

1. Number of substates:

(a) (nsn'l) in LS coupling:

In this case, $L = l$ and there are only two possibilities 1L and 3L . The first has $2L + 1$ substates and the second has $(2S + 1)(2L + 1) = 3(2L + 1)$ substates. The total is $N = 4(2L + 1) \equiv 8l + 4$ substates.

(b) (nsn'l) in jj coupling:

There are two possibilities ($ns_{1/2}n'l_{l-1/2}$) with $2 \times (2(l-1/2)+1) = 4l$ substates and ($ns_{1/2}n'l_{l+1/2}$) with $2 \times (2(l+1/2)+1) = 4l + 4$ substates. The total is again $N = 8L + 4$ substates.

(c) $(nd)^2$ in LS coupling:

Since $L + S$ is even we have 1S , 3P , 1D , 3F , and 1G . The number of states is $N = 1 + 3 \times 3 + 5 + 3 \times 7 + 9 = 45$.

(d) $(nd)^2$ in jj coupling:

We have the jj configurations $(nd_{3/2})^2$, $(nd_{5/2})^2$, and $(nd_{3/2}nd_{5/2})$. For $(nd_{3/2})^2$ only $J = 0$ and 2 are possible with $1 + 5 = 6$ substates; for $(nd_{5/2})^2$ only $J = 0, 2$, and 4 are possible with $1 + 5 + 9 = 15$ substates; for $(nd_{3/2}nd_{5/2})$, $J = 1, 2, 3$, and 4 are possible with $3 + 5 + 7 + 9 = 24$ substates. The total number of substates is thus $N = 6 + 15 + 24 = 45$.

2. Cases (a-d) are normally ordered and have core expectation values 0. Cases (e) and (f) are not normally ordered. The core expectation values for case (e):

$$\langle a_c^\dagger a_d^\dagger a_a a_b \rangle = \delta_{bc} \delta_{ad} - \delta_{bd} \delta_{ac},$$

and for case (f):

$$\langle a_c^\dagger a_b a_d^\dagger a_c \rangle = \delta_{bc} \delta_{dc}.$$

3. We have $|I\rangle = a_a |0_c\rangle$ and $|F\rangle = a_m^\dagger a_b a_c |0_c\rangle$. It follows that

$$\begin{aligned} \langle F|V_I|I\rangle &= \frac{1}{2} \sum_{ijkl} g_{ijkl} \langle 0_c | a_c^\dagger a_b^\dagger a_m : a_i^\dagger a_j^\dagger a_l a_k : a_a | 0_c \rangle \\ &\equiv \frac{1}{2} \sum_{ijkl} g_{ijkl} \langle 0_c | : a_c^\dagger a_b^\dagger a_m : : a_i^\dagger a_j^\dagger a_l a_k : : a_a : | 0_c \rangle \end{aligned}$$

which, using Wick's theorem, reduces to

$$\begin{aligned} &= \frac{1}{2} \sum_{ijkl} g_{ijkl} (\delta_{ia} \delta_{jm} - \delta_{im} \delta_{ja}) (\delta_{kb} \delta_{lc} - \delta_{lc} \delta_{lb}) \\ &= g_{ambc} - g_{amcb} = \tilde{g}_{ambc} \end{aligned}$$

4. The first-order energy is

$$E^{(1)}(j_v j_w, J) = \sum_{m's} - \begin{array}{c} \uparrow j_v m_v \\ \text{---} JM \\ \downarrow j_a m_a \end{array} - \begin{array}{c} \uparrow j_w m_w \\ \text{---} JM \\ \downarrow j_b m_b \end{array} \langle 0_c | a_b^\dagger a_w V_I a_v^\dagger a_a | 0_c \rangle,$$

where w & v and a & b differ only in the values of the magnetic quantum numbers. With the aid of Wick's theorem, we find

$$\langle 0_c | a_b^\dagger a_w V_I a_v^\dagger a_a | 0_c \rangle = g_{w a b v} - g_{w a v b} + \Delta_{w v} \delta_{b a} - \Delta_{b a} \delta_{w v}$$

The two-particle part of the first-order energy is

$$\sum_{m's} - \begin{array}{c} \uparrow j_v m_v \\ \text{---} JM \\ \downarrow j_a m_a \end{array} - \begin{array}{c} \uparrow j_w m_w \\ \text{---} JM \\ \downarrow j_b m_b \end{array} \sum_k \left[- \begin{array}{c} \uparrow j_w m_w \quad \uparrow j_a m_a \\ \text{---} k \\ \downarrow j_b m_b \quad \downarrow j_v m_v \end{array} + X_k(w a b v) \right. \\ \left. - \begin{array}{c} \uparrow j_w m_w \quad \uparrow j_a m_a \\ \text{---} k \\ \downarrow j_v m_v \quad \downarrow j_b m_b \end{array} + X_k(w a v b) \right]$$

After summing over m values, this reduces to

$$\frac{(-1)^{J+j_v-j_a}}{[J]} \left[X_J(w a b v) + [J] \sum_k \left\{ \begin{array}{ccc} j_v & j_a & J \\ j_b & j_w & k \end{array} \right\} X_k(w a v b) \right].$$

This leads to the answer on setting $w = v$ and $b = a$.