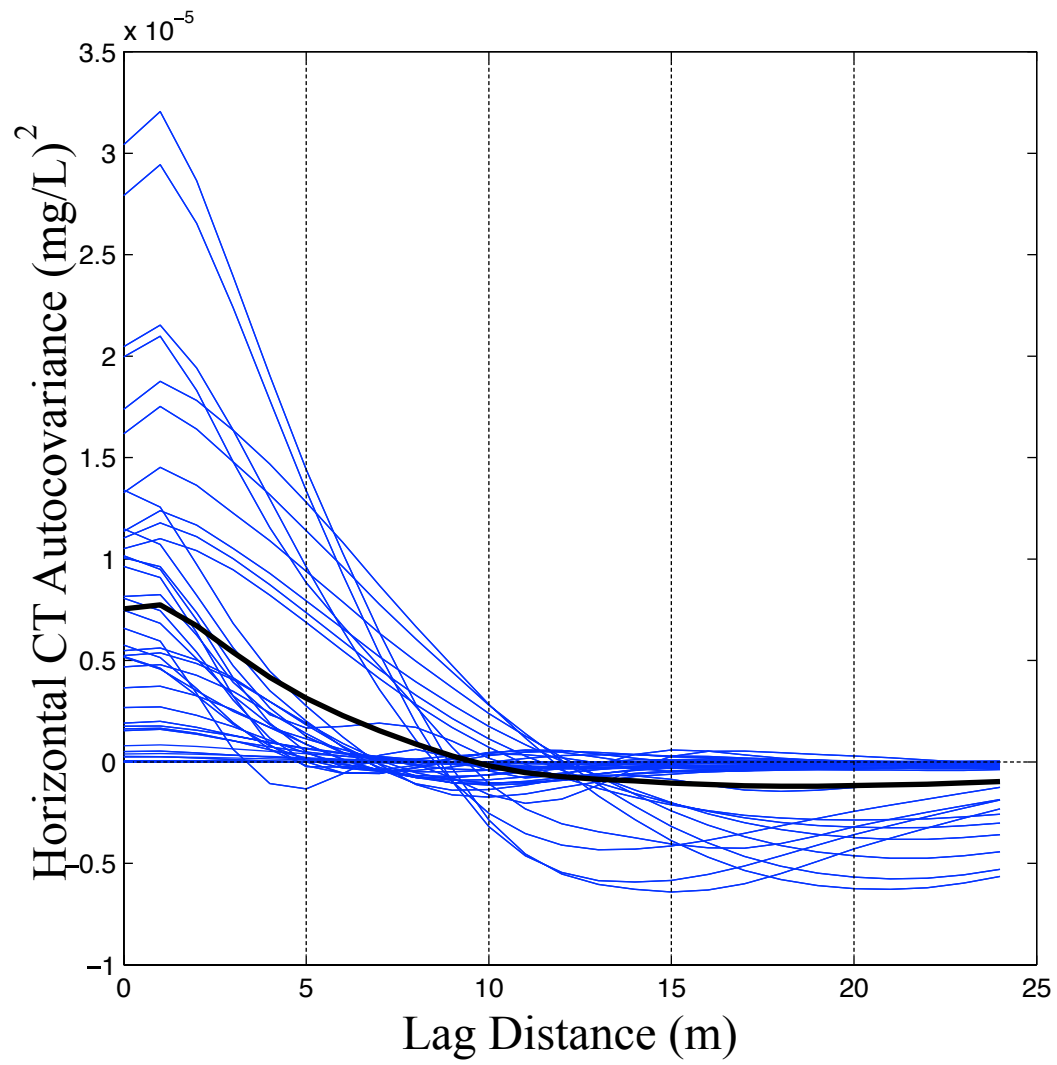
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