

# DICE simplified\*

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**Abstract:** We analyze Nordhaus' DICE model and show that the temperature and CO<sub>2</sub> equations are needlessly complicated and can be simplified without loss of essence. In addition, we argue that the damage function can be altered in such a way that it lends itself to experiments involving extreme risk. We conclude that, within the philosophy of the DICE model, significant simplifications can be made which make the model more transparent, more robust, and easier to apply.

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## 1 Introduction

This paper provides a critical assessment of the 2016 DICE (Dynamic Integrated model of Climate and the Economy) model originally developed by Nordhaus (1992), but since then continuously updated and altered (Nordhaus, 2018b). Integrated assessment models (IAMs) are being used extensively for the analysis of climate change policy and DICE has played an important part in projecting greenhouse gas emissions and temperature under various social and economic scenarios.

The assumptions and functional forms used in DICE have been under considerable scrutiny and criticism. Among others, Pindyck (2013, 2017) and Gerlagh and Liski (2018a, 2018b) argue that IAMs should be made simpler and more transparent, and an editorial in *Nature Climate Change* (Editorial, 2015) also emphasizes this view. DICE is complex, because it employs an equation for radiative forcing, two equations (a two-box model) for the climate system, and a three-reservoir model for the carbon cycle. Some of the other IAMs (for example, the FUND and PAGE models) are simpler, as they employ a single-equation climate model (Calel and Stainforth, 2017).

DICE includes a complex geophysical system, perhaps too complex. The purpose of this paper is to show that the temperature and CO<sub>2</sub> equations are needlessly complicated and can be much simplified. The reason why this simplification is possible lies in the fact that the two dynamic equations describing temperature in the DICE model have one approximate unit eigenvalue, and that the three equations describing CO<sub>2</sub> have one exact and one approximate unit eigenvalue. We derive the exact counterparts after differencing out the auxiliary variables and we provide simplifications.

In response to criticism that IAMs give the impression of being ‘black boxes’ (Editorial, 2015), some studies propose a simpler IAM in which an analytical formula for the optimal carbon price can be derived; see Golosov *et al.* (2014), Traeger (2015), Rezai and Van der Ploeg (2016), Van den Bijgaart *et al.* (2016), Dietz and Venmans (2018), and Van der Ploeg (2018). Among studies employing closed-form IAMs, the specifications of the climate system and the carbon cycle vary. Golosov *et al.* (2014) specify a two-and-a-half-box carbon system consisting of a permanent component (about 20% of carbon) and a transient component. The climate system is omitted. Under assumptions such as logarithmic utility, Cobb-Douglas technology, constant saving rate, and full depreciation of capital, they derive the optimal-tax for-

mula analytically. Rezai and Van der Ploeg (2016) employ a similar two-box carbon system. They allow for a lag between temperature and atmospheric carbon in addition to more general functional forms than those assumed by Golosov *et al.* (2014), and they derive a simple rule for the optimal carbon price.

The omission of the climate system in Golosov *et al.* (2014) is justified by recent findings in climate science that the climate response to a CO<sub>2</sub> emission is nearly instantaneous and remains almost constant over time; see Roe and Bauman (2013) and Ricke and Caldeira (2014). Based on the same findings, Dietz and Venmans (2018) assume that the global mean temperature is linearly proportional to cumulative CO<sub>2</sub> emissions, in contrast to the large thermal inertia of the climate system assumed in DICE and Lemoine and Rudik (2017). These differences may lead to different optimal transition paths of temperature. Thus, the optimal carbon price follows Hotelling’s rule in Dietz and Venmans (2018), whereas it grows more slowly in Lemoine and Rudik (2017).

Our purpose is to stay as close as possible to Nordhaus’ DICE model, but to simplify it using Ockham’s razor quoted above (‘more things should not be used than are necessary’), also in the spirit of Einstein’s dictum: ‘as simple as possible but not simpler’. A complex model is not necessarily a better model than a simple model. Leaving things out is arguably more difficult and more important than putting things in. Many statisticians believe that a more complex model will reduce the bias and increase the variance, but this is only half true. A more complex model does indeed increase the variance, but it does not necessarily reduce the bias (De Luca *et al.*, 2018). Hence, simplicity matters given the large uncertainties in the exogenous variables (such as population and technical knowledge) and the parameters.

The purpose of this paper is therefore not to propose a completely different climate system nor to examine climate sensitivity in the DICE model. Our purpose is more modest. We show that the temperature and CO<sub>2</sub> equations in the DICE model are needlessly complicated and that they can be much simplified. We also argue that the specification of the damage function can be altered in such a way that it lends itself to experiments involving extreme risk. Finally, we briefly discuss the assumption in DICE that the abatement fraction for CO<sub>2</sub> is allowed to become larger than one, which implies that emissions can become negative.

