



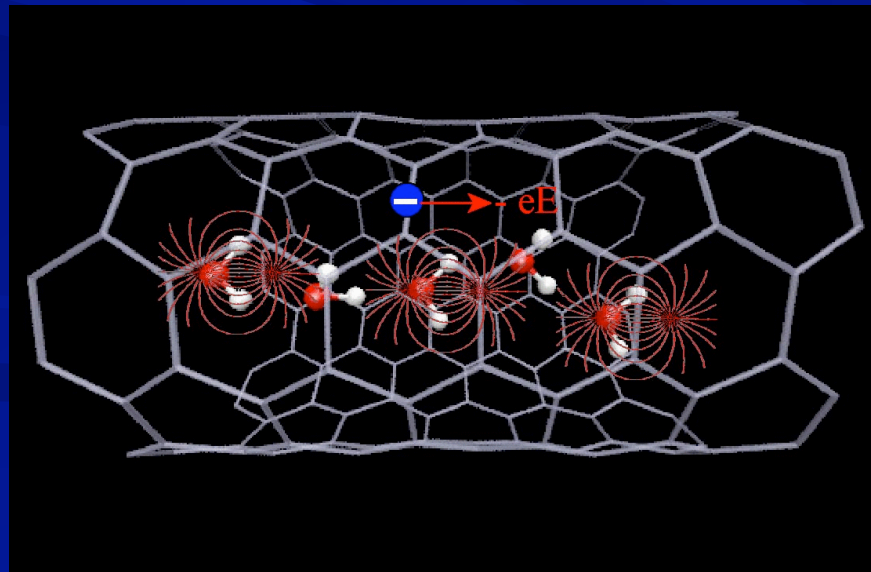
Electronic structure and dielectric behavior of finite-length single-walled armchair carbon nanotubes

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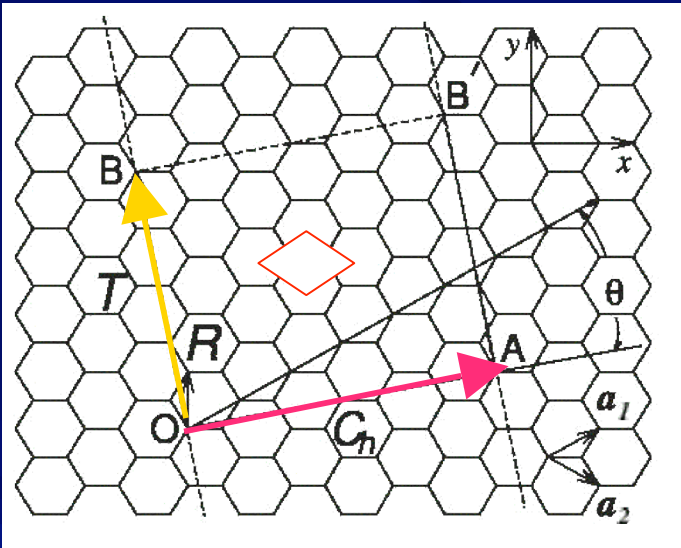
8/10/2004