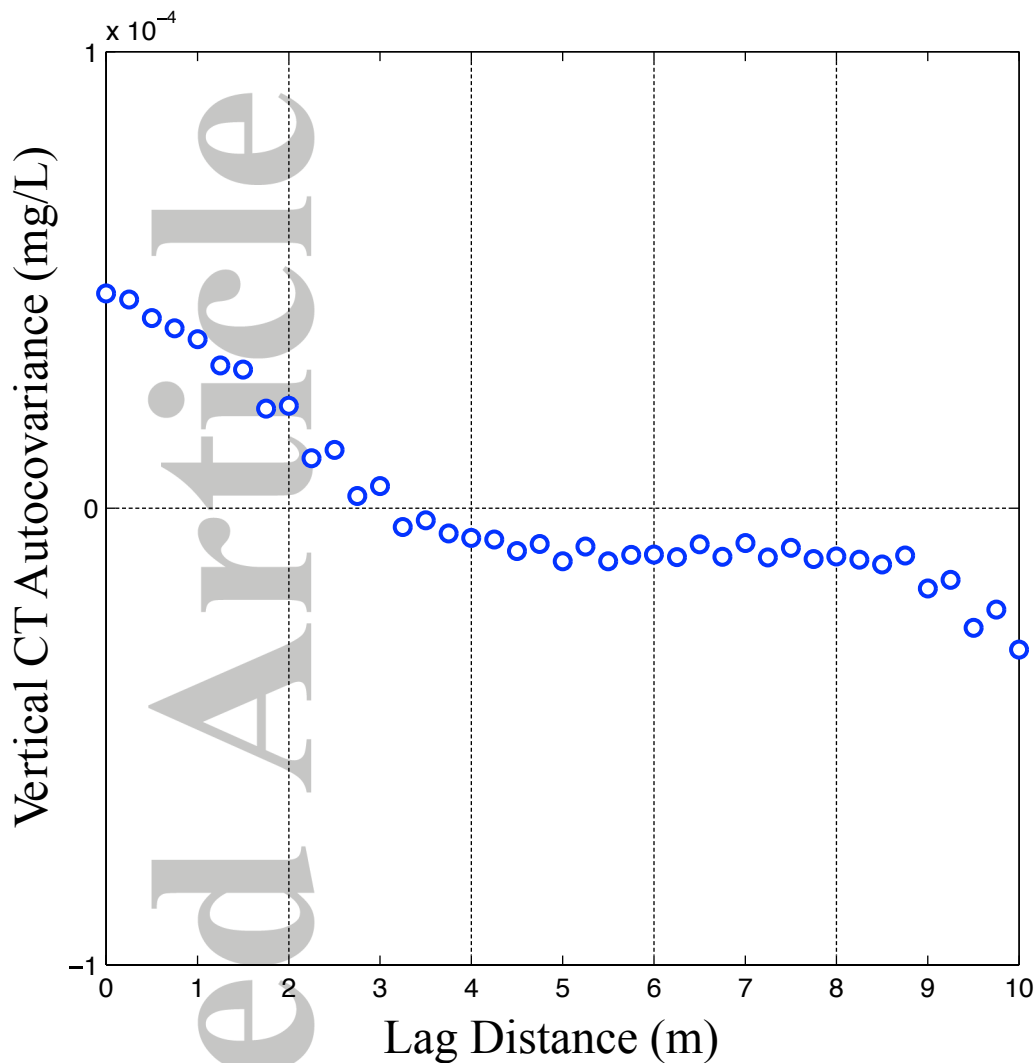
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