In Section 2 we present the Nordhaus DICE 2016R model. In Sections 3 and 4 we discuss and simplify the DICE equations for temperature and CO<sub>2</sub> concentration. In Section 5 we provide an alternative to the DICE damage function. Section 6 presents the SICE (simplified DICE) model and concludes.

## 2 Nordhaus' DICE 2016R model

The following equations are the equations from the beta version of DICE-2016R (Nordhaus, 2017, 2018a), a version with the identification DICE-2016R-091916ap.gms. A number of equations are redundant and have been deleted. A new variable  $\omega_t$  has been introduced, some equations have been combined, and the equations have been reordered; see Ikefuji *et al.* (2019) for the details. Still, this is *precisely* the same model as Nordhaus' 2016R model.

Everybody works. In period  $t$ , the labor force  $L_t$  together with the capital stock  $K_t$  generate GDP  $Y_t$  through a Cobb-Douglas production function

$$Y_t = A_t K_t^\gamma L_t^{1-\gamma} \quad (0 < \gamma < 1), \quad (1)$$

where  $A_t$  represents technological efficiency and  $\gamma$  is the elasticity of capital. Capital is accumulated through

$$K_{t+1} = (1 - \delta)K_t + I_t \quad (0 < \delta < 1), \quad (2)$$

where  $I_t$  denotes investment and  $\delta$  is the depreciation rate of capital.

Carbon dioxide (CO<sub>2</sub>) emissions consist of industrial emissions and non-industrial ('land-use') emissions. We denote the latter type by  $E_t^0$  and consider it to be exogenous to our model. Total CO<sub>2</sub> emissions  $E_t$  are then given by

$$E_t = \sigma_t(1 - \mu_t)Y_t + E_t^0, \quad (3)$$

where  $\sigma_t$  denotes the emissions-to-output ratio for CO<sub>2</sub> and  $\mu_t$  is the abatement fraction for CO<sub>2</sub>. The associated CO<sub>2</sub> concentration increase  $M_t$  in the atmosphere (GtC from 1750) accumulates through

$$M_{t+1} = (1 - b_0)M_t + b_1 X_{1,t} + E_t, \quad (4a)$$

$$X_{1,t+1} = b_0 M_t + (1 - b_1 - b_3)X_{1,t} + b_2 X_{2,t}, \quad (4b)$$

$$X_{2,t+1} = b_3 X_{1,t} + (1 - b_2)X_{2,t}, \quad (4c)$$

where  $X_{1,t}$  and  $X_{2,t}$  are auxiliary variables representing CO<sub>2</sub> concentration increases in shallow and lower oceans, respectively, also measured in GtC from 1750.

Temperature increase  $H_t$  (degrees Celsius from 1900) develops according to

$$H_{t+1} = (1 - a_0)H_t + a_1 \log(M_{t+1}) + a_2 Z_t + F_{t+1}, \quad (5a)$$

$$Z_{t+1} = (1 - a_3)Z_t + a_3 H_t, \quad (5b)$$

where  $Z_t$  is an auxiliary variable representing the temperature increase of the lower oceans, also measured in degrees Celsius from 1900, and  $F_{t+1}$  is exogenous radiative forcing.

In each period  $t$ , the fraction of GDP not spent on abatement or ‘damage’ is either consumed ( $C_t$ ) or invested ( $I_t$ ) along the budget constraint

$$(1 - \omega_t - \xi H_t^2) Y_t = C_t + I_t. \quad (6)$$

A fraction  $\omega_t$  of  $Y_t$  is spent on abatement, and we specify the abatement cost fraction as

$$\omega_t = \psi_t \mu_t^\theta \quad (\theta > 1). \quad (7)$$

When  $\mu_t$  increases then so does  $\omega_t$ , and a larger fraction of GDP will be spent on abatement.

Damage is represented by a fraction  $\xi H_t^2$  of  $Y_t$  and it depends only on temperature. The optimal temperature is  $H_t = 0$ , the temperature in 1900. Deviations from the optimal temperature cause damage. For very high and very low temperatures the fraction becomes large, but (given the value of  $\xi$ ) it will still be a fraction between zero and one, unless in truly catastrophic cases.