DURINT

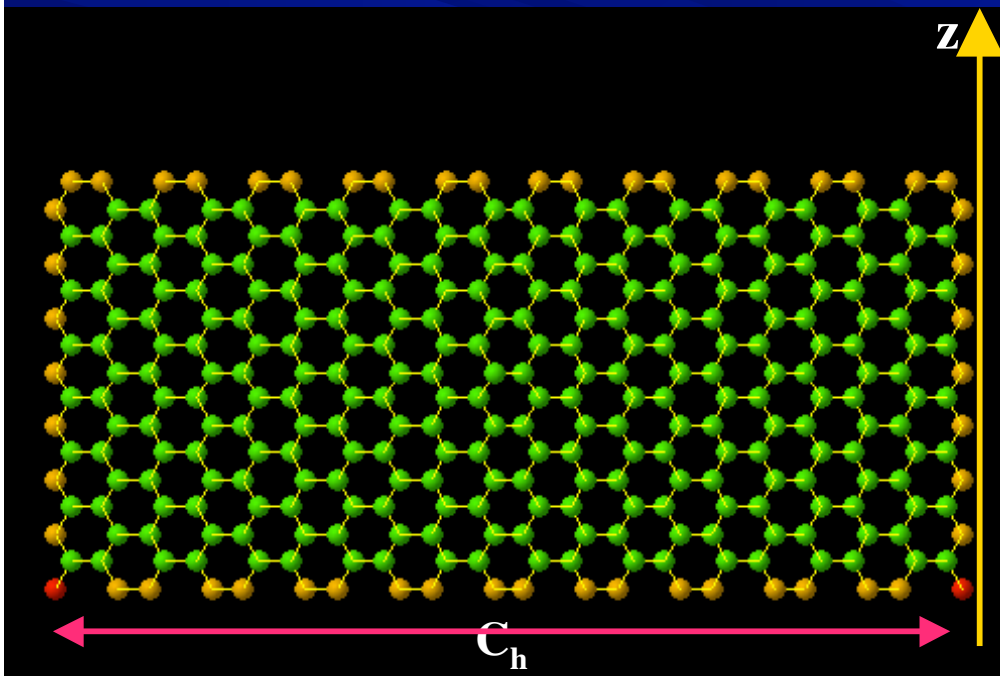


Carbon Nanotube

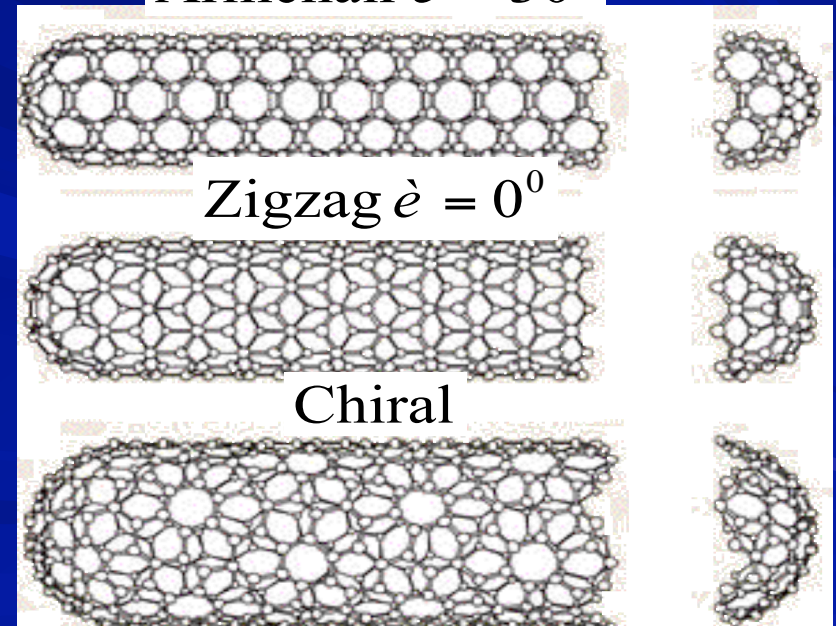


$$\vec{C}_h = n\vec{a}_1 + m\vec{a}_2 \text{ or } (n, m) \begin{cases} \text{metallic: } \text{mod}(n-m, 3) = 0 \\ \text{semiconducting: } \text{mod}(n-m, 3) = 1 \text{ or } 2 \end{cases}$$

