

OPE between primary fields

14.12.11 (3/1)

generally: $\phi_1(z) \phi_2(w) = \sum_{p \in \mathcal{I}} \sum_{\substack{1 \leq k_1 \leq k_2 \leq \dots \leq k_p \\ 1 \leq \bar{k}_1 \leq \bar{k}_2 \leq \dots \leq \bar{k}_p \\ s \geq 0}} C_{12}^{p, \{k_i, \bar{k}_i\}} (z-w)^{-h_1-h_2+h_p+K} (\bar{z}-\bar{w})^{-\bar{h}_1-\bar{h}_2+\bar{h}_p+K} \phi_p(w)$

Annotations:
 - $C_{12}^{p, \{k_i, \bar{k}_i\}}$: from scaling (\mathbb{Z}_0) symmetry
 - $\phi_p(w)$: descendant of ϕ_p at w
 - \sum_{\dots} : sum over descendants

$C_{12}^{p, \{k_i, \bar{k}_i\}} = C_{12p} \beta_{12}^{p, \{k_i\}} \beta_{12}^{p, \{\bar{k}_i\}}$

Annotations:
 - $\beta_{12}^{p, \{k_i\}}$: universal rational functions of C, h_1, h_2, h_p , parameterized by partitions $\{k_i\}$
 - $\beta_{12}^{p, \{k_i\}} = 1$: normalization

Rem: ϕ_p can only appear in $\phi_1 \phi_2$ OPE if ϕ_p appears.

4-point correlator of primary fields

$\langle \phi_1 \phi_2 \phi_3 \phi_4 \rangle = \dots \cdot f(x, \bar{x})$

Annotation: $f(x, \bar{x})$ is cross-ratio of z_1, z_2, z_3, z_4

$f(x, \bar{x}) = \langle \phi_1 | \phi_2(z) \phi_3(x, \bar{x}) | \phi_4 \rangle$

$$\sum_p \sum_{\{k_i, \bar{k}_i\}} C_{34p} \beta_{34}^{p, \{k_i\}} \beta_{34}^{p, \{\bar{k}_i\}} x^{-h_3-h_4+h_p+K} \bar{x}^{-\bar{h}_3-\bar{h}_4+\bar{h}_p+K} |\phi_p, \{k_i, \bar{k}_i\}\rangle$$

Annotation: $|\phi_p, \{k_i, \bar{k}_i\}\rangle$ is descendant state of $|\phi_p\rangle$

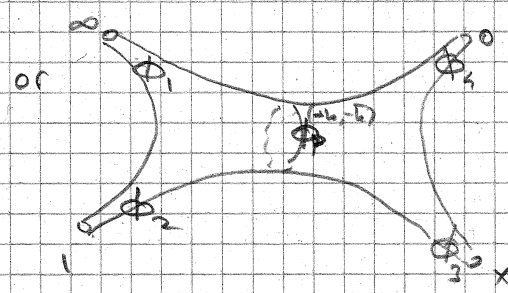
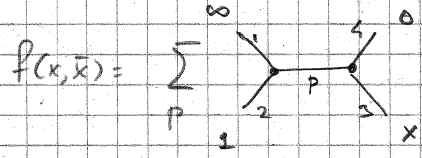
$f(x, \bar{x}) = \sum_{p \in \mathcal{I}} C_{12p} C_{34p} F_{34}^{21}(p|x) F_{34}^{21}(p|\bar{x})$

where $F_{34}^{21}(p|x) = x^{-h_3-h_4+h_p} \sum_{\{k_i\}} x^K \beta_{34}^{p, \{k_i\}} \gamma_{12}^{p, \{k_i\}}$

Annotation: "conformal block for 4-pt function"

where γ 's come from $\langle \phi_1 | \phi_2(z) | \phi_p, \{k_i, \bar{k}_i\} \rangle = \gamma_{12}^{p, \{k_i\}} \beta_{12}^{p, \{k_i\}}$