As in Nordhaus (2017, 2018a) one period is five years. Period 1 refers to the time interval 2015–2019, period 2 to 2020–2024, and so on. Stock variables are measured at the beginning of the period; for example,  $K_1$  denotes capital in the year 2015. We choose the exogenous variables such that  $L_t > 0$ ,  $A_t > 0$ ,  $E_t^0 > 0$ ,  $\sigma_t > 0$ , and  $0 < \psi_t < 1$ . The policy variables must satisfy

$$C_t \geq 0, \quad I_t \geq 0, \quad \mu_t \geq 0. \quad (8)$$

Nordhaus (2018a, 2018b) allows ‘negative-emission technologies’ by setting an upper bound on  $\mu_t$  of 1.2 (rather than 1.0) from period 30 onwards (year 2160), which implies that emissions can become negative by (3), and in fact this upper bound is reached in the DICE output from period 46 onwards (year 2240). The idea of negative emissions is controversial. Anderson and Peters (2016) state that negative-emission technologies are unjust and a high-stake gamble, while a recent editorial in *Nature* (Editorial, 2018) discusses the enormous effort required to carry out such technologies — an effort which would lead to a deterioration of the environment.

Given a utility function  $U$  we define welfare in period  $t$  as

$$W_t = L_t U(C_t/L_t). \quad (9)$$

The policy maker has a finite horizon and maximizes total discounted welfare

$$W = \sum_{t=1}^T \frac{W_t}{(1 + \rho)^t} \quad (0 < \rho < 1), \quad (10)$$

where  $\rho$  denotes the discount rate and  $T = 100$  (500 years). Letting  $x$  denote per capita consumption, the utility function  $U(x)$  is assumed to be defined and strictly concave for all  $x > 0$ . There are many such functions, but a popular choice is

$$U(x) = \frac{x^{1-\alpha} - 1}{1-\alpha} \quad (\alpha > 0), \quad (11)$$

where  $\alpha$  denotes the elasticity of marginal utility of consumption. This is the so-called *power* function. Many authors, including Nordhaus, select this function. In earlier versions of the DICE model, Nordhaus (2008) chooses  $\alpha = 2$  in which case  $U(x) = 1 - 1/x$ . Also popular is  $\alpha = 1$ ; see Kelly and Kolstad (1999) and Stern (2007). In the 2016 version of the DICE model  $\alpha = 1.45$ .

### 3 Temperature

The DICE model thus consists of the seven equations (1)–(7). Four of these, equations (1)–(3) and (7), are not controversial. In the next three sections we shall discuss the CO<sub>2</sub> equation (4), the temperature equation (5), and the budget constraint (6).

We start with the temperature equations in (5), which we now write in matrix form as

$$x_{t+1} = Ax_t + a_{t+1}, \quad (12)$$

where

$$x_t = \begin{pmatrix} H_t \\ Z_t \end{pmatrix}, \quad A = \begin{pmatrix} 1 - a_0 & a_2 \\ a_3 & 1 - a_3 \end{pmatrix}, \quad a_t = \begin{pmatrix} a_1 \log(M_t) + F_t \\ 0 \end{pmatrix}.$$

The matrix  $A$  has two eigenvalues given by

$$1 - \frac{1 - \eta_1}{2} \pm \frac{1}{2} \sqrt{(1 - \eta_1)^2 - 4\eta_2},$$

where  $\eta_1 = 1 - a_0 - a_3 = 0.8468$  and  $\eta_2 = (a_0 - a_2)a_3 = 0.0030$ , so that the eigenvalues are 0.9771 and 0.8697, respectively. The largest eigenvalue is thus close to one and it would be equal to one if (and only if)  $\eta_2 = 0$ .

We can ‘difference out’ the auxiliary variable  $Z$  and this gives

$$\begin{aligned} H_{t+1} &= (1 + \eta_1)H_t - (\eta_1 + \eta_2)H_{t-1} \\ &\quad + \eta_3 \log(M_{t+1}) - (\eta_3 - \eta_4) \log(M_t) + \eta_{0t}, \end{aligned} \quad (13)$$

where  $\eta_3 = a_1 = 0.5338$ ,  $\eta_4 = a_1 a_3 = 0.0133$ , and  $\eta_{0t} = F_{t+1} - (1 - a_3)F_t$ . Equation (13) does not contain  $Z$  but, compared to (5a), it contains an

additional lag in both  $H$  and  $\log(M)$ . Note that (13) is not invariant to scaling in  $M$ .