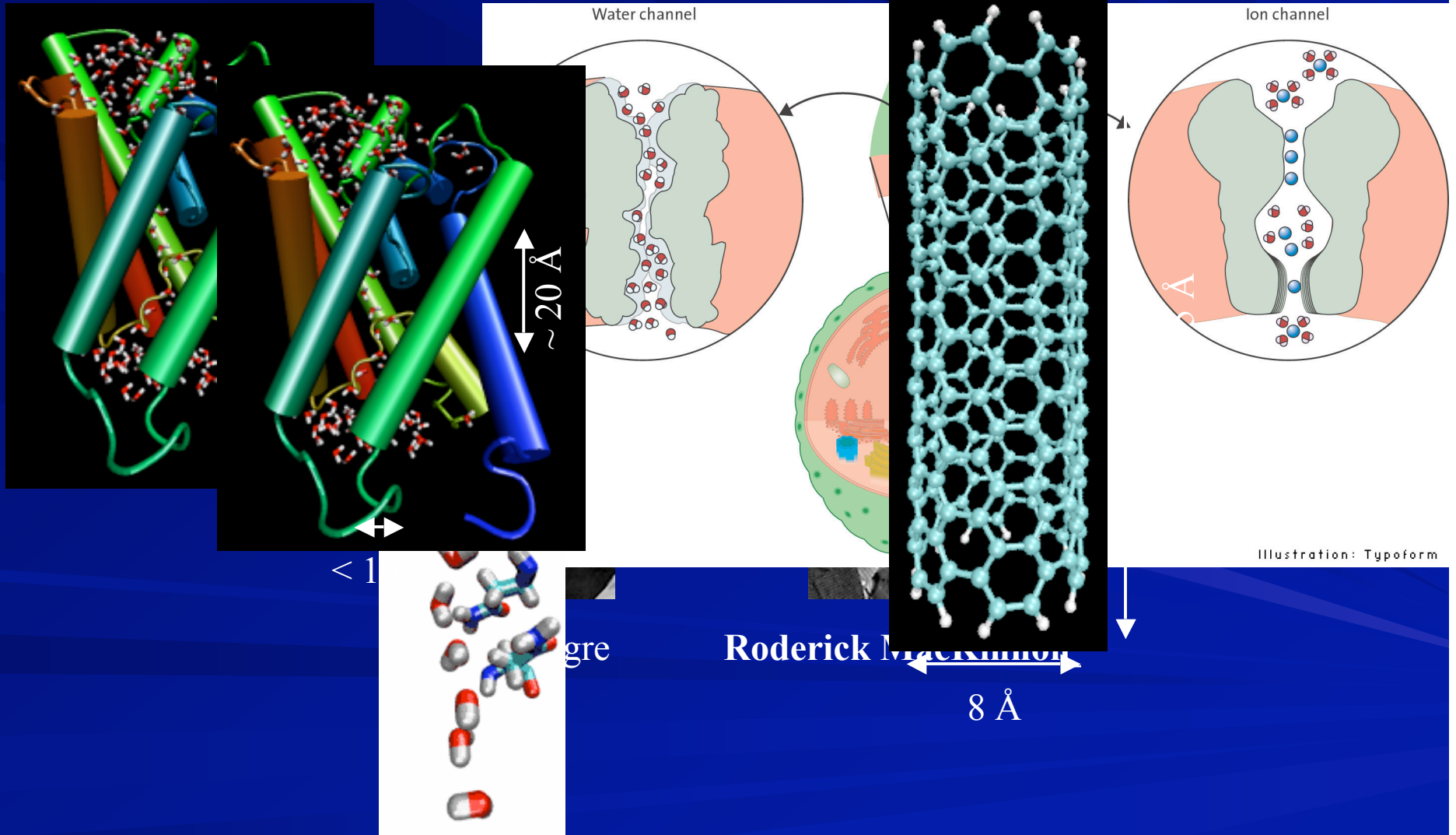
Rolling up a (10,10) nanotube



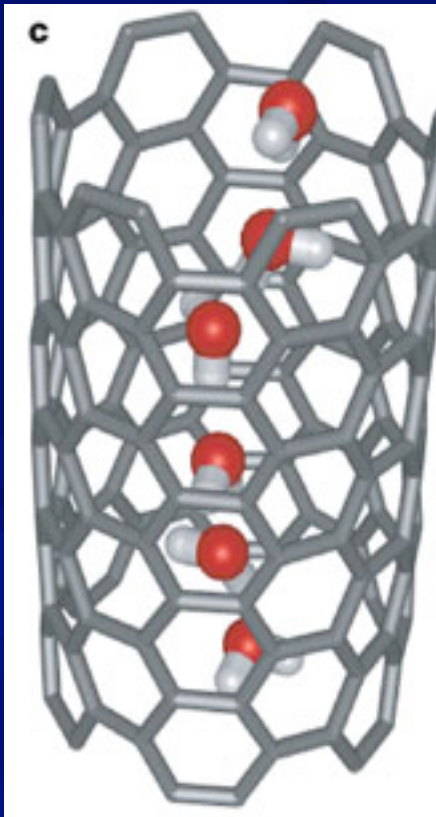
Armchair $\theta = 30^\circ$



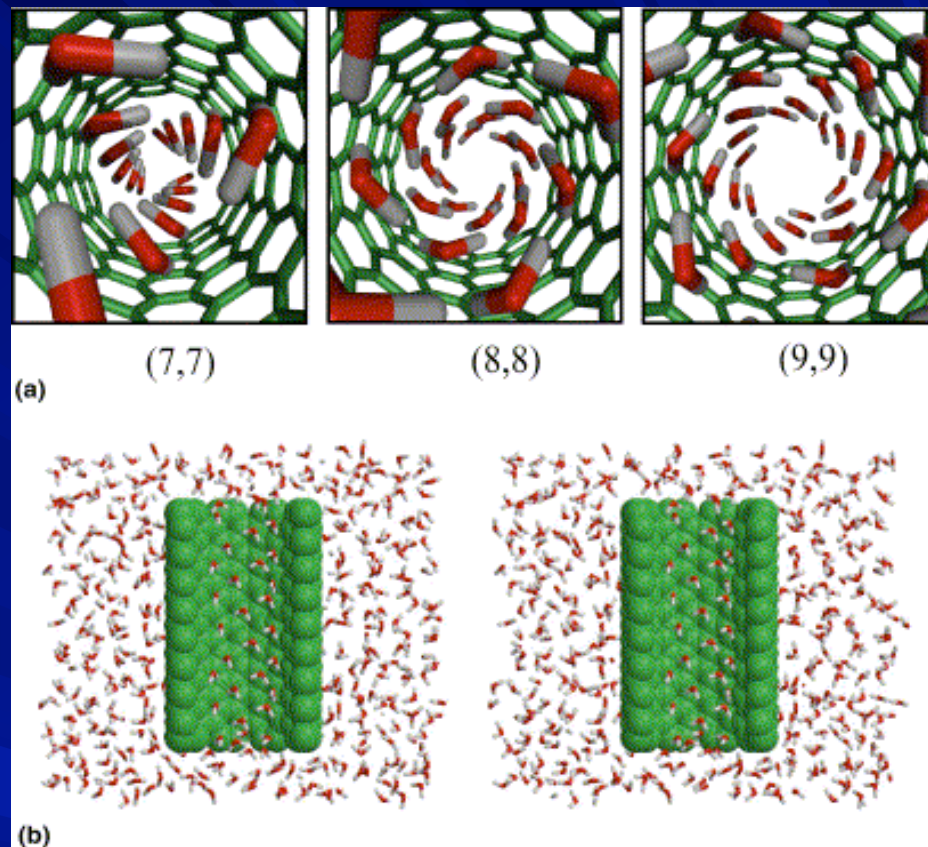
Nanotube & molecular channels



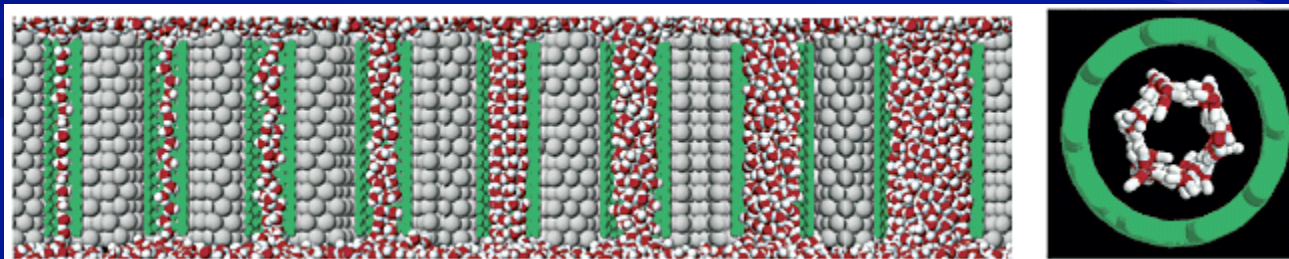
Nano-water



Hummer *et al.*, 2001

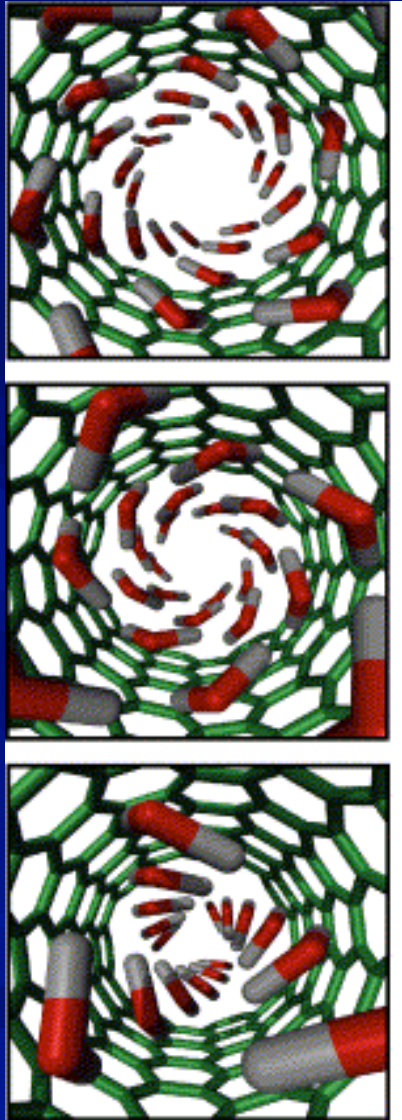


Noon *et al.*, 2002



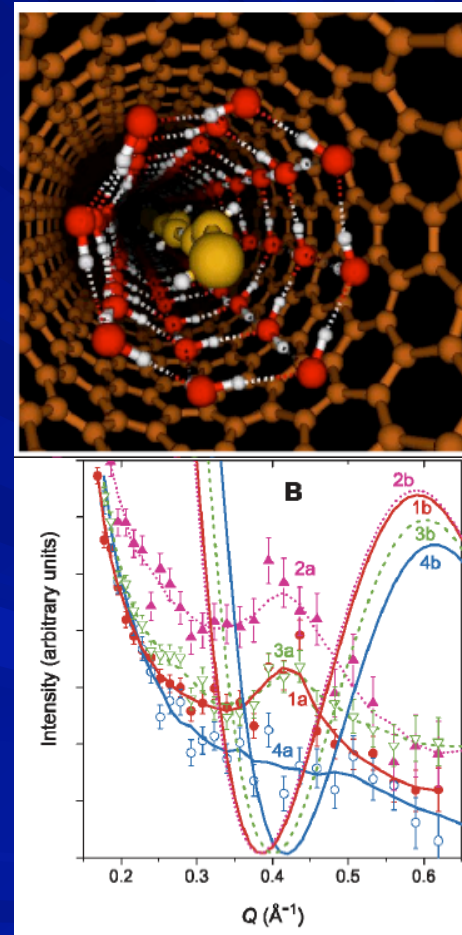
Mashl *et al.*, 2003