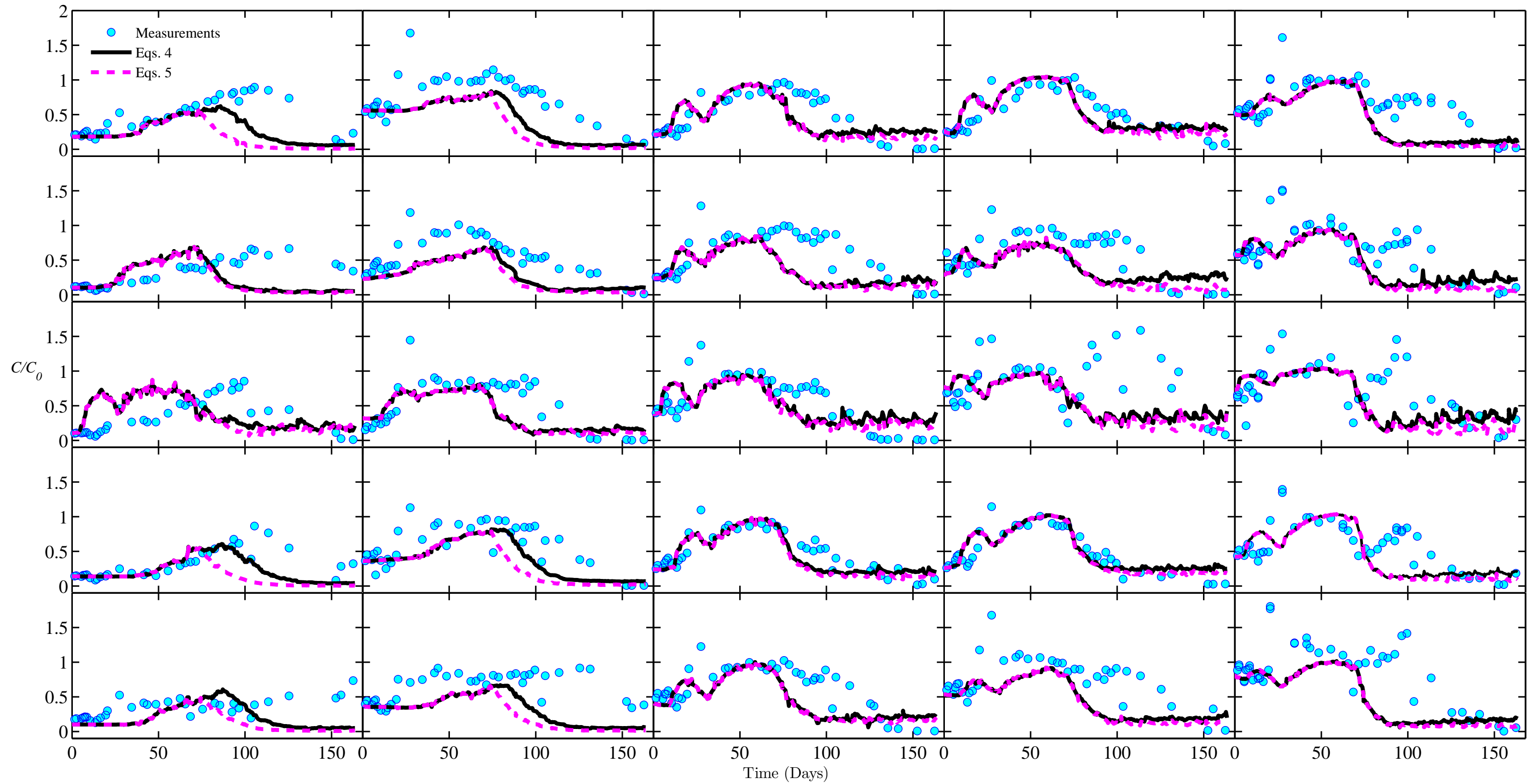
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