Letting  $\Delta$  be the (backward) difference operator defined by  $\Delta x_{t+1} = x_{t+1} - x_t$ , we can write (13) alternatively as

$$\Delta H_{t+1} = \eta_1 \Delta H_t + \eta_3 \Delta \log(M_{t+1}) + \eta_{0t}^*,$$

where  $\eta_{0t}^* = -\eta_2 H_{t-1} + \eta_4 \log(M_t) + \eta_{0t}$ . This equation in first differences can also be written as

$$H_{t+1} = \eta_0 + \eta_1 H_t + \eta_3 \log(M_{t+1}) + \sum_{j=1}^t \eta_{0j}^*, \quad (14)$$

where  $\eta_0 = -3.3291$  is an integration constant. Notice that  $H_{t+1}$  in (14) depends on  $H_t$  but that the effect of  $H_{t-1}$  is negligible, which is another way of saying that the largest eigenvalue of the matrix  $A$  in (12) is close to one. Both are caused by the fact that  $\eta_2$  is small. We emphasize that (12), (13), and (14) are *equivalent* descriptions of the DICE temperature equations. No approximation has yet taken place.

Given the DICE parameter values, in particular the fact that  $\eta_2$  and  $\eta_4$  are small, the partial sums  $\sum_j \eta_{0j}^*$  are well approximated by a linear trend with slope 0.025. This implies that if we run a regression on the equation

$$H_{t+1} = \eta_0^* + \eta_1^* H_t + \eta_2^* \log(M_{t+1}) + \eta_3^* t \quad (15)$$

we will get a good fit. If we leave out the linear trend, then the estimate of  $\eta_1^*$  increases somewhat and the estimate of  $\eta_2^*$  decreases somewhat.

Table 1: Simplified temperature equations

	<i>constant</i>	$H_t$	$\log(M_{t+1})$	<i>trend/10</i>	$s$	$\max  Q_t $
(a)	-3.6772 (0.0511)	0.8707 (0.0014)	0.5826 (0.0076)	0.0064 (0.0003)	0.0032	0.43%
(b)	-2.8672 (0.0561)	0.8954 (0.0012)	0.4622 (0.0082)	—	0.0068	1.70%

More precisely, we obtain the results in Table 1, where we note that in all regressions, figures, and numerical experiments that follow,  $H_t$  (and similarly  $M_t$  and other variables) take the optimal values as obtained from the DICE gams routine, which optimizes welfare in (10).

Under (a) we report the estimated coefficients and standard errors from a regression of  $H_{t+1}$  on a constant,  $H_t$ ,  $\log(M_{t+1})$ , and a time trend, as in (15).



The fit is very good. In particular, letting  $e$  denote the vector of residuals and  $\hat{H}_t$  the predicted value of  $H_t$  from the regression, and defining the regression variance  $s^2$  and the relative deviations  $Q_t$  as

$$s^2 = e'e/(n - k), \quad Q_t = 100(\hat{H}_t - H_t)/H_t,$$

we find that  $s = 0.0032$  and  $\max_t |Q_t| = 0.43$  with  $n = 99$  and  $k = 4$ . This shows that for a temperature increase of, say, 3 degrees Celsius the maximum error will be 0.013 degrees.

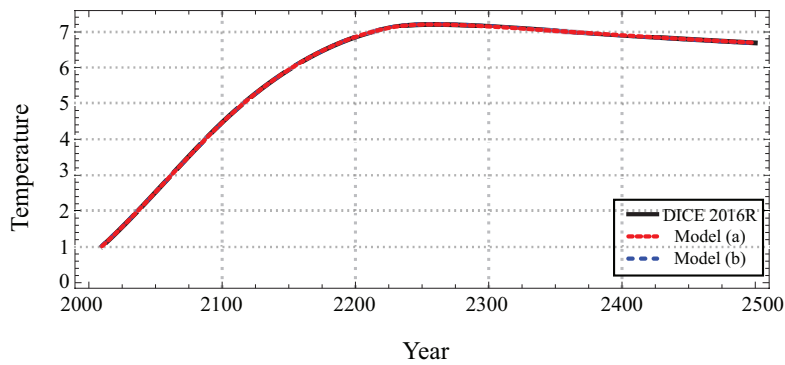


Figure 1: Temperature — time path for DICE 2016R and two simplified models

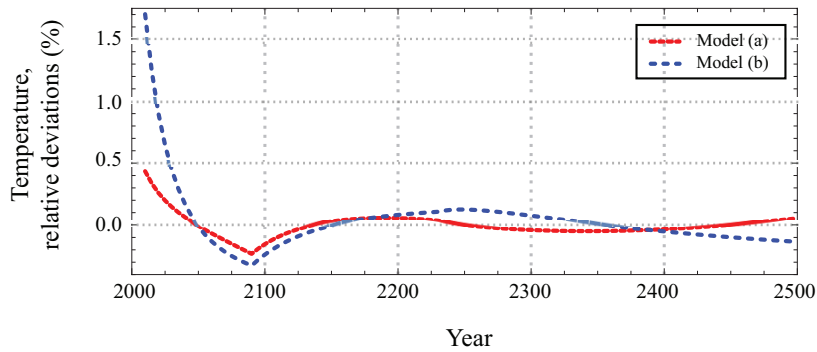


Figure 2: Temperature — Deviations (%) of two simplified models relative to DICE 2016R