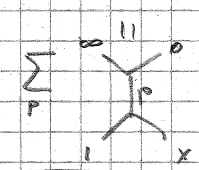
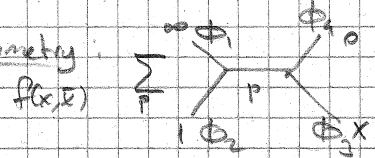
Pictures



$\gamma_{12}^{p, \{k_i\}}$ related to $\beta_{12}^{p, \{k_i\}}$ via $\gamma_{12}^{p, \{k_i\}} = \sum_{\{l_i\}} G_{\{l_i\}, \{k_i\}}^{p, \{k_i\}} \beta_{12}^{p, \{l_i\}}$

Annotation: Gram matrix

Crossing Symmetry



$$\sum_p C_{12p} C_{34p} F_{34}^{21}(p|x) F_{34}^{21}(p|\bar{x}) = \sum_p C_{23p} C_{14p} F_{32}^{14}(p|1-x) F_{32}^{14}(p|1-\bar{x})$$

- a set of quadratic equations on C_{ijk}

More generally, for $\langle \phi_1 \dots \phi_n \rangle$,

one constructs a basis of conf. blocks for any 2-valent tree with leaves labelled by ϕ_1, \dots, ϕ_n . Basis elements are assoc. to decorations of edges by primary fields.

transitions between bases for different trees \leftrightarrow "crossing symmetry"

Null-vectors → diff. equations on correlators

If $V_{c,h}$ contains a null-vector $|x\rangle$ at level N ,

$|x\rangle = \sum_{\{k_j\}} a_{\{k_j\}} \hat{L}_{-k_1} \dots \hat{L}_{-k_n} |h\rangle$, then correlators containing $\phi_{h,h}$ (corresponding primary) satisfy PDEs of degree $\leq N$:
coeffs of the null-vector in standard basis in $V_{c,h}$

$$\left(\sum_{\{k_j\}} a_{\{k_j\}} \hat{L}_{-k_1} \dots \hat{L}_{-k_n} \right) \phi_{h,h} = 0 \Rightarrow 0 = \langle \psi(\phi_{h,h})(z_1, \bar{z}_1) \cdot \phi_2(z_2, \bar{z}_2) \dots \phi_n(z_n, \bar{z}_n) \rangle$$

$\psi \in U(N, \mathbb{R})$

$$= \left(\sum_{\{k_j\}} a_{\{k_j\}} \hat{L}_{-k_1} \dots \hat{L}_{-k_n} \right) \langle \phi_{h,h}(z_1, \bar{z}_1) \phi_2 \dots \phi_n \rangle$$

a diff. operator in z_2, \dots, z_n

where $\hat{L}_{-k} = \sum_{\ell=2}^n \frac{1}{(z_0 - z_\ell)^{k-1}} \frac{\partial}{\partial z_0} + \frac{k-1}{(z_0 - z_1)^k} \frac{\partial}{\partial z_1}$

Ex A null-vector at level 2 in $V_{c,h}$ occurs iff $(h=h_{1,2}(c)$ or $h_{2,1}(c)$)

and has the form $|x\rangle = \left(L_{-2} - \frac{3}{2(2h+1)} L_{-1}^2 \right) |h\rangle$

$h = \frac{5-c \pm \sqrt{(c-1)(25-c)}}{16}$

Corresponding diff. eq.:

$$\left(L_{-2} - \frac{3}{2(2h+1)} L_{-1}^2 \right) \langle \psi(z_1, \bar{z}_1) \cdot \phi_2(z_2, \bar{z}_2) \dots \phi_n(z_n, \bar{z}_n) \rangle = 0$$

explicitly: $\left(\sum_{\ell=2}^n \left(\frac{1}{z-z_\ell} \frac{\partial}{\partial z} + \frac{h_\ell}{(z-z_\ell)^2} \right) - \frac{3}{2(2h+1)} \frac{\partial^2}{\partial z^2} \right) \langle \psi \phi_2 \dots \phi_n \rangle = 0$