Nano-water



Noon *et al.*, 2002

simulation

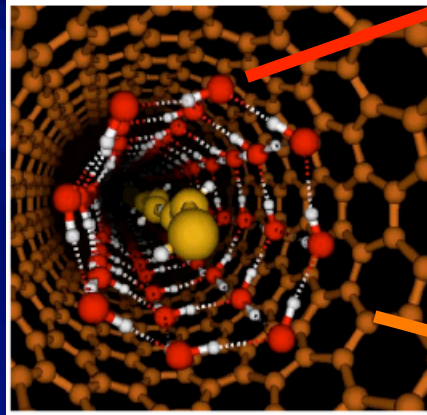


Kolesnikov *et al.*, 2004

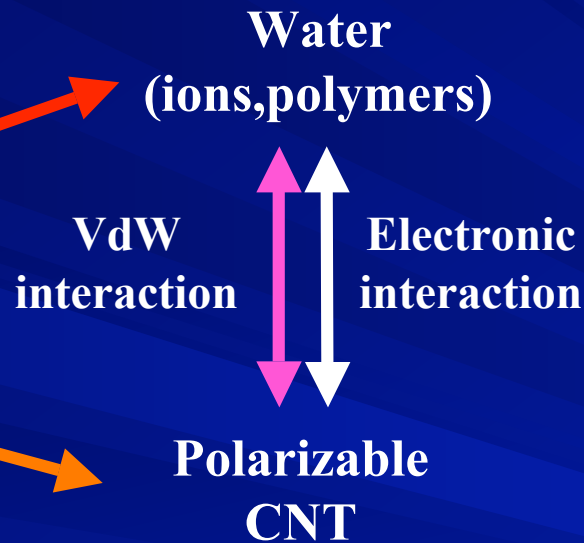
Experiment + simulation

Motivation: Model nano-water system

A polarizable model



Kolesnikov *et al.*, *PRL*, 2004

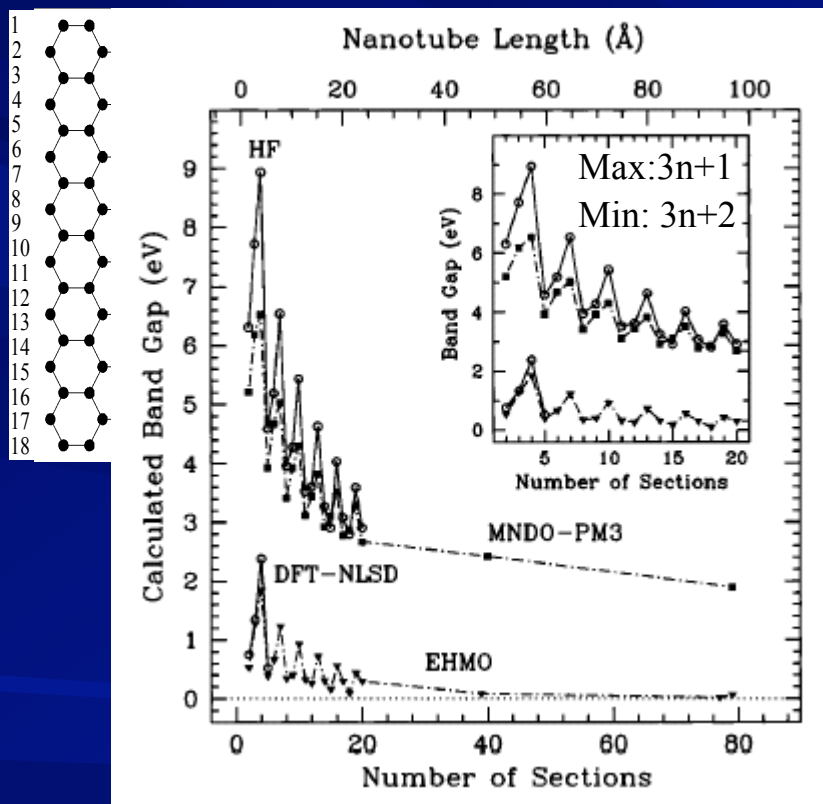


CNT does not only serve as a nm-size pore. The polarization effect of CNT may also modify the channel transport dynamics.