If we leave out the linear trend we obtain (b) which is almost as good, except that the error in the first few periods is somewhat higher. This is illustrated in Figures 1 and 2. In Figure 1 the time paths of temperature

of DICE and the two models (a) and (b) are indistinguishable, reaching a maximum of 7.2 in 2270. The relative deviations  $Q_t$  are graphed in Figure 2. They are all below 0.5% except the first four periods in Model (b). Even though the estimated coefficient on the time trend is ‘significant’ it is not important, and the fit is essentially the same.

Summarizing, the simplified equation (15) provides a good approximation because (a) the coefficients in (13) correspond approximately to a first-order difference equation; and (b) the omitted variable is essentially constant. The second approximation (without trend) is almost as good as the approximation in (15), and suffices for practical applications.

## 4 CO<sub>2</sub> concentration

Next we consider the CO<sub>2</sub> equations in (4), which we also write in matrix form as

$$x_{t+1} = Ax_t + a_t, \quad (16)$$

where now

$$x_t = \begin{pmatrix} M_t \\ X_{1,t} \\ X_{2,t} \end{pmatrix}, \quad A = \begin{pmatrix} 1 - b_0 & b_1 & 0 \\ b_0 & 1 - b_1 - b_3 & b_2 \\ 0 & b_3 & 1 - b_2 \end{pmatrix}, \quad a_t = \begin{pmatrix} E_t \\ 0 \\ 0 \end{pmatrix}.$$

One of the three eigenvalues of  $A$  equals one, and the remaining eigenvalues are given by

$$1 - \frac{1 - \phi_1}{2} \pm \frac{1}{2} \sqrt{(1 - \phi_1)^2 - 4\phi_2},$$

where  $\phi_1 = 1 - b_0 - b_1 - b_2 - b_3 = 0.675535$  and  $\phi_2 = b_0b_2 + b_0b_3 + b_1b_2 = 0.001303$ , so that the two remaining eigenvalues take the values 0.995933 and 0.679602, respectively. The largest eigenvalue is thus equal to one and the next eigenvalue is close to one; it would be equal to one if (and only if)  $\phi_2 = 0$ . This suggests that we should difference not once (as in the previous section) but twice, and this is precisely what we shall do.

As in the previous section we can ‘difference out’ the auxiliary variables  $X_1$  and  $X_2$ , and this gives

$$M_{t+1} = (\phi_1 + 2)M_t - (1 + 2\phi_1 + \phi_2)M_{t-1} + (\phi_1 + \phi_2)M_{t-2} + E_t^*, \quad (17)$$

where

$$E_t^* = E_t - (1 + \lambda_1)E_{t-1} + (\lambda_1 + \lambda_2)E_{t-2}.$$

This equation does not contain  $X_1$  and  $X_2$  but it contains two additional lags in both  $M$  and  $E$ . We can write (17) alternatively as

$$\Delta M_{t+1} = (\phi_1 + 1)\Delta M_t - (\phi_1 + \phi_2)\Delta M_{t-1} + E_t^*,$$

where we notice that there is no remainder term in  $M_{t-2}$  because the largest eigenvalue of  $A$  equals one exactly given the DICE parameters. This leads to

$$M_{t+1} = \phi_0 + (\phi_1 + 1)M_t - (\phi_1 + \phi_2)M_{t-1} + E_t^{**}, \quad (18)$$

where  $E_t^{**} = \sum_{j=1}^t E_j^*$  and  $\phi_0 = 0.8761$  is an integration constant. This, in turn, can be written as

$$\Delta M_{t+1} = \phi_0 + \phi_1\Delta M_t - \phi_2M_{t-1} + E_t^{**},$$

so that

$$M_{t+1} = \phi_{00} + \phi_0 t + \phi_1 M_t - \phi_2 \sum_{j=1}^{t-1} M_j + \sum_{j=1}^t w_{tj} E_j, \quad (19)$$

where  $\phi_{00} = 263.2837$  is an integration constant and

$$\begin{aligned} \sum_{j=1}^t w_{tj} E_j &= \sum_{j=1}^t E_j^{**} = \sum_{j=1}^t \sum_{i=1}^j E_i^* = \sum_{j=1}^t (t-j+1) E_j^* \\ &= E_t + (1-\lambda_1) \sum_{j=1}^{t-1} E_j + \lambda_2 \sum_{j=1}^{t-2} (t-j-1) E_j, \end{aligned}$$

so that

$$w_{tt} = 1, \quad w_{tj} = 1 - \lambda_1 + (t-j-1)\lambda_2 \quad (j = 1, \dots, t-1).$$