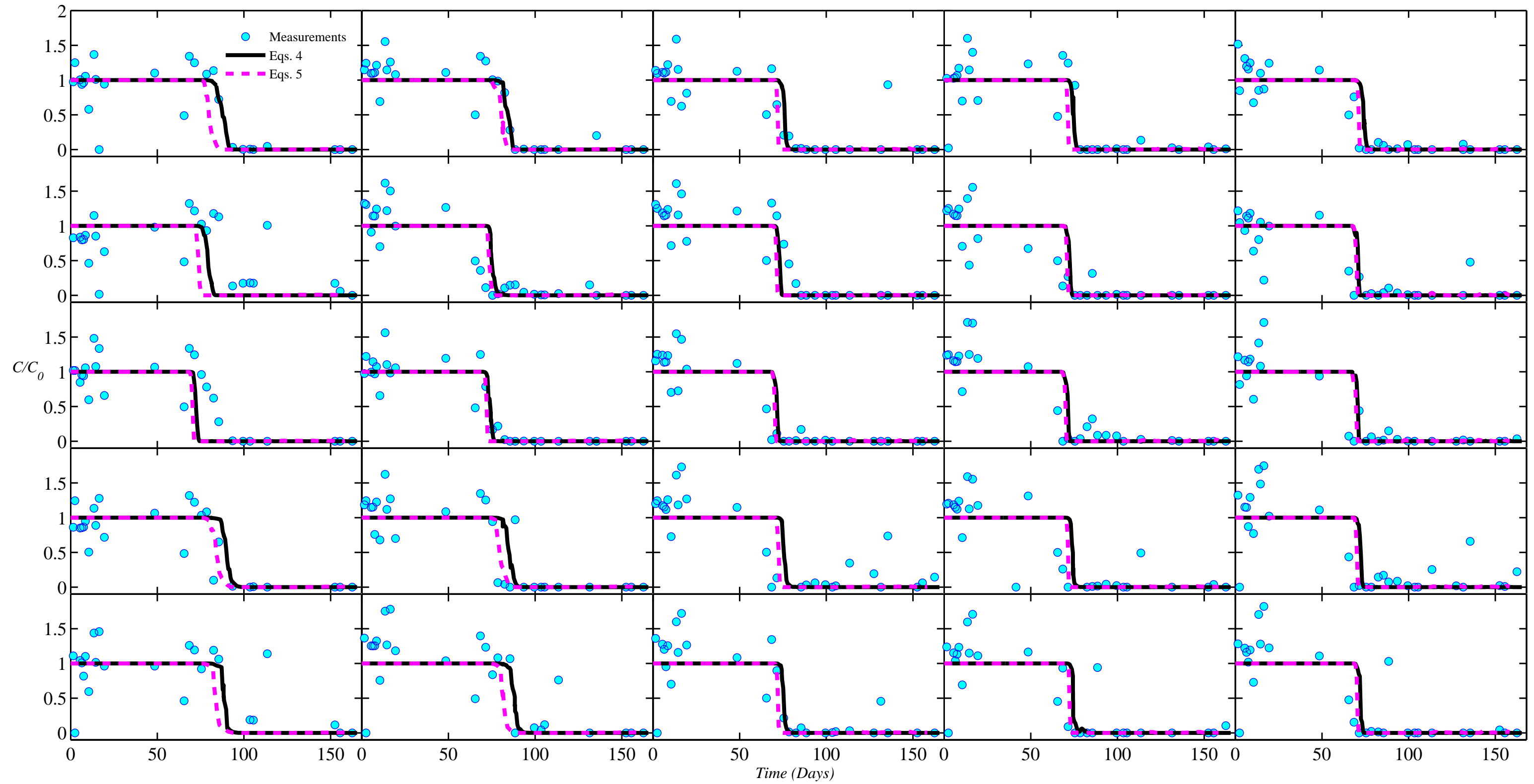
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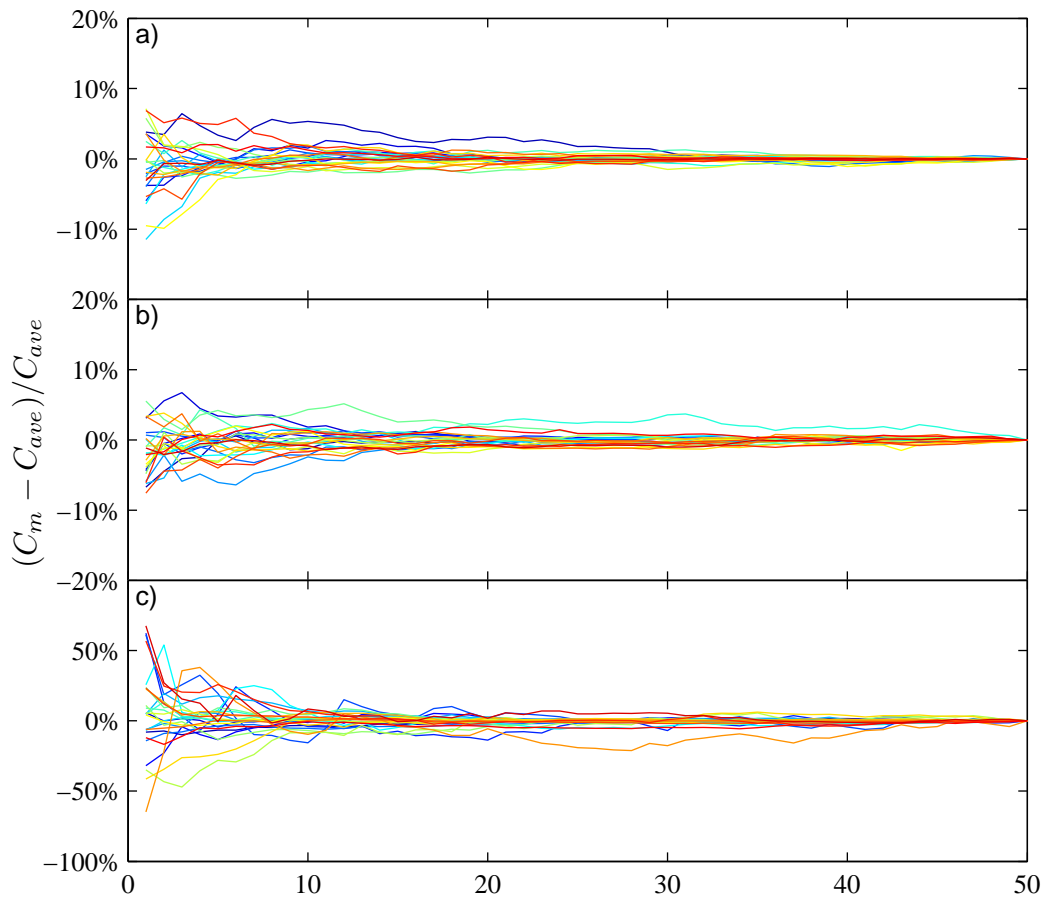