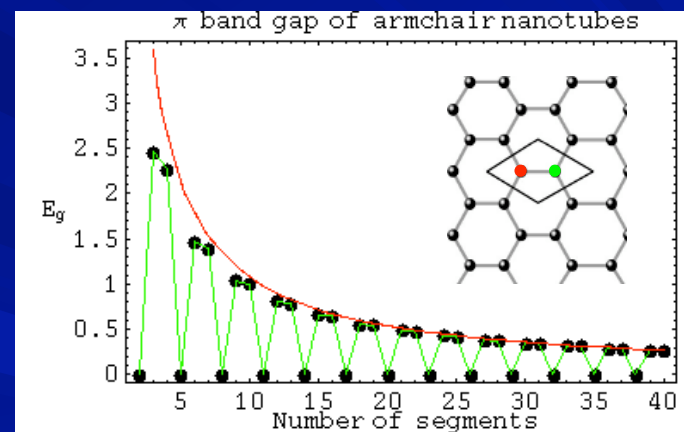
- Classic molecular dynamics: non-polarized CNT
- *ab initio* method: e.g. CPMD polarizable, but slow
- Develop a fast MD method including polarization
 - ❖ CNT polarization: Tight-binding model
 - ❖ Atomic motion: Classic force field

Band gap of finite armchair CNT

(6,6) armchair



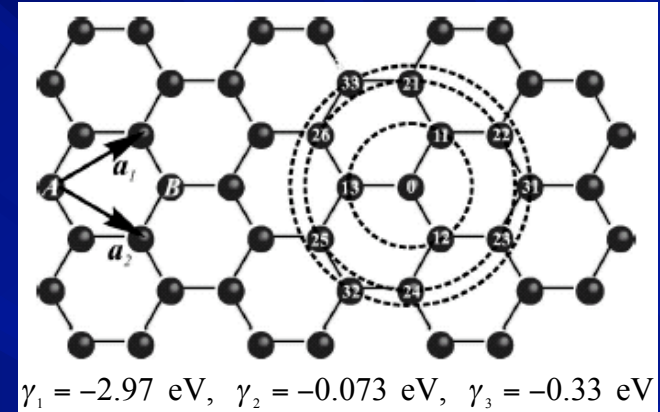
Particle in the box:



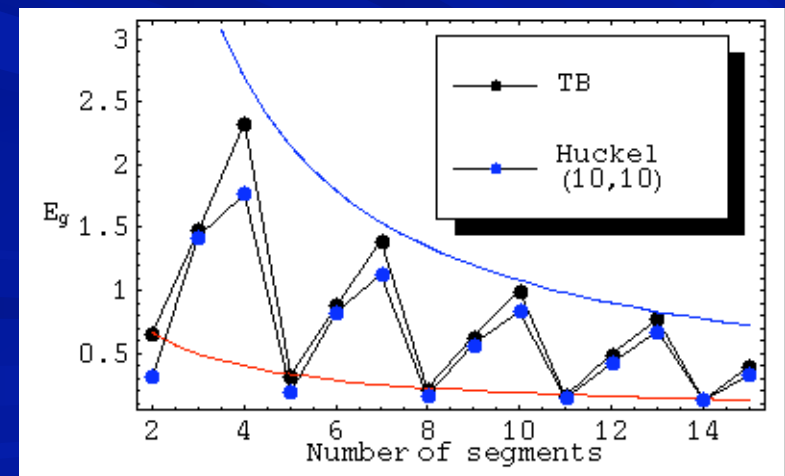
Rochefort, et al., *J. Phys. Chem. B*, 103:641-646, 1999.

Electronic properties of finite armchair CNT

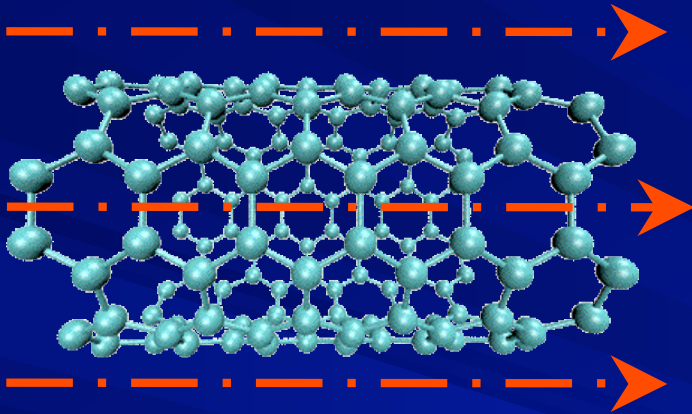
- Frontier wave functions are from π -orbitals
- 3rd n.n. hopping integral need to be included to break the **A-B degeneracy** at $3n+2$
- The 3rd n.n. TB calculation is able to generate the same band gap pattern as *ab initio* results
- Curvature effect is not large for HOMO/LUMO gap
- The frontier orbitals from the 3rd n.n. TB are identical with Huckel orbitals ($n>7$)