The DICE weights  $w_{tj}$  are thus slightly *increasing* rather than decreasing, which is a little awkward. Notice that equations (16)–(19) are *equivalent* descriptions of the DICE CO<sub>2</sub> equations. No approximation has yet taken place.

Since  $\phi_2 = 0.0013$  and  $\lambda_2 = 0.0003$  are close to zero, (19) will be well approximated by

$$M_{t+1} \approx \phi_{00} + \phi_0 t + \phi_1 M_t + E_t + \sum_{j=1}^{t-1} (1 - \lambda_1) E_j.$$

In fact we will run regressions on the equation

$$M_{t+1} = \phi_0^* + \phi_1^* M_t + \phi_2^* E_t + \phi_3^* t \quad (20)$$

Table 2: Simplified CO<sub>2</sub> equations

	<i>constant</i>	$M_t$	$E_t$	<i>trend</i>	$s$	$\max  Q_t $
(a)	16.27 (1.08)	0.9900 (0.0004)	0.6166 (0.0067)	0.0317 (0.0098)	1.52	0.35%
(b)	17.94 (0.99)	0.9902 (0.0004)	0.6001 (0.0046)	—	1.59	0.39%
(c)	—	0.9975 (0.0002)	0.6549 (0.0073)	—	3.33	0.92%
(d)	—	0.9942 (0.0007)	1.0000 (—)	—	16.35	1.66%

and simplifications thereof.

This leads to the results in Table 2. Under (a) we regress  $M_{t+1}$  on all four variables, under (b) we delete the trend, under (c) we also delete the constant term, and under (d) we restrict the coefficient of  $E_t$  to be one. The last model is the simplest and mirrors capital accumulation in (2). The fit is very good in all cases, as can be seen from the values of  $s$  and  $\max |Q_t|$ , and also from Figures 3 and 4.

In Figure 3 the time paths of CO<sub>2</sub> of DICE and the four models (a)–(d) are indistinguishable, reaching a maximum of 2707 in 2230. The relative deviations  $Q_t$  are graphed in Figure 4. In Models (a) and (b) the relative deviations are all below 0.4% in absolute value. In Model (c) the relative deviations are all below 0.6% in absolute value, except in the first three periods. In Model (d) the relative deviations are larger than 0.6% up to period 36 (year 2190) and smaller than 0.6% afterwards, with a maximum of 1.7% in period 12 (year 2070) where  $M_t = 1402$  (the DICE output) and  $\hat{M}_t = 1425$  (the predicted value of  $M_t$  from the regression).

The simplified equation (20) thus provides an excellent approximation to the DICE results because the coefficients in (17) correspond approximately to a second-order difference equation. Model (b) is possibly the preferred approximation although the simplest Model (d) will suffice for most practical applications.

## 5 Damage and abatement

The damage-abatement function is DICE specifies two fractions,  $\omega_t$  (abatement) and  $\xi H_t^2$  (damage), of  $Y_t$  which reduce  $Y_t$  so that less money is available for investment and consumption along the budget constraint. In DICE this

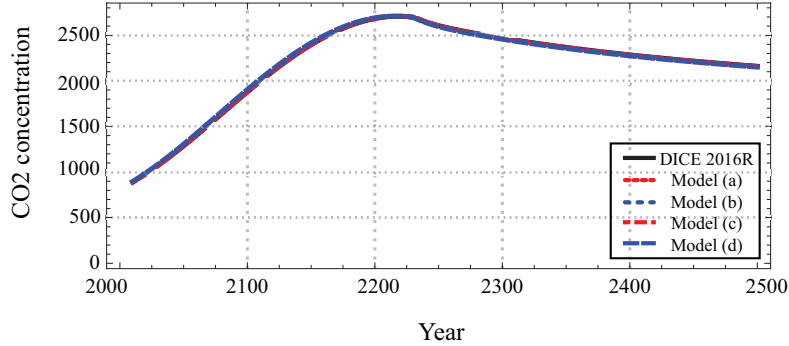


Figure 3: CO<sub>2</sub> concentration — time path for DICE 2016R and four simplified models