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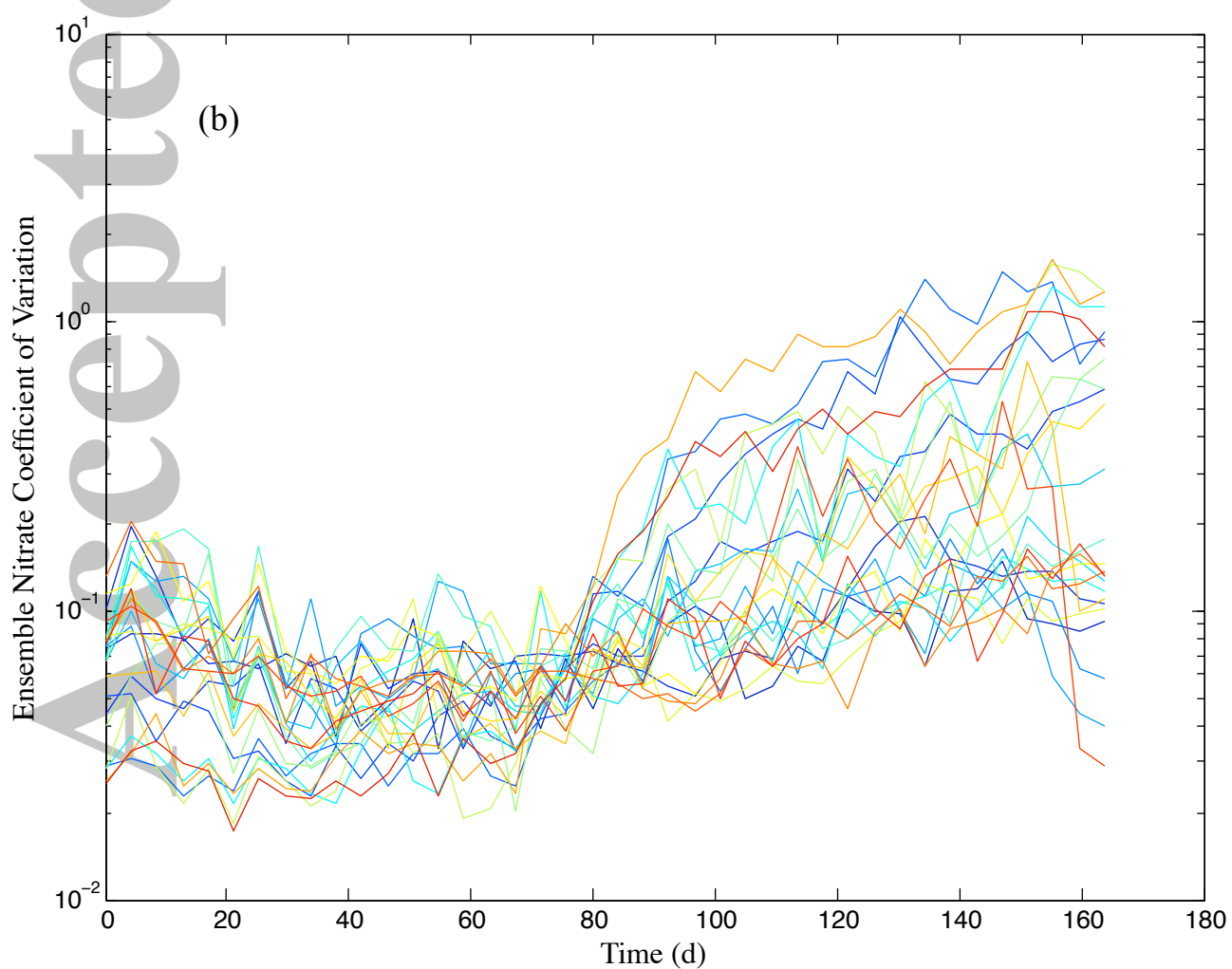