$n=2 \rightarrow$ a tautologically satisfied equation

$n=3 \rightarrow$ a quadratic relation between $h_1, h_2, h_3 \rightarrow$ selection rule for $\phi_1 \cdot \phi_2$ OPE ("fusion rules")

$n=4 \rightarrow$ hypergeometric ODE on $f(x, \bar{x})$ - non-triv. part of $\langle \psi \phi_2 \phi_3 \phi_4 \rangle$, depending on cross-ratio

Ex: 4-spin correlator in Ising model

$$G^{(4)} = \langle \sigma(z_1, \bar{z}_1) \dots \sigma(z_4, \bar{z}_4) \rangle = \left(\frac{z_{13} z_{24}}{z_{12} z_{34} z_{14} z_{23}} \right)^{1/8} \left(\frac{z_{13} z_{24}}{z_{12} z_{34} z_{14} z_{23}} \right)^{1/8} F(x, \bar{x})$$

$(\frac{1}{16}, \frac{1}{16})$ -field, $c = \frac{1}{2} = \bar{c}$

from global conformal invariance

null-vector at level 2 for $V_{\frac{1}{2}, h_{1,2} = \frac{1}{16}} \Rightarrow$

$$\Rightarrow \left(x(1-x) \frac{\partial^2}{\partial x^2} + \left(\frac{1}{2} - x \right) \frac{\partial}{\partial x} + \frac{1}{16} \right) F(x, \bar{x}) = 0$$

solutions: $f_{1,2}(x) = (1 \pm \sqrt{1-x})^{1/2}$ ← conformal blocks for 4-spin function

$$G^{(4)} = \left| \frac{z_{13} z_{24}}{z_{12} z_{34} z_{14} z_{23}} \right|^{1/8} \sum_{i,j=1}^2 a_{ij} f_i(x) f_j(\bar{x})$$

$a(1f_1(x)^2 + 1f_2(x)^2)$

due to single-valuedness of $G^{(4)}$ when $\bar{x} = x^*$

coefficient $a = \frac{1}{2}$

determined from OPE $\sigma(z_1, \bar{z}_1) \sigma(z_2, \bar{z}_2) \sim \frac{1}{z_{12}^{1/8} \bar{z}_{12}^{1/8}} + \dots$ (from normalization of σ wrt 2-pt function)

Fusion rules

if $h = h_{p,q}$, then there is an algebraic equation on conf. weights h', h'' which can appear in $\phi \cdot \phi' = \sum_{\phi''} \sum_{\{h', h''\}} \phi''(h', h'')$ OPE - from degree pq diff. equation on 3-pt. function

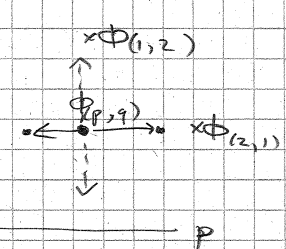
(notation: $\phi \times \phi' = \sum_{\phi'' \in \mathcal{I}' \subset \mathcal{I}} \phi''$)
 ↑ all primaries allowed conf. parameters

E.g. 1) if $h = h_{1,1} = 0$ then equation is $h' = h''$ (indeed, $\mathbb{1} \times \phi' = \phi'$)

2) if $h = h_{2,1}$ then (quadratic equation on h', h'') $\Leftrightarrow (\alpha'' = \alpha' \pm \alpha_{\pm})$
 with substitution $h = \frac{\alpha^2 - 1}{4m(m+1)}$, $\alpha_{\pm} = m \pm 1$, $m = m(c)$
 if $h = h_{1,2}$ ----- $\alpha'' = \alpha' \pm \alpha_{\pm}$, $\alpha_{\pm} = m$

In minimal models:

denote $\phi_{(p,q)}$ the primary field with $h = h_{p,q}$
 $\phi_{(3,2)} \times \phi_{(p,q)} = \phi_{(p,q-1)} + \phi_{(p,q+1)}$
 $\phi_{(2,1)} \times \phi_{(p,q)} = \phi_{(p-1,q)} + \phi_{(p+1,q)}$