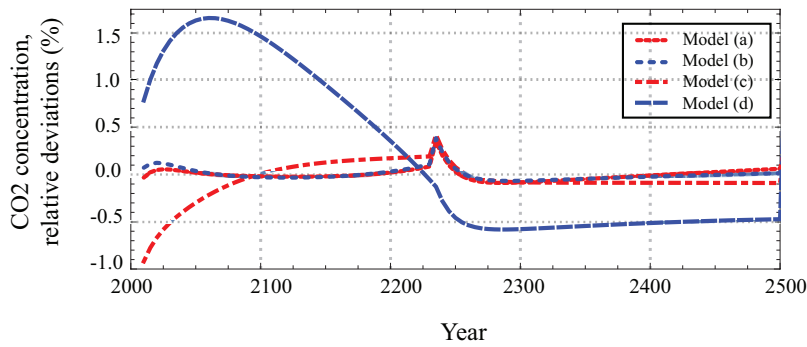


Figure 4: CO<sub>2</sub> concentration — Deviations (%) of four simplified models relative to DICE 2016R

fraction is specified as

$$1 - \omega_t - \xi H_t^2.$$

For very high and very low temperatures the fraction becomes large, but (given the value of  $\xi$ ) it will still be a fraction between zero and one, unless in truly catastrophic cases when  $H_t > 20.58$ , that is, when the temperature in period  $t$  is more than 20 degrees Celsius higher than in 1900. Of course, other forms of the damage function are possible; see Stern (2007), Weitzman (2009), Ackerman *et al.* (2010), and Nordhaus (2013). Howard and Sterner (2017) emphasize the importance of the damage function in accurately estimating coefficient and standard error bias.

In Figure 5 the graph labeled DICE 2016R contains the time path of this

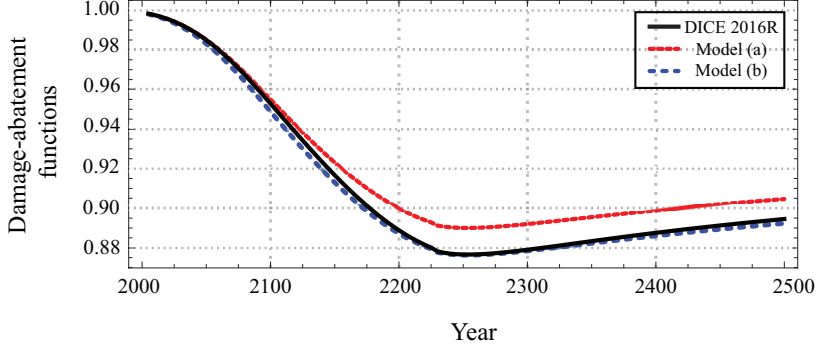


Figure 5: Damage-abatement functions

fraction. Models (a) and (b) use an alternative specification, namely

$$\frac{1 - \omega_t}{1 + \xi H_t^2},$$

based on the fact that the difference

$$D_t = 1 - \omega_t - \xi H_t^2 - \frac{1 - \omega_t}{1 + \xi H_t^2} = \frac{-(\omega_t + \xi H_t^2)(\xi H_t^2)}{1 + \xi H_t^2}$$

is small, about 1% in relative terms.

The alternative specification is of interest because we may wish to randomize  $\xi$ , as in Ikefuji *et al.* (2019). This is difficult under the DICE specification because  $1 - \omega_t - \xi H_t^2$  could become negative (under extreme circumstances), while the alternative specification is always positive provided  $\xi > 0$ .

In Model (a) we use the same value for  $\xi$  as in the DICE model, while in Model (b) we use an ‘optimal’ value  $\xi^* = 0.00265$  which brings the lines closer together; in fact, DICE and Model (b) are indistinguishable in the figure. The value  $\xi^*$  is obtained by minimizing the sum of squares

$$\sum_{t=1}^T \left( 1 - \omega_t - \xi H_t^2 - \frac{1 - \omega_t}{1 + \xi^* H_t^2} \right)^2.$$

with respect to  $\xi^*$ . In summary, for a suitable choice of  $\xi$  we obtain an alternative for the DICE damage function which lends itself better to studying situations of uncertainty or catastrophe.