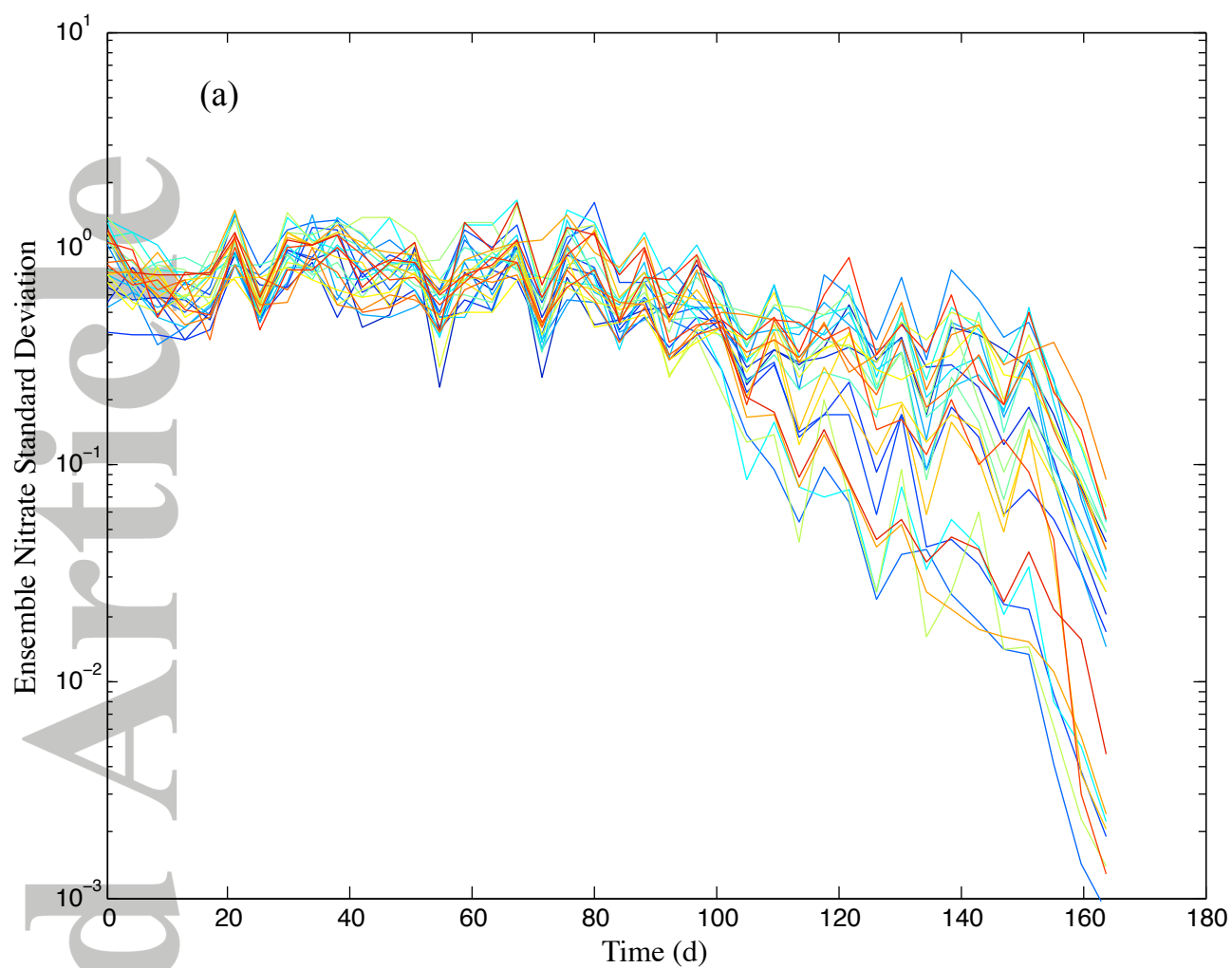
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