- combining these two, we may arrive to more restrictive fusion rules:

E.g. $\phi_{(1,2)} \times \phi_{(2,1)} = \phi_{(2,0)} + \phi_{(2,2)}$
 $\phi_{(2,1)} \times \phi_{(1,2)} = \phi_{(0,2)} + \phi_{(2,2)}$
 $\Rightarrow \phi_{(1,2)} \times \phi_{(2,1)} = \phi_{(2,2)}$

Generally: for degenerate primary fields $\phi_{(p,q)} \times \phi_{(p',q')} = \sum_{\substack{p''=1+p-p', \\ p+p'+p''=1 \text{ mod } 2}}^{p+p'-1} \sum_{\substack{q''=1+q-q', \\ q+q'+q''=1 \text{ mod } 2}}^{q+q'-1} \phi_{(p'',q'')} \quad (*)$

- only degenerate fields with $p'', q'' \geq 1$ occur in table "truncation of operator algebra"

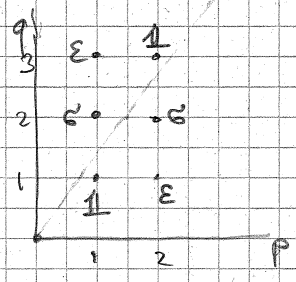
in minimal model $\mathcal{M}(P, Q)$, we additionally have $h_{p,q} = h_{P-p, Q-q}$

which results in fusion rule:

$$\phi_{(p,q)} \times \phi_{(p',q')} = \sum_{p''=1+p-p'}^{\min(p+p'-1, 2P-p-p'-1)} \sum_{q''=1+q-q'}^{\min(q+q'-1, 2Q-q-q'-1)} \phi_{(p'',q'')}$$

\rightarrow fields in Kac table form a closed algebra under OPE

E.g. for $\mathcal{M}(3,4)$:
 $\mathbb{1} := \phi_{(1,1)} = \phi_{(2,3)} \quad h=0$
 $\sigma := \phi_{(1,2)} = \phi_{(2,2)} \quad h=1/6$
 $\varepsilon := \phi_{(2,1)} = \phi_{(3,3)} \quad h=1/2$



Fusion rules:

\times	$\mathbb{1}$	σ	ε
$\mathbb{1}$	$\mathbb{1}$	σ	ε
σ	σ	$\mathbb{1} + \varepsilon$	σ
ε	ε	σ	$\mathbb{1}$

Vertex algebras

def a vertex algebra is the following data:

- a vector space V (space of states)
- "vacuum vector" $|0\rangle \in V$
- "translation operator" $T \in \text{End}(V)$
- state-field correspondence $Y(\cdot, z): V \rightarrow \text{End}(V)[[z^{\pm 1}]]$

formal Laurent series

$$A \mapsto Y(A, z) = \sum_{n \in \mathbb{Z}} A_{(n)} z^{-n-1}$$

$\text{End}(V)$

Axioms:

(vacuum axiom)

$$Y(|0\rangle, z) = \text{Id}_V$$

$$\forall A \in V \quad Y(A, z)|0\rangle \in V[[z]] \quad , \quad Y(A, z)|0\rangle|_{z=0} = A$$

- no negative powers

(translation axiom) $\forall A \in V$

$$[T, Y(A, z)] = \partial_z Y(A, z)$$

$$T|0\rangle = 0$$

(locality) $\forall A, B \in V$, $Y(A, z)$ and $Y(B, w)$ are "mutually local", i.e.