## 6 The SICE model

We summarize our proposed SICE (simplified DICE) model with the relevant parameters. The SICE model is the DICE model, but with the temperature equation replaced by the new (simplest) temperature equation

$$H_{t+1} = \eta_0^* + \eta_1^* H_t + \eta_2^* \log(M_{t+1}) \quad (21)$$

with

$$\eta_0^* = -2.8672, \quad \eta_1^* = 0.8954, \quad \eta_2^* = 0.4622,$$

and the CO<sub>2</sub> equation replaced by

$$M_{t+1} = \phi_1^* M_t + E_t, \quad \phi_1^* = 0.9942. \quad (22)$$

For the damage equation we propose

$$\frac{1 - \omega_t}{1 + \xi^* H_t^2}, \quad \xi^* = 0.00265, \quad (23)$$

instead of the DICE specification

$$1 - \omega_t - \xi H_t^2, \quad \xi = 0.00236.$$

In addition, one may wish to set  $\mu \leq 1.0$  instead of the upper bound  $\mu \leq 1.2$  as used in DICE.

We conclude that, within the philosophy of Nordhaus' DICE model, significant simplifications can be made which make the model more transparent, more robust, and easier to apply.

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## Appendix

In this Appendix we present three tables which together contain all variable and parameter definitions required to compute the optimum in DICE (the Nordhaus model) and SICE (our simplified version of DICE): the variables employed in SICE and DICE and their relationship (Table 3), the initial values of the six state variables of DICE (Table 4), and the parameters employed in SICE and DICE and their relationship (Table 5).

Table 3: Variables in SICE and DICE

SICE	DICE
<i>State variables</i>	
$H_t$	TATM(t)
$K_t$	K(t)
$M_t$	MAT(t)
<i>Policy variables</i>	
$I_t$	$tstep \times I(t)$
$\mu_t$	MIU(t)
$C_t$	$tstep \times C(t)$
<i>Exogenous variables</i>	
$A_t$	$tstep \times AL(t)$
$E_t^0$	$tstep \times ETREE(t)/3.666$
$F_t$	$-a_1 \log(588) + 0.1005 \times FORCOTH(t)$
$L_t$	L(t)/1000
$\psi_t$	COST1(t)
$\sigma_t$	SIGMA(t)/3.666
<i>Auxiliary variables</i>	
$E_t$	$tstep \times E(t)/3.666$
$Y_t$	$tstep \times YGROSS(t)$
$\omega_t$	defined in (7)
$X_{1,t}$	MU(t)
$X_{2,t}$	ML(t)
$Z_t$	TOCEAN(t)

Table 4: Initial values of state variables in SICE and DICE

$H_1$	$K_1$	$M_1$	$X_{1,1}$	$X_{2,1}$	$Z_1$
0.85	223	851	460	1740	0.0068

Table 5: Parameters in SICE and DICE

SICE	DICE	Value	Description
<i>Basic parameters</i>			
	<i>a1</i>	0	
$\xi$	<i>a2</i>	0.00236	
	<i>a3</i>	2	
	<i>tstep</i>	5	
	<i>dk</i>	0.1	
$\delta$		0.40951	$1 - \delta = (1 - dk)^{tstep}$
	<i>prstp</i>	0.015	
$\rho$		0.077284	$1 + \rho = (1 + prstp)^{tstep}$
$\theta$	<i>expcost2</i>	2.6	
$\gamma$	<i>gama</i>	0.3	
$\alpha$	<i>elasmu</i>	1.45	
$a_0$		0.128189	$a_0 = (c1)(fco22x)/(t2xco2) + (c1)(c3)$
$a_1$		0.533755	$a_1 = (c1)(fco22x)/\log 2$
$a_2$	<i>c1 × c3</i>	0.008844	
$a_3$	<i>c4</i>	0.025	
$b_0$	<i>b12</i>	0.12	
$b_1$	<i>b21</i>	0.196	
$b_2$	<i>b32</i>	0.001465	
$b_3$	<i>b23</i>	0.007	
<i>Additional parameters</i>			
$\eta_0$		-3.329093	
$\eta_1$		0.846811	$\eta_1 = 1 - a_0 - a_3$
$\eta_2$		0.002984	$\eta_2 = (a_0 - a_2)a_3$
$\eta_3$		0.533755	$\eta_3 = a_1$
$\eta_4$		0.013344	$\eta_4 = a_1a_3$
$\phi_0$		0.876134	
$\phi_{00}$		263.283680	
$\phi_1$		0.675535	$\phi_1 = 1 - b_0 - b_1 - b_2 - b_3$
$\phi_2$		0.001303	$\phi_2 = b_0b_2 + b_0b_3 + b_1b_2$
$\lambda_1$		0.795535	$\lambda_1 = 1 - b_1 - b_2 - b_3$
$\lambda_2$		0.000287	$\lambda_2 = b_1b_2$