S. Reich, *et al.*, *PRB*,66,035412,2002.



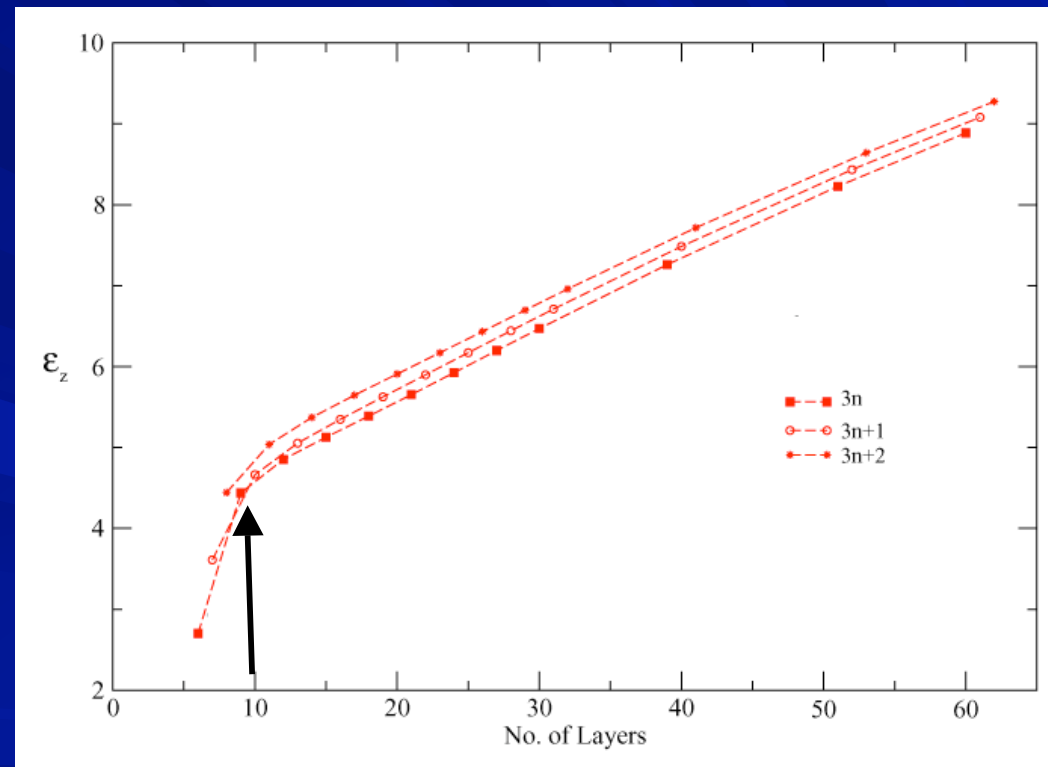
Axial Screening



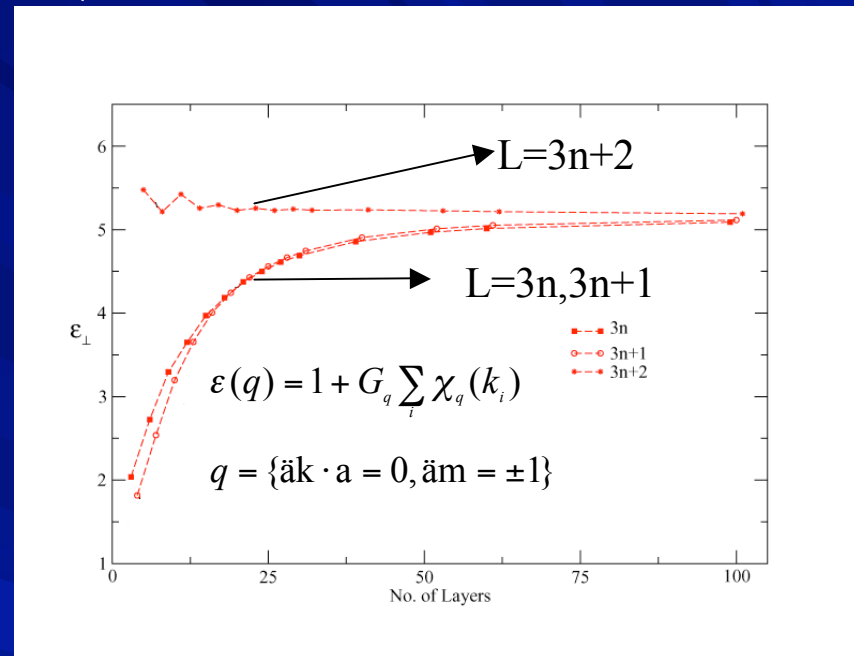
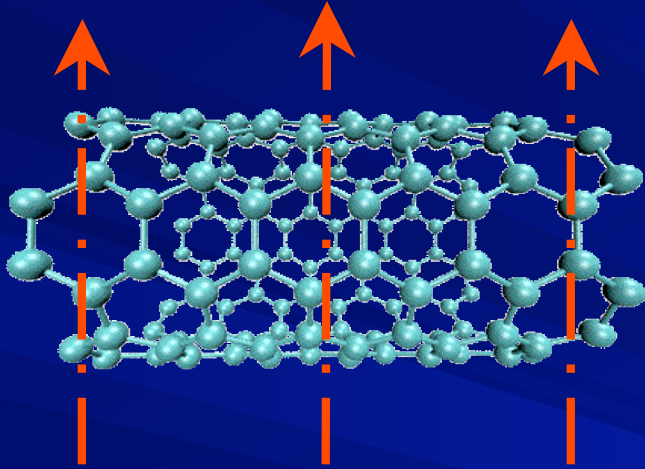
- Infinite long CNT: external field is well screened
- Finite length CNT: expect the screening effect becomes weaker

→ Define:

$$\epsilon_{\parallel} = \frac{V^{ext}}{V^{tot}} \Big|_{mid}$$

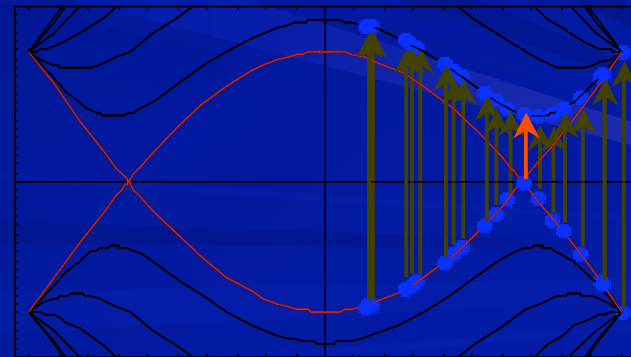


Perpendicular screening



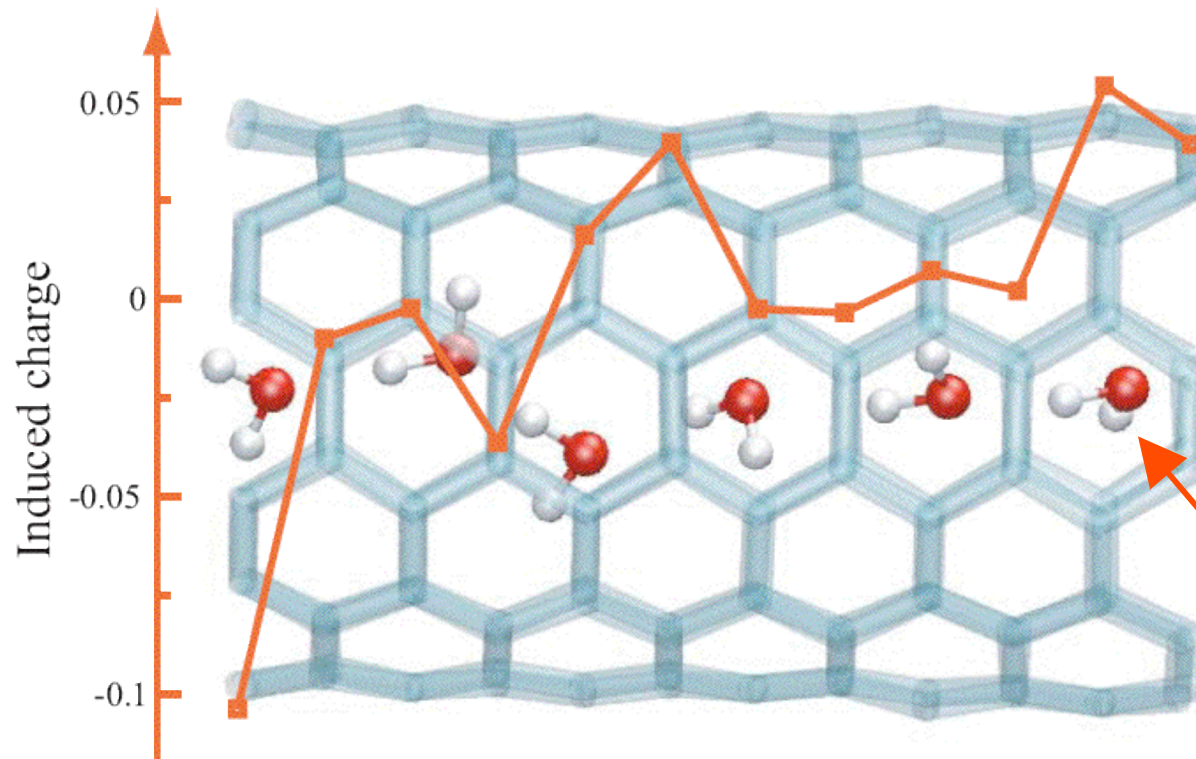
- $\epsilon_{\perp} \sim 5$ for infinite long CNT
- For finite length CNT, the screening is the strongest in the middle.

→ Define:
$$\epsilon_{\perp} = \left. \frac{V^{ext}}{V^{tot}} \right|_{mid}$$



Water chain in CNT channel

Induced charges along the axis of a (6,6)



Partial charge

H: 0.417

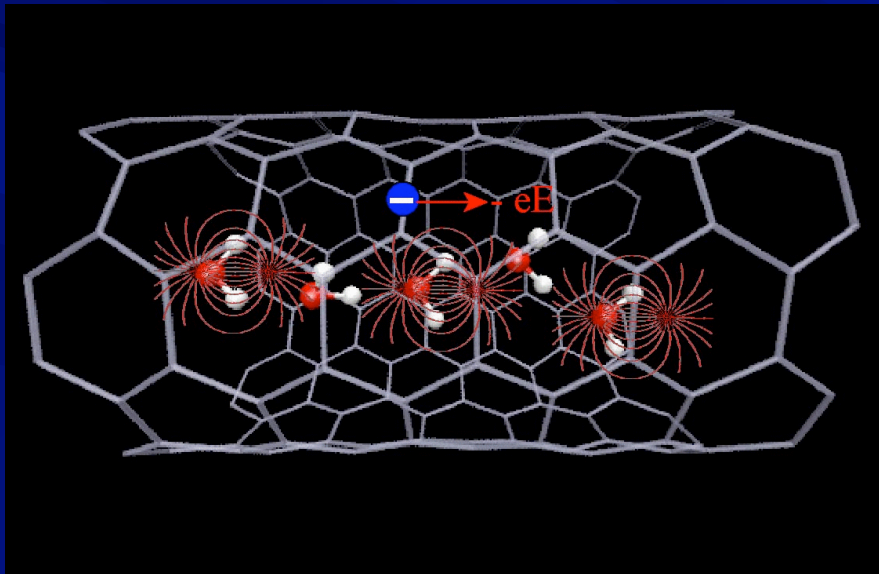
O: -0.834

water profile
from MD
simulation

Water chain in CNT channel

dipole moments of a water-filled

(6,6) CNT segment



H ₂ O wire length	μ_x	μ_y	μ_z	$ \mu $
6 (no CNT)	2.13	-0.54	-12.03	12.23
	-0.49	-0.79	16.07	16.09
6 (with CNT)	0.43	-0.12	-5.22	5.24
	0.05	-0.15	4.53	4.53

ab initio calculation: D. J. Mann and M.D. Halls, *PRL*, 2003

Summary

- **Electronic properties of finite length armchair CNT**
 - HOMO/LUMO gap oscillation, frontier orbitals
 - Break A/B degeneracy by γ_3
- **Factors affecting dielectric behavior**
 - Tube length, cross-over of 1D and 0D
 - Different subclasses ($l=3n, 3n+1, 3n+2$)
- **Induced charge distribution**
 - Screening of water dipoles (1/2), large induced Coulomb energy
- **Implement the polarizable carbon nanotube model into MD**
 - Molecular transport dynamics
 - Control channel by external field and functionalization