$$\forall u \in V^* \quad \langle u, Y(A, z)Y(B, w)v \rangle \quad \text{and} \quad \langle u, Y(B, w)Y(A, z)v \rangle$$

are expansions of the same element

$$F_{u, A, B, v} \in \mathbb{C}[[z, w]][z^{-1}, w^{-1}, (z-w)^{-1}]$$

in $\mathbb{C}((z))((w))$ and $\mathbb{C}((w))((z))$ resp.; order of pole at $z=w$ is uniformly bounded $\forall u, v$

or equivalently: $\exists N \in \mathbb{N}$ s.t. $(z-w)^N Y(A, z)Y(B, w) = (z-w)^N Y(B, w)Y(A, z)$ (*)

Rem: it follows from def. that

$$T(A) = A_{(-2)}|0\rangle$$

$$\frac{T^n}{n!}(A) = A_{(-n-1)}|0\rangle \rightsquigarrow Y(A, z)|0\rangle = e^{zT}(A)$$

(trivial) example: Holomorphic $\mathbb{C}V, A_0$:

$$(N=0 \text{ in } (*) \forall A, B) \Rightarrow Y(A, z) \cdot B = Y(A, z)Y(B, w)|0\rangle|_{w=0} = Y(B, w)Y(A, z)|0\rangle|_{w=0} \in \mathbb{C}[[z]]$$

- all vertex operators $Y(A, z)$ are regular at $z=0$

(Holomorphic V.A.) \iff (commutative unital algebras with derivation)

$$(V, |0\rangle, T, Y) \longmapsto V, \quad A \circ B = \underbrace{Y(A, 0)}_{=A_{-1}} \cdot B \quad \text{- commutative, associative}$$

$|0\rangle$ -unit, T -derivation

$$Y(A, z) = e^{zT}(A) \circ \bullet$$

def V.A. is called $\mathbb{Z}_{\geq 0}$ -graded if $V = \bigoplus_{n \geq 0} V_n$,
 $A \in V_k, B \in V_p \Rightarrow A_{(n)}(B) \in V_{k+p-n-1}$

def conformal V.A. = $\mathbb{Z}_{\geq 0}$ -graded V.A. equipped with a "conformal vector" $\omega \in V_2$ s.t. coefficients in

$$Y(\omega, z) = \sum_{n \in \mathbb{Z}} L_n z^{-n-2} \quad \text{s.t.}$$

- $\forall A \in V_n, L_0 A = n A$
- $L_{-1} = T$
- $[L_m, L_n] = \dots$

Virasoro comm. relation

$Y(\omega, z)$ is called "energy-momentum field"

Example: vertex algebra associated to Heisenberg Lie algebra:

$V = V_0^{\text{Heis}}$ - Fock space for Heisenberg Lie algebra $[a_n, a_m] = n \delta_{n,-m} \mathbb{1}$
 $a_0 |0\rangle = 0$

$$Y(a_{-k_1} \dots a_{-k_n} |0\rangle, z) = \frac{1}{(k_1-1)! \dots (k_n-1)!} : \partial^{k_1} \varphi(z) \dots \partial^{k_n} \varphi(z) :$$

where $\varphi(z) = -\sum_{n \in \mathbb{Z}} \frac{a_n z^{-n}}{n}$

$$\omega = (a_{-1})^2 |0\rangle$$

$$Y(\omega, z) = \frac{1}{z} : \partial \varphi(z)^2 :$$

$$T = L_{-1} = \frac{1}{2} \sum_{n \in \mathbb{Z}} a_n a_{-n}$$

e.g. $Y(a_{-1} |0\rangle, z) = \partial \varphi(z)$

Example: Virasoro vertex algebra

$$V = \text{Span}(L_{-k_1} \dots L_{-k_n} |0\rangle)_{k_i \geq 2} = V_{c,0} / V_{c,1}$$

$$Y(L_{-k_1} \dots L_{-k_n} |0\rangle, z) = \frac{1}{(k_1-2)! \dots (k_n-2)!} : \partial^{k_1-2} L(z) \dots \partial^{k_n-2} L(z) :$$

where $L(z) = \sum_{n \in \mathbb{Z}} L_n z^{-n-2} = Y(L_{-2} |0\rangle, z)$
 - stress-energy tensor